Ordinary Differential EquationsODEs Part 2

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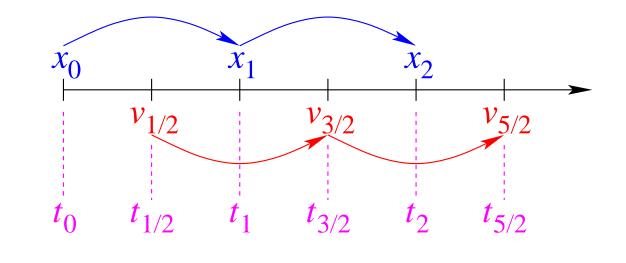
The Leapfrog Integrator

- Very useful for second-order DEs in which $d^2 \mathbf{x}/dt^2 = \mathbf{f}(\mathbf{x})$, e.g., SHM, N-body, etc.
 - NOTE: Now dropping the prime (') from f...
- Suppose x is position, so d^2x/dt^2 is acceleration.
- Procedure: define $\mathbf{v} = d\mathbf{x}/dt$ at the *midpoints* of the steps, i.e., velocities staggered wrt positions.

• Set
$$\mathbf{v}_{n+1/2} = \mathbf{v}(t_n + \frac{1}{2}h)$$
.

• Then advance \mathbf{x}_n to \mathbf{x}_{n+1} and $\mathbf{v}_{n+1/2}$ to $\mathbf{v}_{n+3/2}$:

$$\mathbf{x}_{n+1} = \mathbf{x}_n + h\mathbf{v}_{n+1/2},$$
$$\mathbf{v}_{n+3/2} = \mathbf{v}_{n+1/2} + h\mathbf{f}(\mathbf{x}_{n+1}).$$



Complication: need to "jumpstart" and "resync"...

$$\begin{aligned} \mathbf{v}_{n+1/2} &= \mathbf{v}_n + (h/2) \mathbf{f}(\mathbf{x}_n) & \text{[opening "kick": Euler]} \\ \mathbf{x}_{n+1} &= \mathbf{x}_n + h \mathbf{v}_{n+1/2} & \text{["drift"]} \\ \mathbf{v}_{n+1} &= \mathbf{v}_{n+1/2} + (h/2) \mathbf{f}(\mathbf{x}_{n+1}) & \text{[closing "kick": resync]} \end{aligned}$$

• Note
$$\mathbf{v}_{n+3/2} = \mathbf{v}_{n+1} + (h/2)\mathbf{f}(\mathbf{x}_{n+1}) = \mathbf{v}_{n+1/2} + h\mathbf{f}(\mathbf{x}_{n+1}).$$

- Also have "drift-kick-drift" (DKD) scheme.
- Like midpoint method, Leapfrog is <u>second order</u>:

$$\mathbf{x}(t+h) = \mathbf{x}(t) + h\mathbf{v}(t+h/2),$$

but

$$\mathbf{v}(t+h/2) = \mathbf{v}(t) + (h/2)\mathbf{f}(t) + \mathcal{O}(h^2).$$

Therefore

$$\mathbf{x}(t+h) = \mathbf{x}(t) + h\mathbf{v}(t) + (h^2/2)\mathbf{f}(t) + \mathcal{O}(h^3).$$



- So why is Leapfrog so great?...
- Answer: Leapfrog is time reversible.
- Suppose we step back from $(t_{n+1}, \mathbf{x}_{n+1}, \mathbf{v}_{n+3/2})$ to $(t_n, \mathbf{x}_n, \mathbf{v}_{n+1/2})$. Applying the algorithm:

$$\mathbf{v}_{n+1/2} = \mathbf{v}_{n+3/2} + (-h)\mathbf{f}(\mathbf{x}_{n+1}),$$
$$\mathbf{x}_n = \mathbf{x}_{n+1} + (-h)\mathbf{v}_{n+1/2}.$$

- These are precisely the steps (in reverse) that we took to advance the system in the first place!
- Hence if we Leapfrog forward in time, then reverse to t = 0, we're back to where we started, *precisely*.

- Leapfrog is time reversible because of the symmetric way in which it is defined, unlike the other schemes.
 - In Euler, forward and backward steps do not cancel since they use different derivatives at different times.
 - In Midpoint, the estimate of the derivative is based on an extrapolation from the left-hand side of the interval. On time reversal, the estimate would be based on the right-hand side, not the same.
 - Similarly, RK4 is not time reversible.
- Time reversibility is important because it guarantees conservation of energy, angular momentum, etc. (in many cases).

Suppose the integrator makes an error ε after one orbital period. Now reverse. Is the error -ε? No! The time-reversed orbit is a solution of the original ODE (with v replaced with -v), so the energy error is still +ε. But we've returned to our starting point, so we know the final energy error is zero.
Hence ε = 0!

- Leapfrog is only second order, but very stable.
- Leapfrog is an example of a class of "symplectic" integrators that conserve phase-space volume: exactly solves an approximate Hamiltonian system.

$$H = H_D + H_K + H_{\text{err}} = \frac{1}{2}v^2 + V(\mathbf{r}) + H_{\text{err}},$$

or D(h/2)K(h)D(h/2), with $H_{err} \sim O(h^3)$. You can also construct the usual kick-drift-kick scheme, K(h/2)D(h)K(h/2), because the Hamiltonian is separable.

Adaptive Stepsize Control

- \checkmark Up to now, have assumed stepsize h is constant.
- Clearly prefer choosing h small when |f'| is large, and h large when |f'| is small. (We've reintroduced the prime (') notation, just to be confusing...)
- The tradeoff is extra trial steps to determine optimum h, but may achieve factor of 10 to 100 increase in stepsize, so it's often worth it.
- NRiC provides a routine odeint() for RK4 with adapative stepsize control. Complicated to use!