Local Dark Matter

Konrad Kuijken Leiden Observatory, Leiden University



Leiden Observatory

Galaxy haloes



flat rotation curves: galaxies are not WYSIWYG



Rotation curves

 $\frac{GM(< r)}{r^2} \simeq \frac{V_{\rm circ}^2}{r}$

Even in very flattened systems, the rotation curve mostly probes the enclosed mass (monopole) of the mass distribution

The example: exponential disk $\Sigma = \Sigma_0 e^{-r/h}$



In finite systems, Keplerian falloff V=r^{-1/2} inevitable

Monday, July 16, 2012

eiden Observatory

Rotation curve decomposition





Assign fraction of potential to 'visible' components
'Maximum disk' decomposition ('minimum extra')
Requires knowledge of Mass-to-Light ratio
WYSIWYG models not feasible in most cases

Stellar Mass Function



Most light emitted
 by most massive
 stars

Most of the mass is 'invisible'

Starlight is a poor tracer of stellar mass



Monday, July 16, 2012

-eiden Observatory

Direct measurements of M/L in disk galaxies

Opposition of self-gravitating thin disks gravity field ~ vertical ø vertical motions decouple from horizontal ones \varnothing stars move in 1-D vertical potential $\Psi(z)$ $\nabla^2 \psi = 4\pi G \rho_m \Rightarrow \frac{d^2 \psi(z)}{dz^2} = 4\pi G \rho_m \Rightarrow \frac{d\psi(z)}{dz} = 2\pi G \Sigma_m (\langle |z|)$

Direct measurements of M/L in disk galaxies

Essentially compare mean orbital v_z^2 at z=0 with

Ineed disk thickness (edge-on view)

velocity dispersion (face-on view)

Leiden Observatory 0 Problem in practice:

mean z reached



Solar Neighbourhood

X

☆

Can measure 3D structure from star counts

Can measure dynamics from individual stars

Kapteyn (1922) – Oort (1932)



eiden Observatory

Kapteyn 1922

FIRST ATTEMPT AT A THEORY OF THE ARRANGEMENT AND MOTION OF THE SIDEREAL SYSTEM¹

BY J. C. KAPTEYN²



the relative velocity is also in the plane of the Milky way and about 40 km/sec. It is incidentally suggested that when the theory is perfected it may be possible to determine the amount of dark matter from its gravitational effect. (5) The chief defects

Leiden Observatory

Oort 1932

BULLETIN OF THE ASTRONOMICAL INSTITUTES OF THE NETHERLANDS.

1932 August 17

Volume VI.

No. 238.

COMMUNICATION FROM THE OBSERVATORY AT LEIDEN.

The force exerted by the stellar system in the direction perpendicular to the galactic plane and some related problems, by \mathcal{F} . H. Oort.

(14)

11. The amount of dark matter. From the results found for the decrease of K(z) with z we may derive an approximate value of the total density of matter, Δ , in the neighbourhood of the sun. Let us suppose that we are situated inside a homogeneous ellipsoid of revolution with semi-axes a and c, and density Δ . For z = 0 there will then be the following relation:

$$\partial K(z)/\partial z = -4\pi\gamma x\Delta$$

Case	Δ
I a)	.108
I b)	.093
II a)	·089
II b).	·079

total mass density (M_o/pc³)

Stellar dynamics The Describe stars with phase space density $f(x_i, v_j)$ $\int f d^3 v = \rho; \qquad \int v_z^2 f d^3 v = \rho \sigma_z^2;$ Collisionless Boltzmann equation $\frac{Df}{Dt} = \frac{\partial f}{\partial t} + v_i \frac{\partial f}{\partial x_i} - \frac{\partial \psi}{\partial x_i} \frac{\partial f}{\partial v_i} = 0$ Ignore all x,y gradients $v_z \frac{df}{dz} = -K_z \frac{df}{dy}$

x v_z, integrate over v_z:

 $\frac{d(\rho\sigma_z^2)}{dz} = K_z\rho$

Monday, July 16, 2012

-eiden Observatory

Stellar dynamics

More generally, for an axisymmetric system in cylindrical polar coordinates (R, ϕ, z) $\frac{Df}{Dt} = \frac{\partial f}{\partial t} + v_R \frac{\partial f}{\partial R} + v_z \frac{\partial f}{\partial z} + \left(K_R + \frac{v_\phi^2}{R}\right) \frac{\partial f}{\partial v_R} - \frac{v_R v_\phi}{R} \frac{\partial f}{\partial v_\phi} + K_z \frac{\partial f}{\partial v_z} = 0$ x v_i, integrate gives Jeans equations $\frac{1}{R}\frac{\partial(R\rho\sigma_R^2)}{\partial R} + \frac{\partial(\rho\sigma_{Rz})}{\partial z} - \frac{\rho(v_\phi^2 + \sigma_\phi^2)}{R} = \rho K_R$ $\frac{\partial(\rho\sigma_z^2)}{\partial z} + \frac{\partial(\rho\sigma_{Rz})}{\partial R} = \rho K_z$

Solve NB: in a thermalized system velocity dispersion is isotropic: $\sigma_{Rz}=0$; $\sigma_{R}=\sigma_{z}=\sigma_{\varphi}$

Monday, July 16, 2012

-eiden Observatory

Stellar dynamics

More generally, for an axisymmetric system in cylindrical polar coordinates (R, ϕ, z) $\frac{Df}{Dt} = \frac{\partial f}{\partial t} + v_R \frac{\partial f}{\partial R} + v_z \frac{\partial f}{\partial z} + \left(K_R + \frac{v_\phi^2}{R}\right) \frac{\partial f}{\partial v_R} - \frac{v_R v_\phi}{R} \frac{\partial f}{\partial v_\phi} + K_z \frac{\partial f}{\partial v_z} = 0$ x v_i, integrate gives Jeans equations $\frac{1}{R}\frac{\partial(R\rho\sigma_R^2)}{\partial R} + \frac{\partial(\rho\sigma_{Rz})}{\partial z} - \underbrace{\begin{array}{c}\rho(v_{\phi}^2 + \sigma_{\phi}^2)\\R\end{array}}_{R} = \rho K_R$ $\frac{\partial(\rho\sigma_z^2)}{\partial z} + \frac{\partial(\rho\sigma_{Rz})}{\partial R} = \rho K_z$

eiden Observatory

Solution NB: in a thermalized system velocity dispersion is isotropic: $\sigma_{Rz}=0$; $\sigma_{R}=\sigma_{z}=\sigma_{\varphi}$

Stellar dynamics

More generally, for an axisymmetric system in cylindrical polar coordinates (R, ϕ, z) $\frac{Df}{Dt} = \frac{\partial f}{\partial t} + v_R \frac{\partial f}{\partial R} + v_z \frac{\partial f}{\partial z} + \left(K_R + \frac{v_\phi^2}{R}\right) \frac{\partial f}{\partial v_R} - \frac{v_R v_\phi}{R} \frac{\partial f}{\partial v_\phi} + K_z \frac{\partial f}{\partial v_z} = 0$ $\frac{1}{R}\frac{\partial(R\rho\sigma_R^2)}{\partial R} + \frac{\partial(\rho\sigma_{Rz})}{\partial z} - \underbrace{\frac{\rho(v_\phi^2 + \sigma_\phi^2)}{R} = \rho K_R}_{R}$ $\underbrace{\frac{\partial(\rho\sigma_z^2)}{\partial z}}_{\partial z} + \underbrace{\frac{\partial(\rho\sigma_{Rz})}{\partial R}}_{\partial R} \neq \overbrace{\rho K_z}^{\partial K_z}$

Solution NB: in a thermalized system velocity dispersion is isotropic: $\sigma_{Rz}=0$; $\sigma_{R}=\sigma_{z}=\sigma_{\varphi}$

Monday, July 16, 2012

eiden Observatory

Vertical kinematics

 $v_z \frac{df}{dz} + \frac{d\psi(z)}{dz} \frac{df}{dv_z} = 0 \quad \Rightarrow \quad f(z, v_z) = f(\psi(z) + v_z^2) = f(E_z)$

☆

Velocity distribution at z=0 + ψ (z)
 ==> velocity distribution (and density) at height z

 (avoids derivatives needed to apply the Jeans equation)

 $\frac{d(\rho\sigma_z^2)}{dz} = K_z\rho$

eiden Observatory

In practice

Pick a good tracer population that probes the disk (z < 1 kpc)

numerous

sufficiently old, well-mixed

well-calibrated distances

good radial velocity measurements

Lower main sequence stars (G – K dwarfs)

Parameterize possible potentials

Known star populations + gas + dark disk + halo

Monday, July 16, 2012

Leiden Observatory

Trial potentials



Universitett Leiden

predicted ρ(z) in a WYSIWYG potential

Example (KK+Gilmore 1989)



UPGREN K GIANT SAMPLE



Disk dark matter?

Since 1990's, most analyses agree:

disk kinematics below
 z=500pc consistent with
 known populations' gravity

 disk volume density at z=0
 is about 0.1M₀/pc³, 50-50

 stars and gas

disk surf. density 50M_☉/pc²







Garbari et al 2012

Monday, July 16, 2012

-eiden Observatory

Halo dark matter

Milky Way rotation is a bit awkward to model

Fix local disk surf. density

Vary scale height, solar orbit speed, fit V_{circ} curve









Using disk stars to probe halo dark matter

 $\frac{1}{R}\frac{\partial(R\rho\sigma_{R}^{2})}{\partial R} + \frac{\partial(\rho\sigma_{Rz})}{\partial z} - \frac{\rho(v_{\phi}^{2} + \sigma_{\phi}^{2})}{R} = \rho K_{R}$ $\frac{\partial(\rho\sigma_{z}^{2})}{\partial z} + \frac{\partial(\rho\sigma_{Rz})}{\partial R} = \rho K_{z} \qquad (A\pi G\rho_{m} = -\frac{\partial K_{z}}{\partial z} - \frac{1}{R}\frac{\partial(RK_{R})}{\partial R}$

At high z, no longer vertical-only kinematics
Velocity ellipsoid tilts, is anisotropic
Need to measure radial gradient of this tilt
Also Poisson equation not separable

Monday, July 16, 2012

eiden Observatory

Orbit tilt

 $The \ tilt \ of \ the \ velocity \ ellipsoid \ from \ RAVE$



Monday, July 16, 2012

Leiden Observatory

Orbit tilt

 $\frac{\partial(\rho\sigma_z^2)}{\partial z} + \frac{\partial(\rho\sigma_{Rz})}{\partial R} = \rho K_z$

For a 2:1 radial:vertical velocity dispersion ratio $\sigma_{Rz} \simeq (\sigma_R^2 - \sigma_z^2) \frac{z}{R}$

 $F_h \frac{z}{R} K_R \simeq F_h \frac{z}{R^2} V_{\text{circ}}^2$

So Radial gradient set by disk scale length H $(\rho\sigma_{Rz})_{,R} \simeq (\sigma_R^2 - \sigma_z^2) \frac{2z}{HR}$

Ø Vertical force due to halo is

σ_z≈40km/s V_{circ}≈200km/s H≈R/2



-eiden Observatory

Ratio halo K_z : 'fake tilt force'

 $F_h \frac{V_{\rm circ}^2}{4\sigma_z^2} \frac{H}{2R} \simeq F_h$

Moni Bidin et al. 2012

Kinematical and chemical vertical structure of the Galactic thick disk^{1,2} II. A lack of dark matter in the solar neighborhood

Study of the vertical kinematics using red giant stars up to 4kpc from the disk plane.

Large number of assumptions on missing terms; reasonable for disk population but uncertain

o criticised by Bovy & Tremaine (2012)



Bovy & Tremaine 2012



Bovy & Tremaine 2012



Leiden Observatory

Moni Bidin et al. 2012

Rotational velocity profile from MB & al Paper I
 v_{rot} < 0.5v_{circ} ! Are these disk stars?



Monday, July 16, 2012

Leiden Observatory

Not discussed here

Halo flattening: barely affects rotation curve, but increases local density

ø prolate halo would lower it

 Local substructure: could lead to higher or lower densities

Overall dark halo structure (size, extent, shape)



Summary

Solution Long history of measuring local DM density

No convincing evidence for 'cold' DM component in the disk (as thin as a few 100 pc)

Halo density derived from vertical kinematics unreliable:

⊘ velocity ellipsoids tilt, radial Poisson eq. terms
 ⊘ Needs a full picture of 3-D kinematics (Gaia!)
 ⊘ Rotation curve decompositions lead to estimates of 0.005 - 0.01 M_☉/pc³.

Universitett L