

Stellar Populations

Physics of Galaxies 2012

part 3



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Why do we need to know about stars to study galaxies?

- ✦ We can't see the *vast majority of the mass* in galaxies
 - ✦ dark matter!
- ✦ We only see the “baryons” in galaxies
 - ✦ stars
 - ✦ gas
- ✦ To understand the properties of galaxies, we need to study these components



Stellar atmospheres



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- When we look at galaxies in the UV, visible, and NIR, we see (mostly) the atmospheres of their stars
 - their continua: blackbody curves
 - their absorption lines
 - ...and sometimes the emission lines of the gas, if there is star formation or an active galactic nucleus



- Stellar spectra depend mainly on **three variables**:
 - *Effective temperature*, where $L \equiv 4\pi R^2 \sigma T_{\text{eff}}^4$
 - To first order, the emitted continuum of a star is a blackbody curve with $T \sim T_{\text{eff}}$
 - and a peak at $\lambda_{\text{max}} = [2.9/T \text{ (K)}] \text{ mm}$



- *Composition*: more-or-less the mean abundance of elements heavier than He

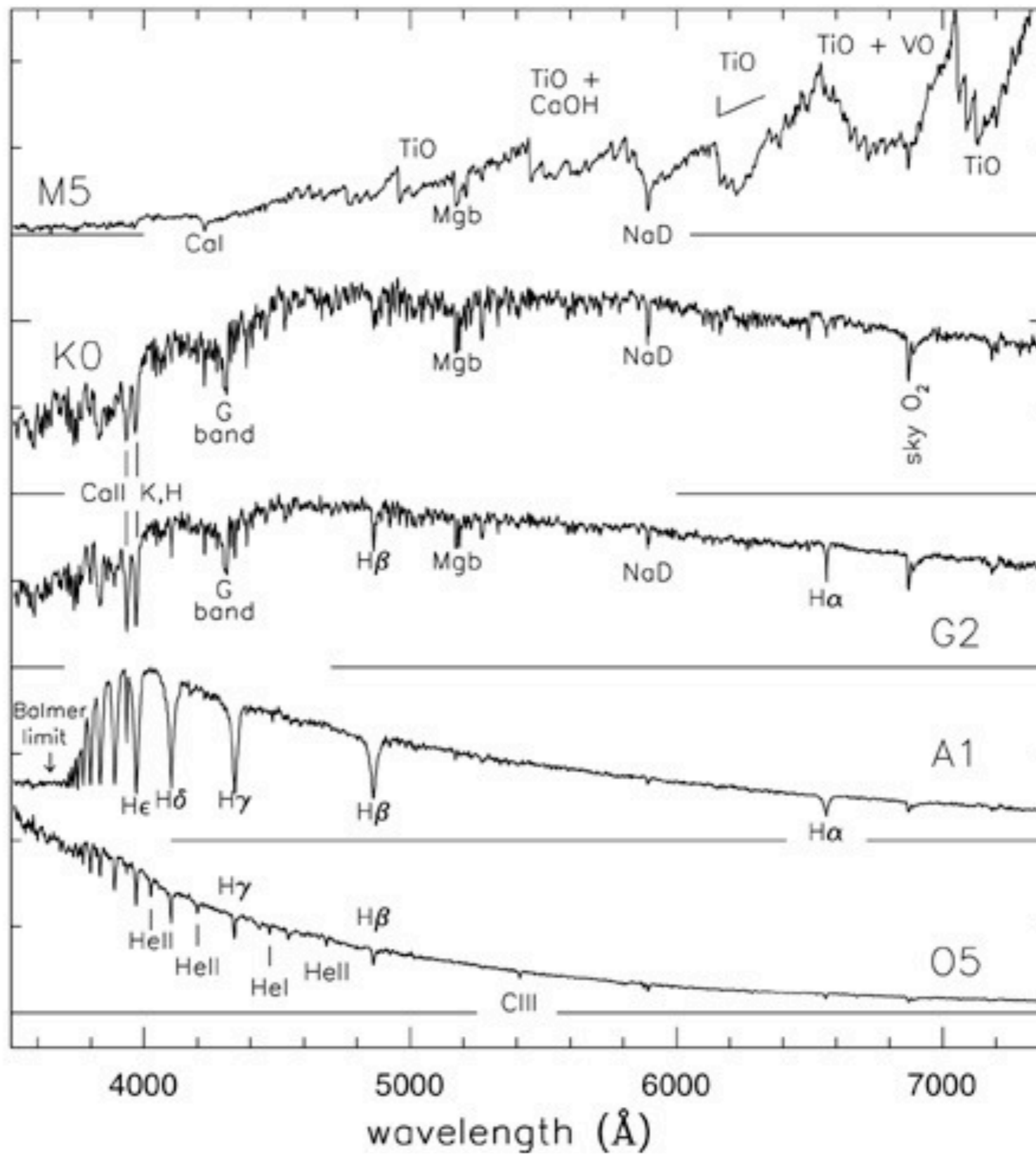
$$[\text{Fe}/\text{H}] \equiv \log_{10}[\epsilon(\text{Fe}/\text{H})_*] - \log_{10}[\epsilon(\text{Fe}/\text{H})_{\odot}]$$

where $\epsilon(\text{Fe}/\text{H})$ is the number of Fe atoms relative to the number of H atoms

- causes absorption lines where atoms “intercept” continuum light of star
- also “line blanketing”: when absorption lines “pile up”, like below 4000 Å (not to be confused with Balmer decrement...)



flux F_λ (arbitrary units)

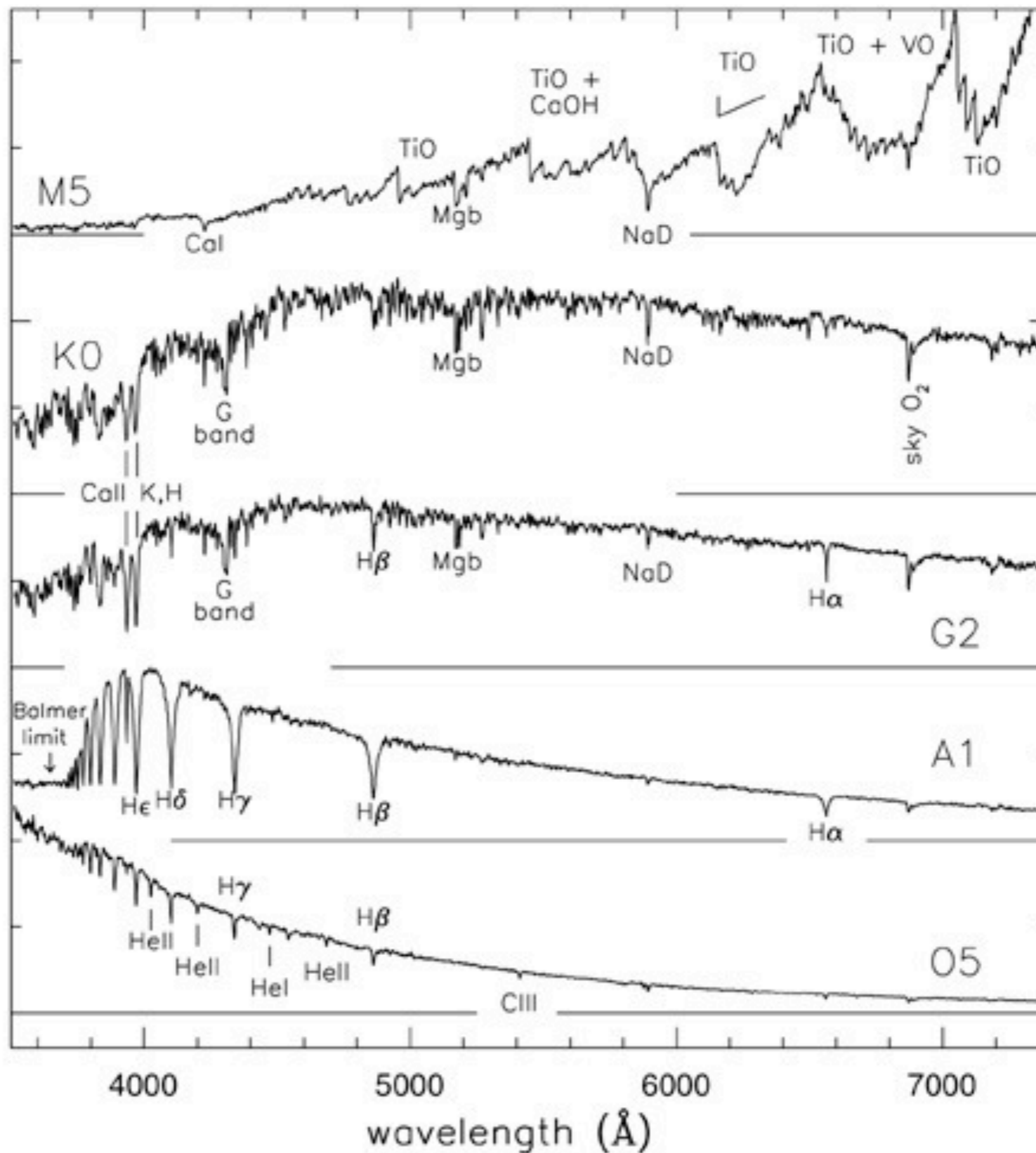


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flux F_λ (arbitrary units)



Cooler



Hotter



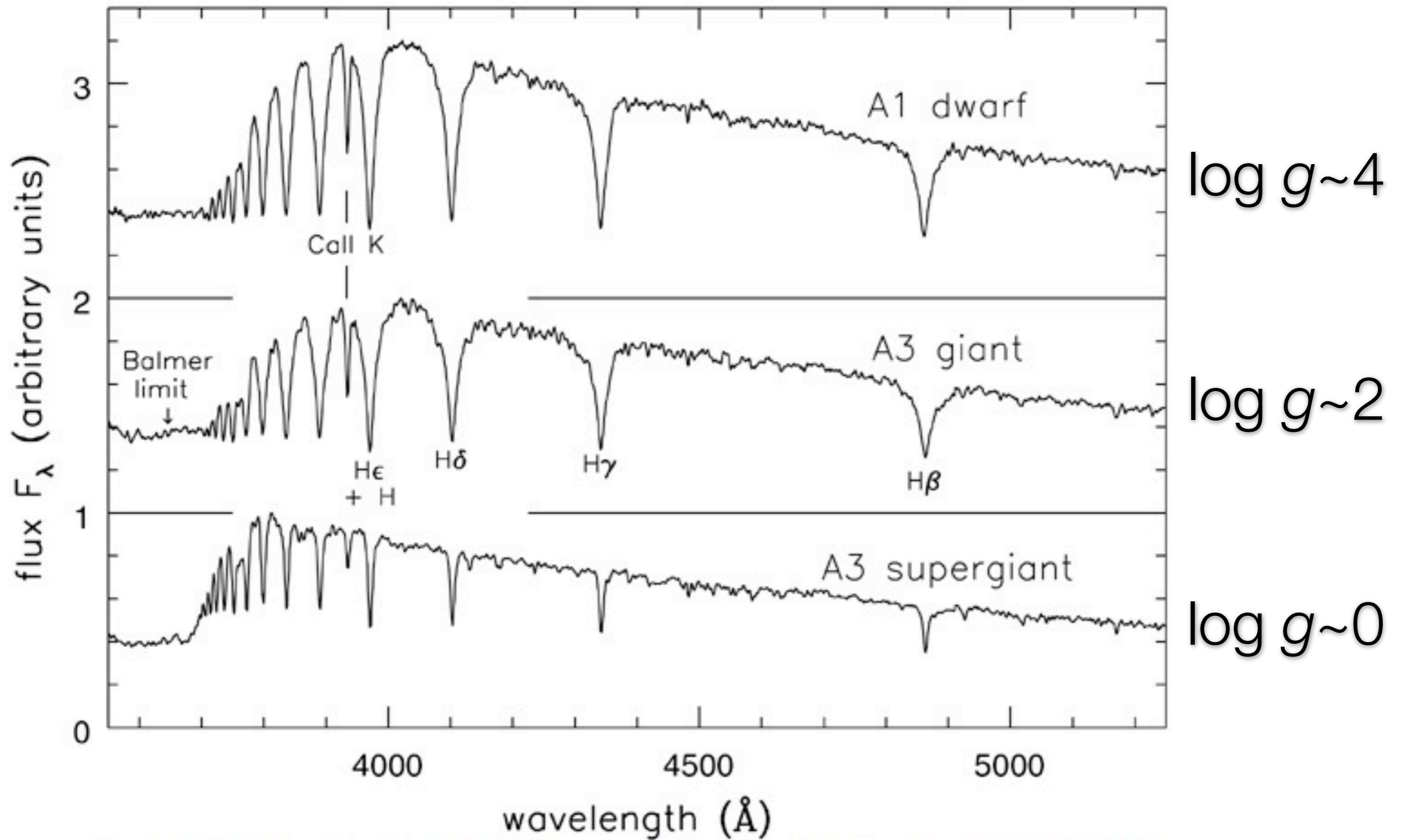
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- *Surface gravity:* $g \equiv \frac{GM}{R^2}$
- set gas pressure
- ionization balance (e.g., singly- vs. doubly-ionized Fe)
- pressure broadened lines (like H lines)
- abundances of certain molecules (in cool stars: CN, CO, etc.)





Stellar classification (quickly)

- Morgan-Keenan spectral classes
 - based on ratios of strong absorption lines
 - OBAFGKM(LT) $\Rightarrow T_{\text{eff}}$
 - Luminosity class I-V $\Rightarrow \log g$
 - Sun is G2V star
- Only works within $\sim 2x$ of solar abundance; qualitative; discrete steps; no abundance calibration



Magnitudes and colors

- Unfortunately, spectra are *expensive* in telescope time
 - need to spread photons out into lots of pixels – takes a lot of time to do!
- Need faster way to get stellar properties



- Consider two objects with fluxes f_1 and f_2
- Then the *magnitude difference* between these objects is

$$m_1 - m_2 = -k \log_{10} \left(\frac{f_1}{f_2} \right)$$
 - where $k = 2.5$ means that $m_1 - m_2 = 5$ when $f_2/f_1 = 100$
- so $m_1 - m_2 = -2.5 \log_{10} \left(\frac{f_1}{f_2} \right)$
- or $\frac{f_1}{f_2} = 10^{-0.4(m_1 - m_2)}$

note: from here on, I mean log-base-10 for "log" and log-base-e for "ln"



- This defines the **apparent magnitude**, the magnitude of the flux received by the detector at the telescope,

where
$$f \equiv \int_0^{\infty} f_{\nu} F_{\nu} R_{\nu} T_{\nu} d\nu$$

- here
 - f_{ν} is the flux of the object (in frequency units)
 - F_{ν} is the transmission of any filter used to isolate the (frequency) region of interest
 - R_{ν} is the transmission of the telescope, optics, and detector
 - T_{ν} is the transmission of the atmosphere (if any)



- Most (but not all) magnitude systems are based on taking a magnitude with respect to a star with a known (or predefined) magnitude
- the “Vega” system defines a set of A0V stars as having apparent magnitude 0 in *all* bands of a system (see below)
- So to get a “magnitude on system X”, one observes stars with known magnitudes and *calibrates* the “instrumental magnitudes” onto the “standard system”



- Some useful properties and “factoids” about magnitudes...
- The magnitude system is roughly based on *natural logarithms*: $m_1 - m_2 = 0.921 \ln(f_1 / f_2)$
- If $\Delta f \ll 1$, then $\Delta m = m_2 - m_1 \approx 1.086 \Delta f$
 - so the magnitude difference between two objects of nearly-equal brightness is equal to the fractional difference in their brightnesses – i.e., a difference of 0.1 magnitudes is $\sim 10\%$ in brightness
- A factor of 2 difference in brightness is a difference of 0.75 magnitudes



- We define the **color** of an object as the magnitude difference of the object in two different filters (“bandpasses”)

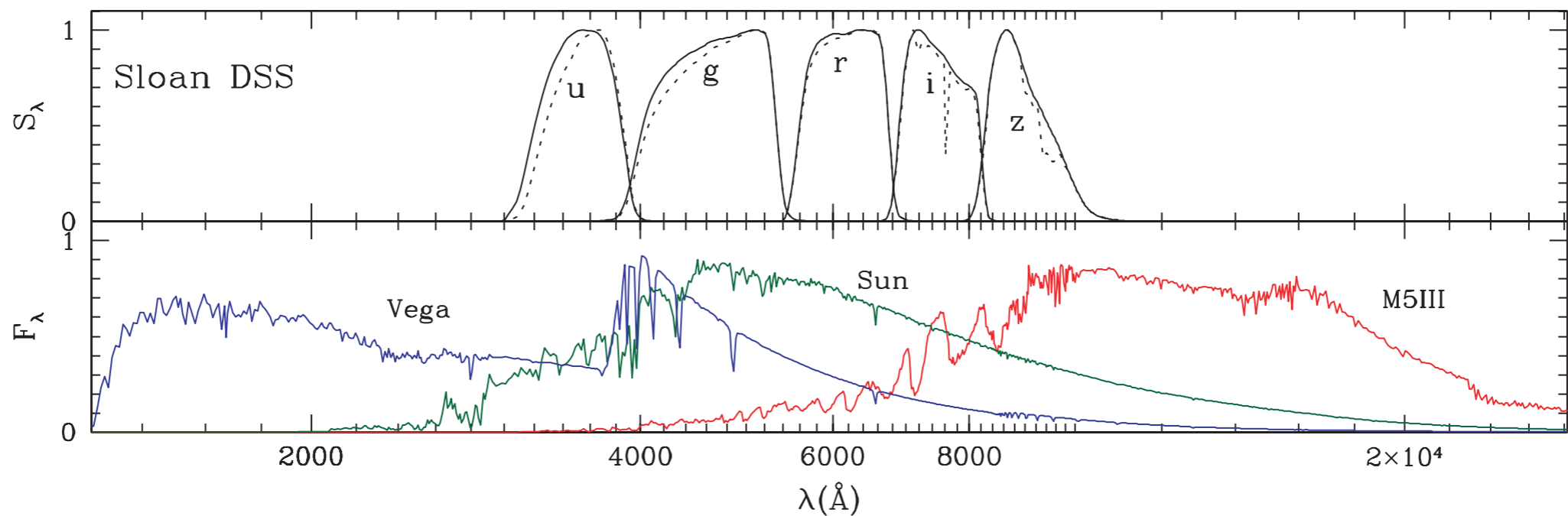
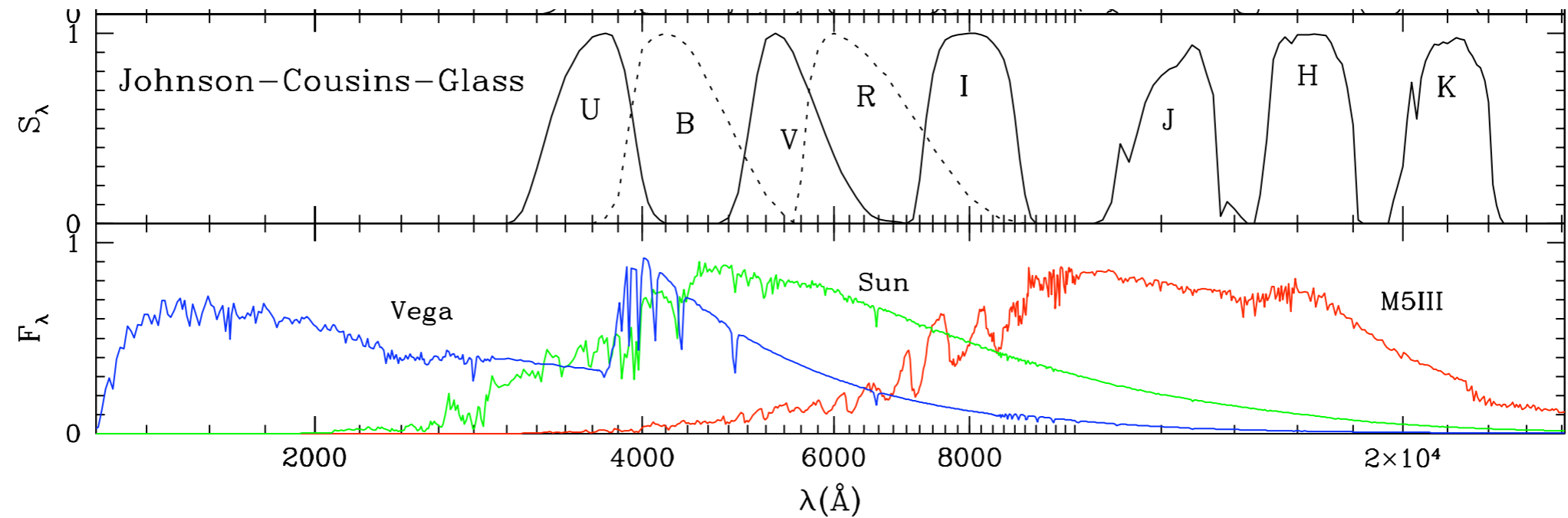
- if the filters are X and Y , then the color ($X-Y$) is

$$\begin{aligned}
 (X - Y) \equiv m_X - m_Y &= \text{const.} - 2.5 \log \frac{\int_0^\infty d\lambda S_\lambda(X) f_\lambda}{\int_0^\infty d\lambda S_\lambda(Y) f_\lambda} \\
 &= \text{const.} - 2.5 \log \frac{f_X}{f_Y}
 \end{aligned}$$

- where S_λ is the combined telescope-detector-filter sensitivity



Two common magnitude systems



- Notice that stars of different temperatures have very different spectra, so they have different fluxes through the filters
 - therefore they have *different colors*
 - and therefore well-chosen colors can be good **temperature indicators**
 - problems: calibration is tricky – requires good stellar radii; line blanketing makes bands bluer than $\sim V$ sensitive to $[\text{Fe}/\text{H}]$ and $\log g$



Apparent to absolute magnitudes

strictly true only if M is constant, but this is almost true on the giant branch.

- At fixed effective temperature, stars with lower $\log g$ must be bigger, since $g \equiv \frac{GM}{R^2}$
 - so they must have higher luminosity, since $L \equiv 4\pi R^2 \sigma T_{\text{eff}}^4$
- But we measure the flux, not the luminosity: $F = \frac{L}{4\pi D^2}$
- so we need the *distance D to the stars*



- We will come back to how to determine distances in the next lecture. In the meantime...
- We receive flux f at our telescopes from an object at distance d . If we wanted instead to know what the flux F is of that object if it were at distance D , then

$$f = \left(\frac{D}{d}\right)^2 F$$

- The difference in magnitudes m of f at d and M of F at D is then

$$m - M = -2.5 \log \left(\frac{f}{F}\right) = 5 \log \left(\frac{d}{D}\right)$$



- Now, we *always* pick D to be 10 parsecs (10 pc), so if d is measured in parsecs, then

$$m - M = 5 \log(d[\text{pc}]) - 5$$

- If d is measured in 10^6 pc=1 Mpc, then

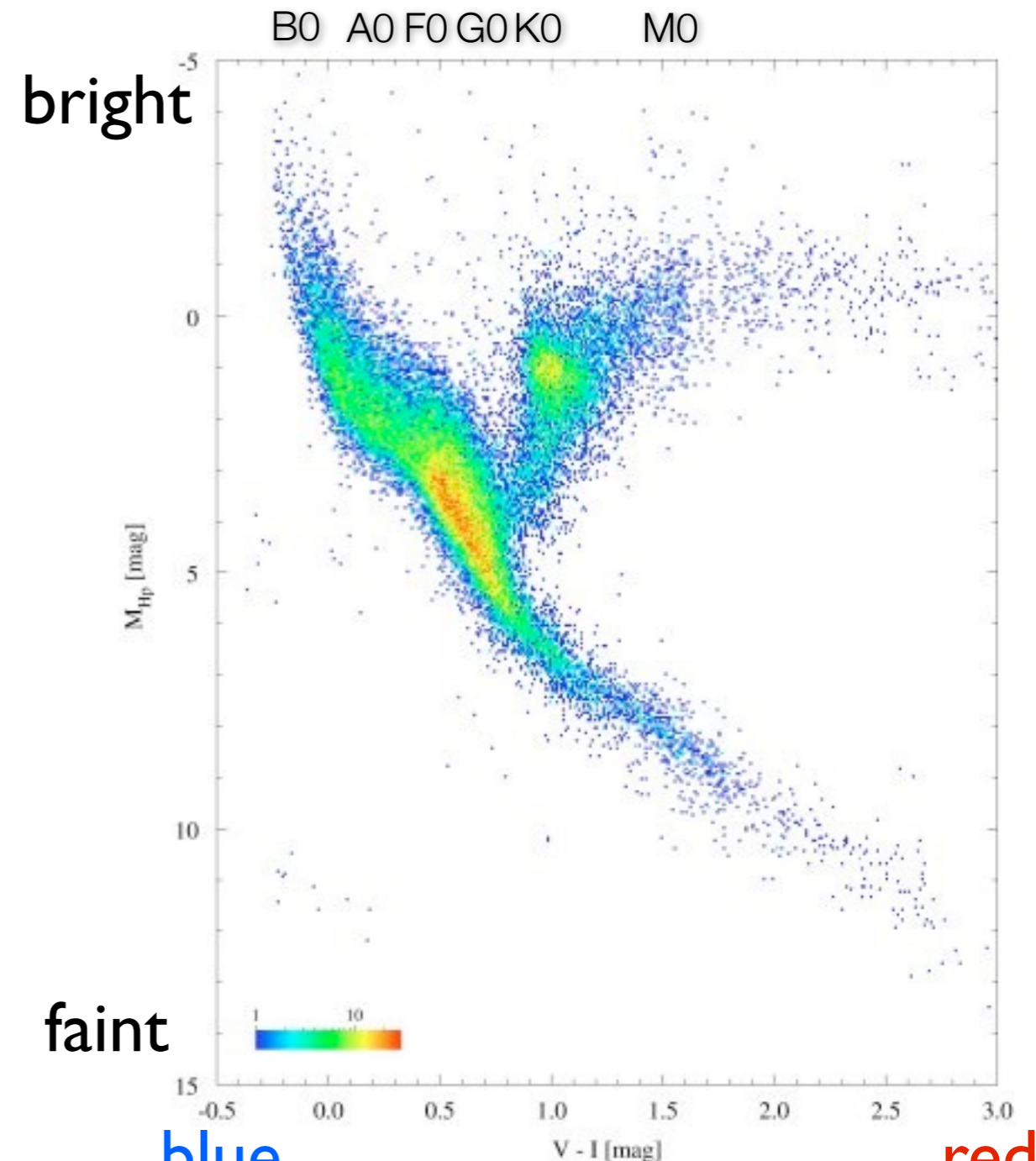
$$m - M = 5 \log(d[\text{Mpc}]) + 25$$

- $m-M$ is called the **distance modulus** and is sometimes written $\mu=m-M$
- ...and M is called the **absolute magnitude**



Color-magnitude diagrams

- By combining precisely-measured distances with apparent magnitudes and colors, we can plot a **color-magnitude diagram** (CMD) of stars
- This CMD is for stars within 100 pc of the Sun with accurate distances and colors



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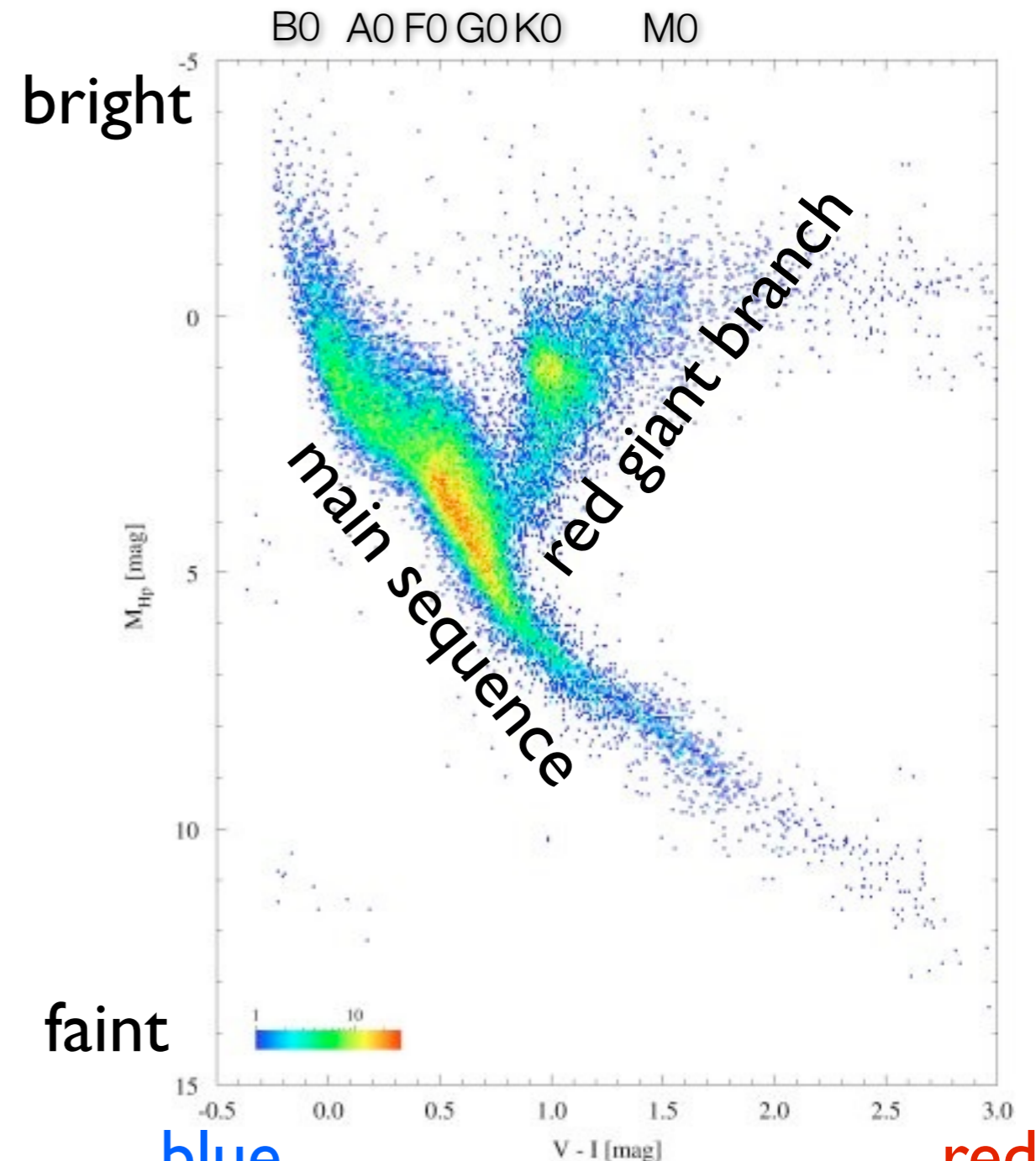
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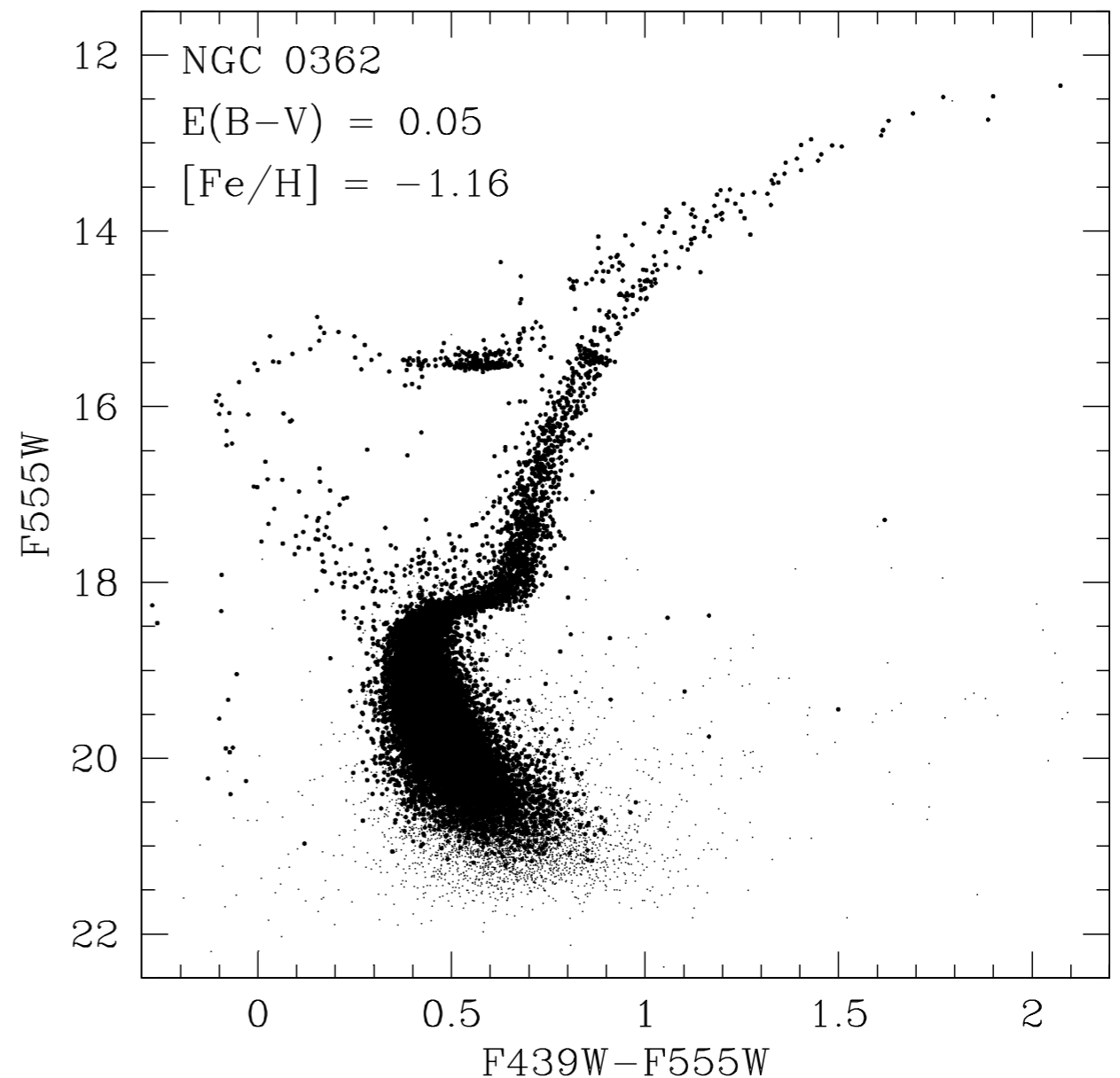
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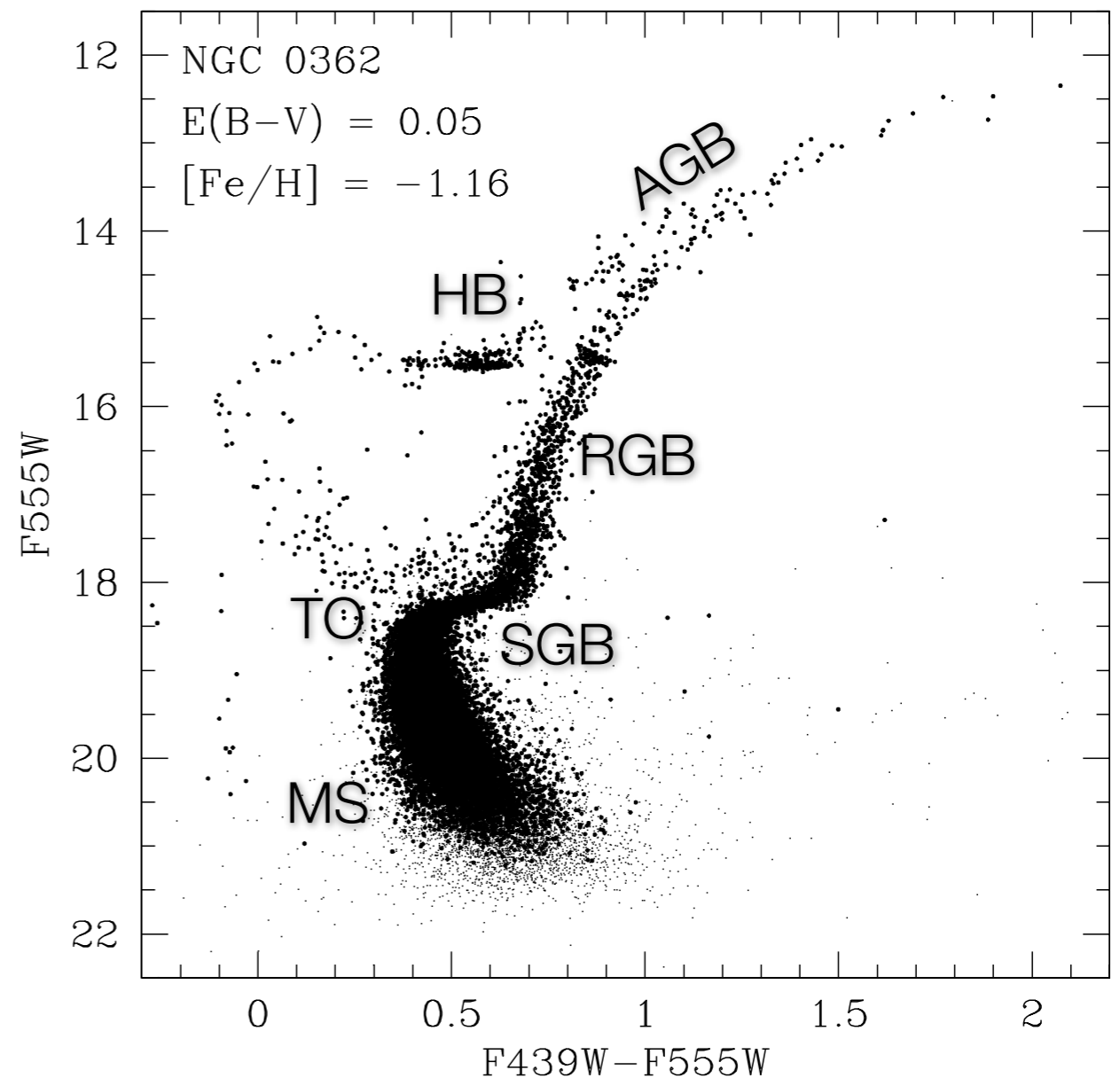
- ✦ Note!
 - ✦ A **Hertzsprung-Russell diagram** is a plot of luminosity as a function of effective temperature
 - ✦ A **color-magnitude diagram** is a plot of magnitude as a function of color
 - ✦ *They are not the same!*
 - ✦ The conversion from color to temperature depends on knowing the *radii* of stars, which (as mentioned) are difficult to measure!
 - ✦ Magnitude \neq luminosity unless you know the distance and the *bolometric correction*, which depends on stellar type (or color)



- ✦ The local CMD is made up of stars of many ages (we'll return to this point!)
- ✦ When we look at an object like a globular cluster, we see a very distinct pattern



- The local CMD is made up of stars of many ages (we'll return to this point!)
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Evolutionary timescales

- Stars burn hydrogen into helium for most of their lifetimes, until they exhaust the H in their cores
 - How long does this take?
 - Let's call E the amount of energy released by H burning during its *main-sequence lifetime* τ_{MS}
 - If the star's luminosity on the main sequence is L , then $E = L\tau_{\text{MS}}$
 - So if we can determine E , we can determine τ_{MS}



- Now, the amount of energy released by converting a mass dM of H into He is $dE=0.007c^2dM$
- So if a fraction α of the total mass of the star M is burned into He before core H-exhaustion, then $E=0.007c^2 \alpha M$
- so the main-sequence lifetime is $\tau_{\text{MS}} = 0.007\alpha Mc^2 / L$
- Typically, $\alpha=0.1$, so $\tau_{\text{MS}} = 10M/M_{\odot} (L/L_{\odot})^{-1}$ Gyr



- It is straightforward to work out from the equations of stellar structure that $L \propto M^4$ over a good chunk of the main sequence (for stars near the solar composition; this is also true observationally!)
- Then finally we have the main-sequence lifetime as a function of mass: $\tau_{\text{MS}} = 10A(M/M_{\odot})^{-3}$ Gyr
 - where A is a constant of ~ 1
- So a star ten times the mass of the sun lives for 1/1000 of the time: 10 Myr!



- The main-sequence lifetime is a crucial timescale: it is a *clock*
- By measuring the magnitude and color of the **main sequence turnoff** (MSTO or just TO) we can determine the age of a stellar population – and therefore something about its evolution
- The brighter and bluer the MSTO, the younger the population
 - brighter, bluer MS stars are more massive



Open or Galactic Clusters

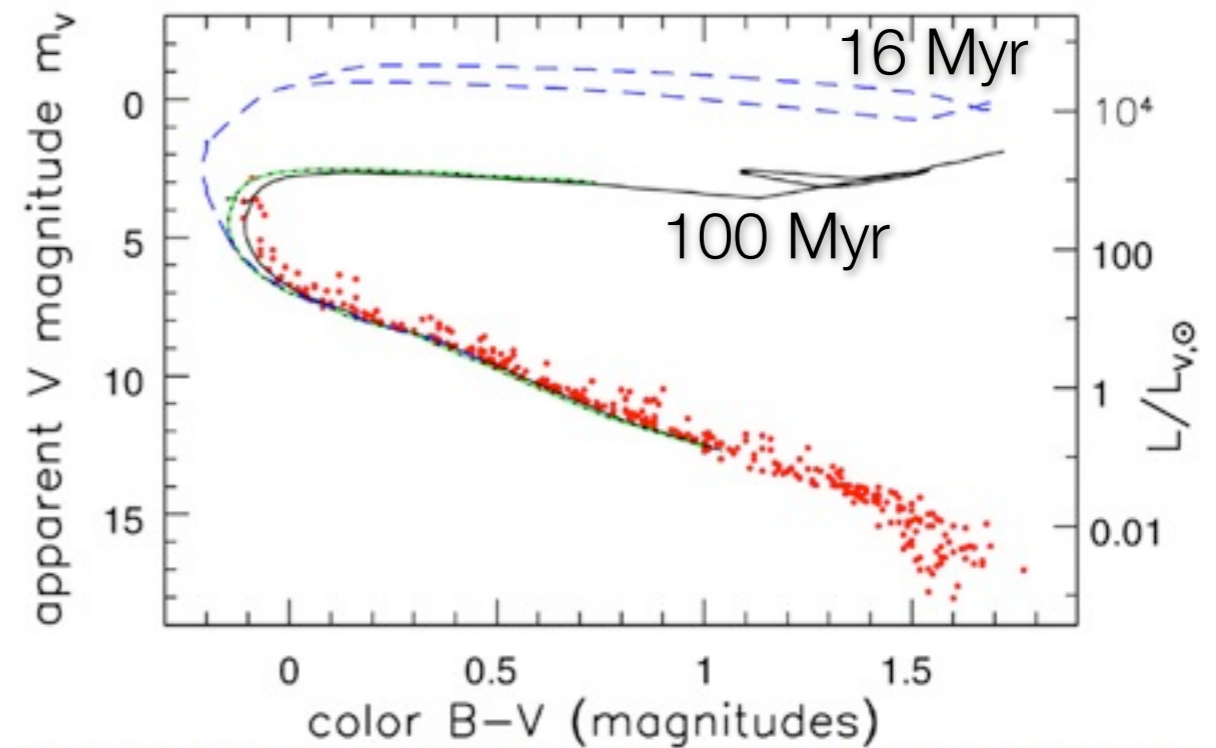
- “Open” or Galactic clusters are low mass, relatively small (~ 10 pc diameter) clusters of stars in the Galactic disk containing $< 10^3$ stars
- The Pleiades cluster is a good example of an open cluster
 - the “fuzziness” is starlight reflected from interstellar dust



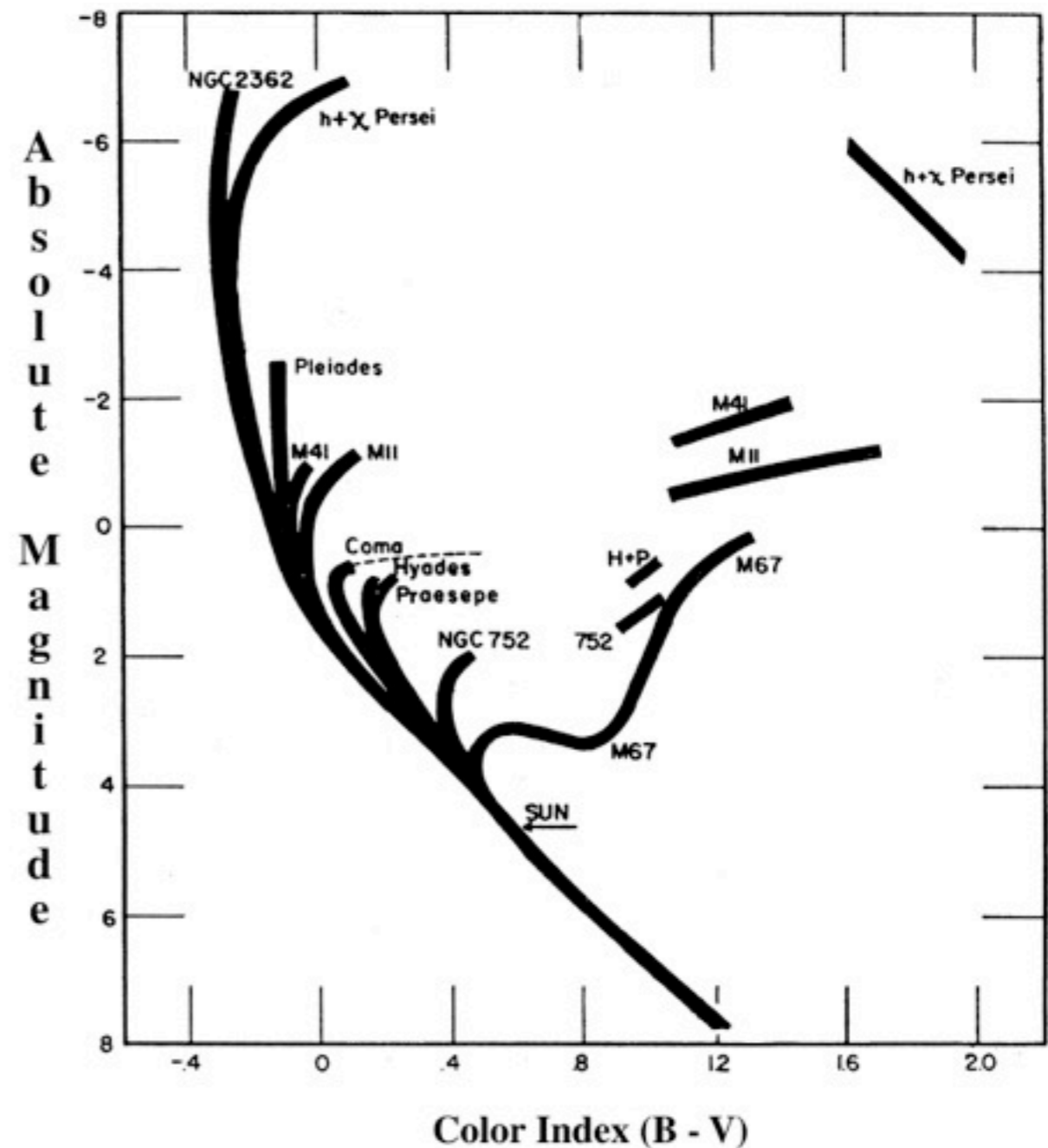
- Because open clusters live in the disk of our Milky Way, they are subject to strong tides and shearing motions
- Because they are so small and contain few stars, they also *evaporate* quickly
- Therefore they do not live very long unless they are very massive — so most of them are quite *young*



- ✦ All stars in an open cluster
 - ✦ are at the same distance
 - ✦ formed at the same time
 - ✦ have the same composition
- ✦ Very useful for testing stellar evolution models!



- They are also very useful for studying the evolution of the properties of the MW's disk
- Which of these clusters is the youngest?
- Which is the oldest?



Globular clusters

Globular clusters are named for their spherical shape and contain $\sim 10^4$ - 10^6 stars and are bigger than open clusters, with diameters of 20-100 pc

In the Milky Way, *all* globular clusters are **old**:
>10 Gyr!

M15 imaged with JKT at La Palma
core of M15 observed with HST

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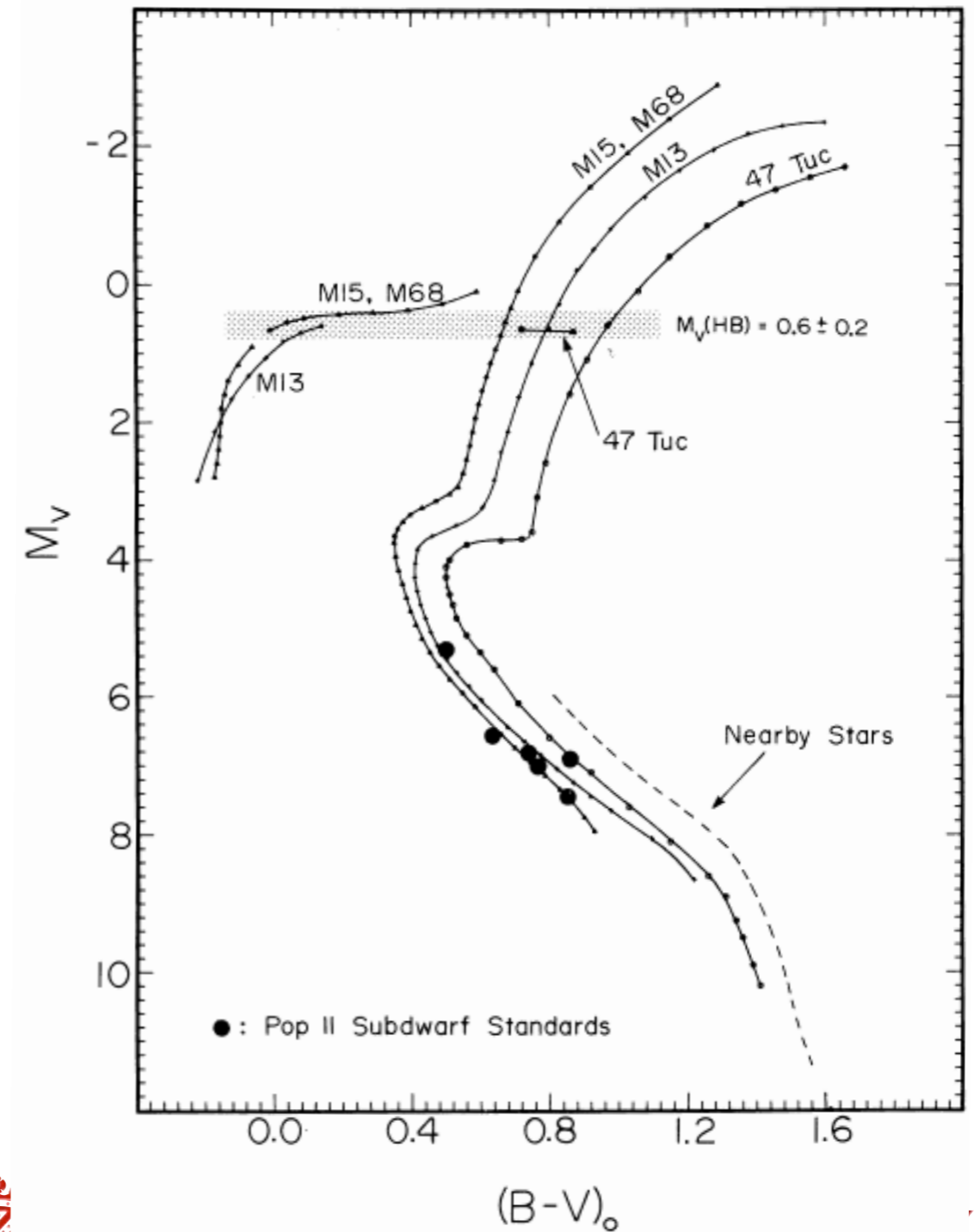
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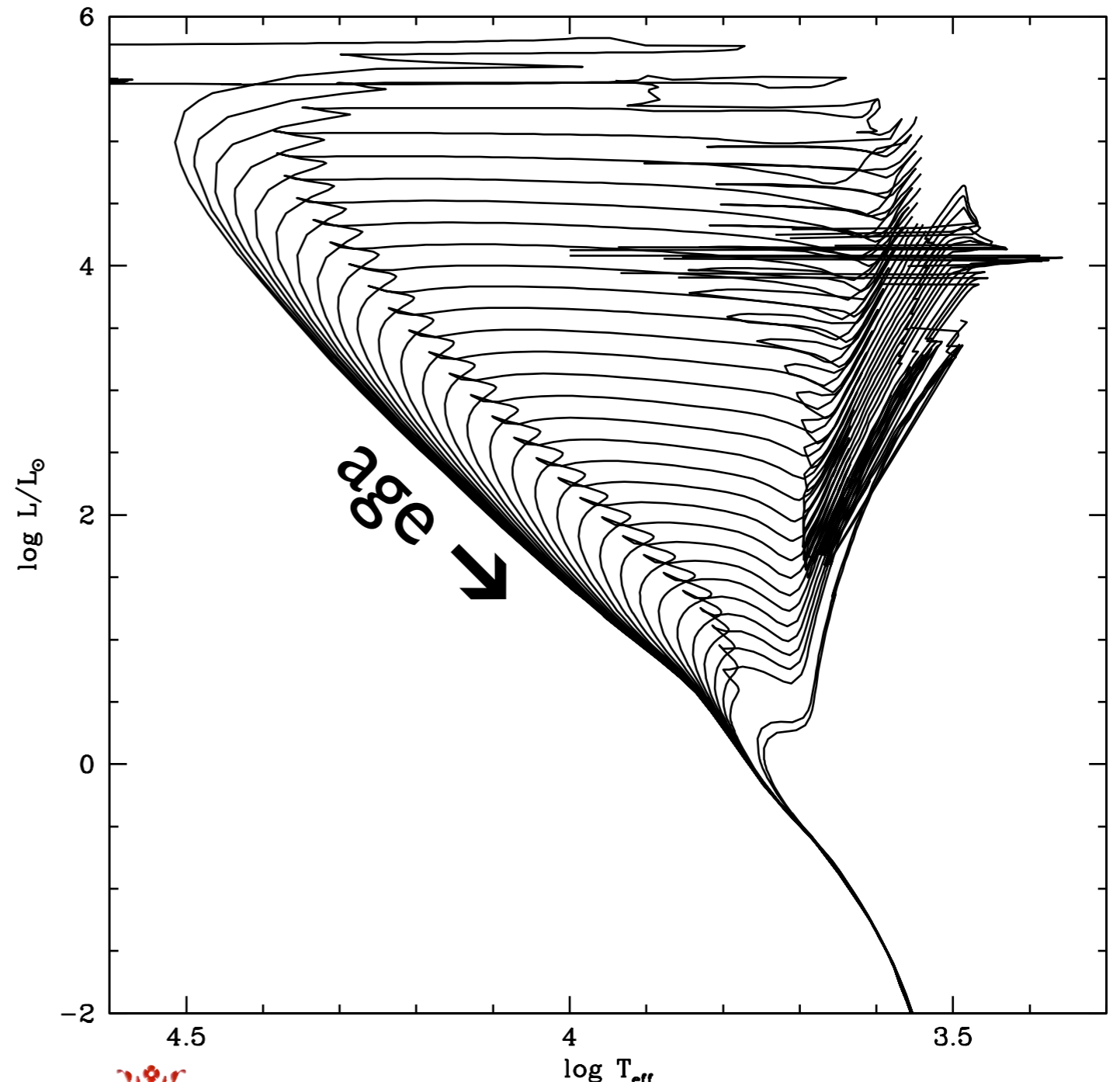
Globular cluster CMDs

- As *all* MW GCs are **old**, their CMDs are very similar. They vary primarily due to *composition differences*
 - “metallicity” [Fe/H]
 - He content
 - variations of other elements
- Age is an important but secondary consideration



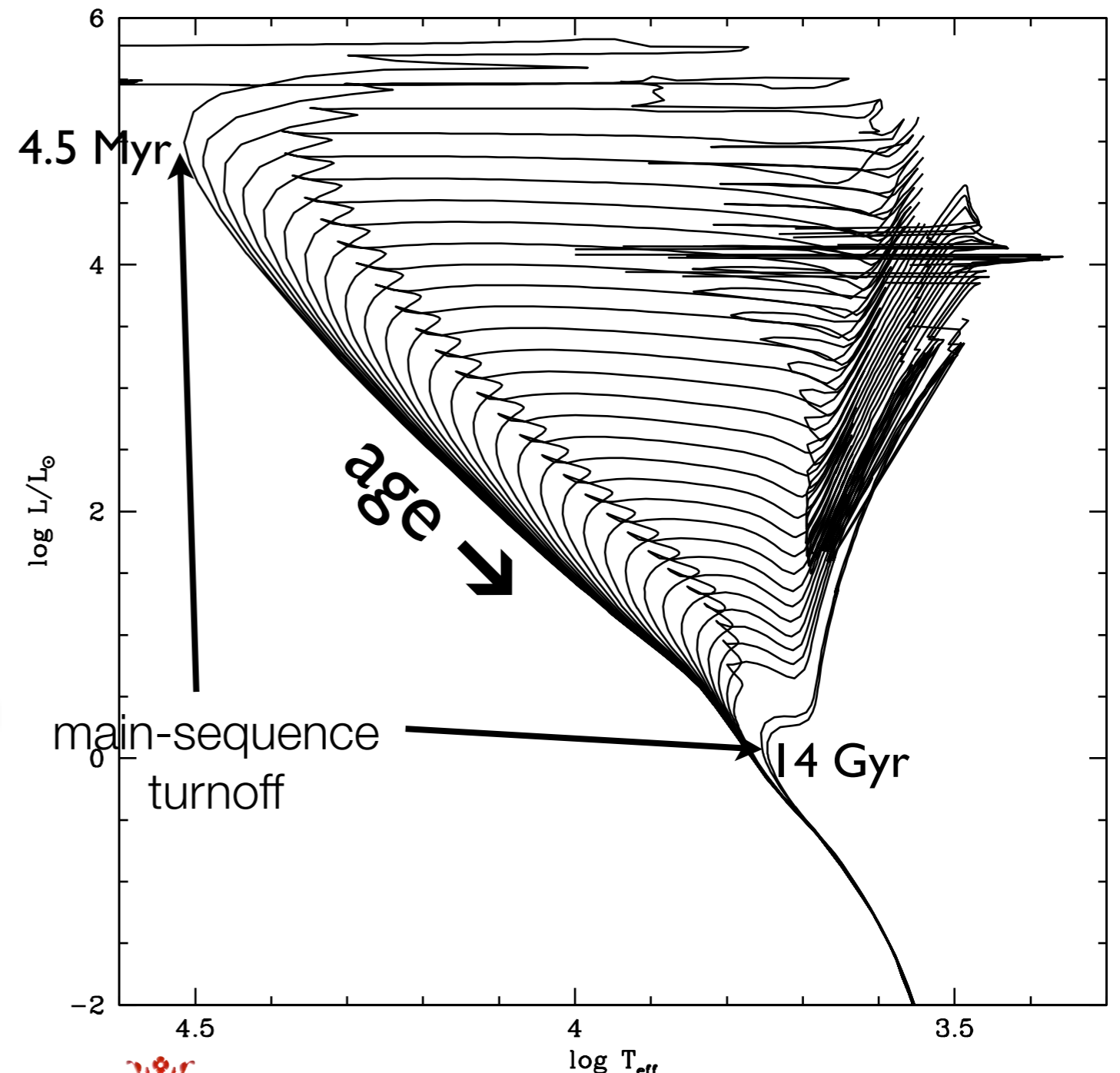
Isochrones: “single burst” stellar populations

- Take a set of stellar models all with the **same age** and **same composition** but **different masses**
- the resulting track in an HR diagram or a CMD is called an *isochrone*

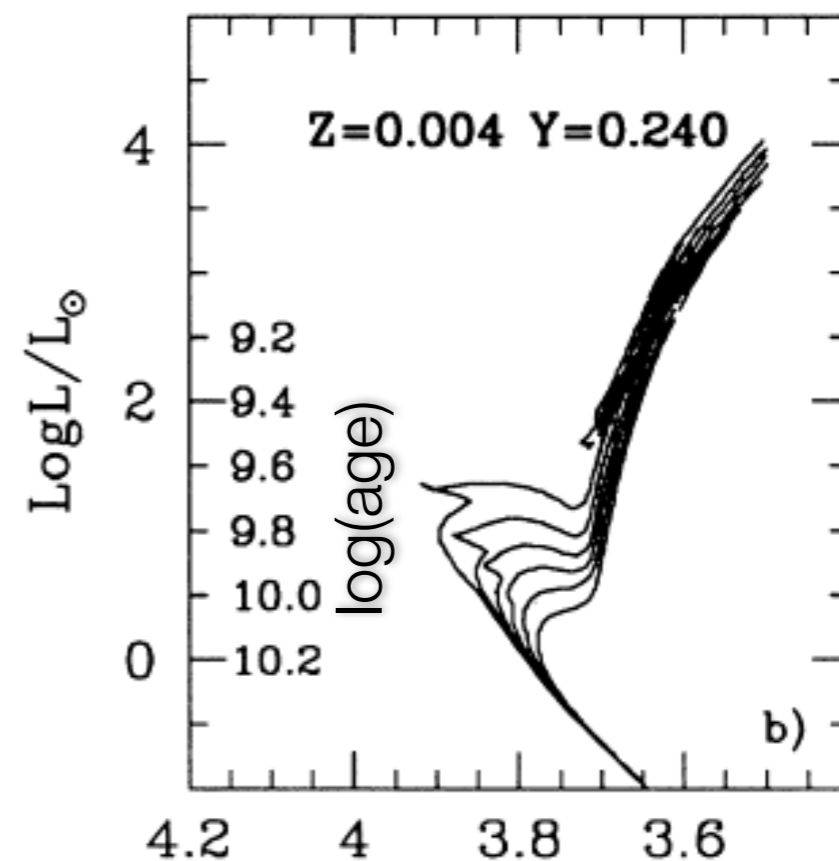
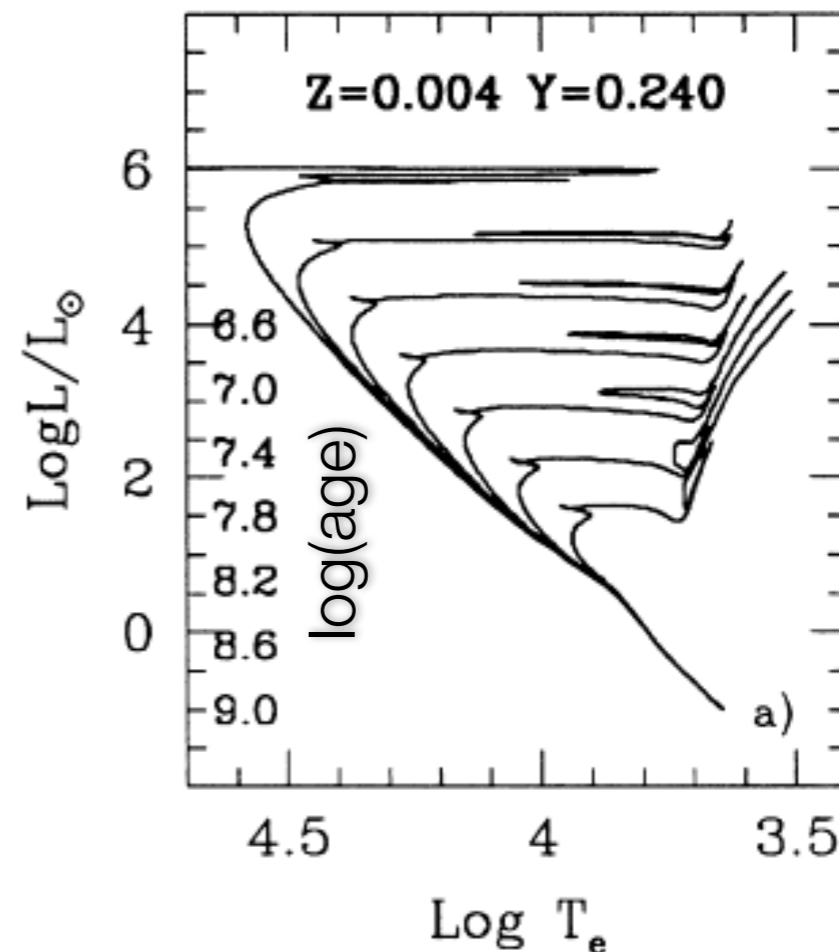


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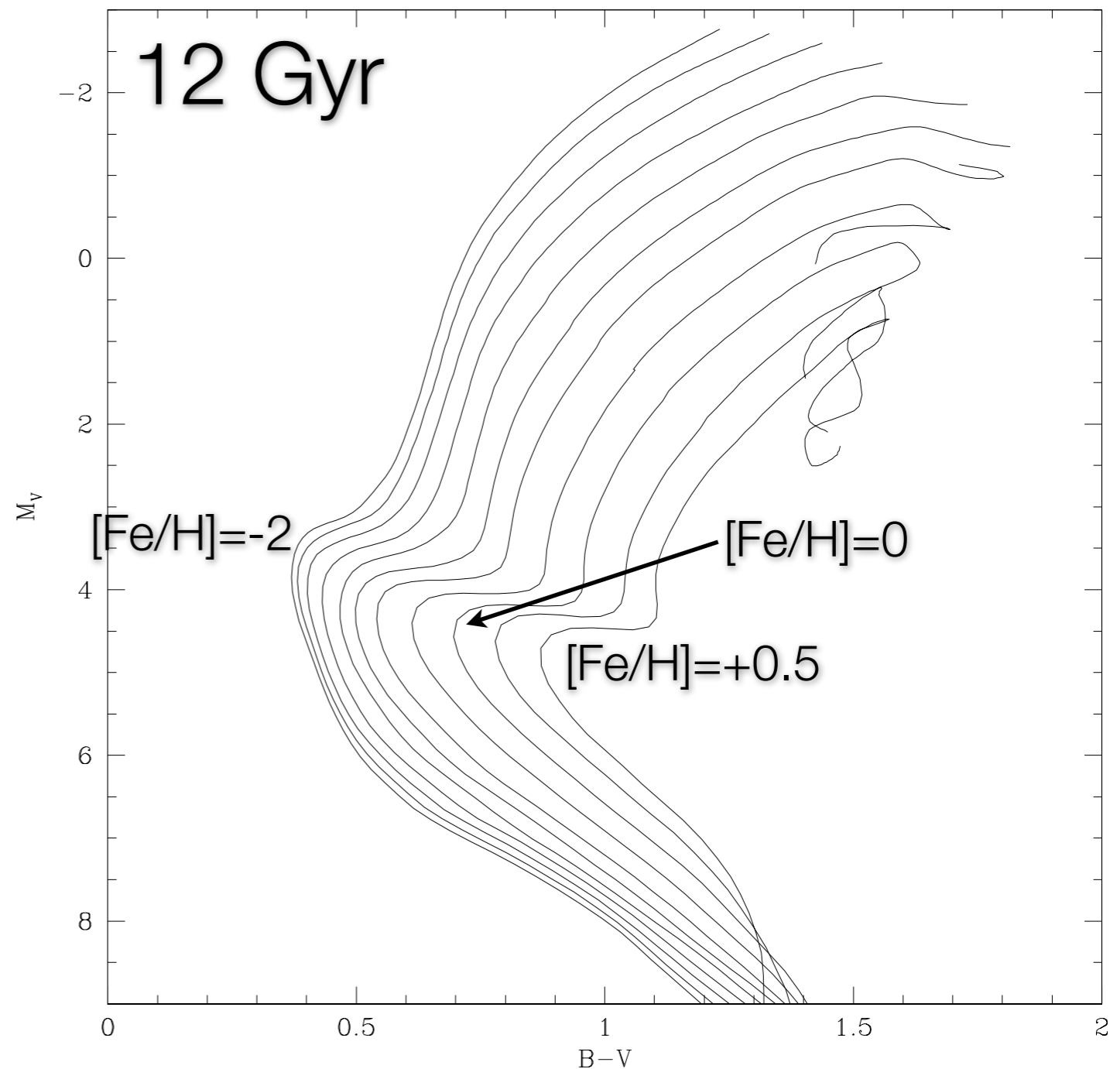


- ✦ When populations get old, the isochrones “pile up” at very similar luminosities and temperatures
- ✦ It is easy to determine ages for **young** populations
- ✦ ...but not for **old** populations!
 - ✦ “Age dating” GCs is difficult!



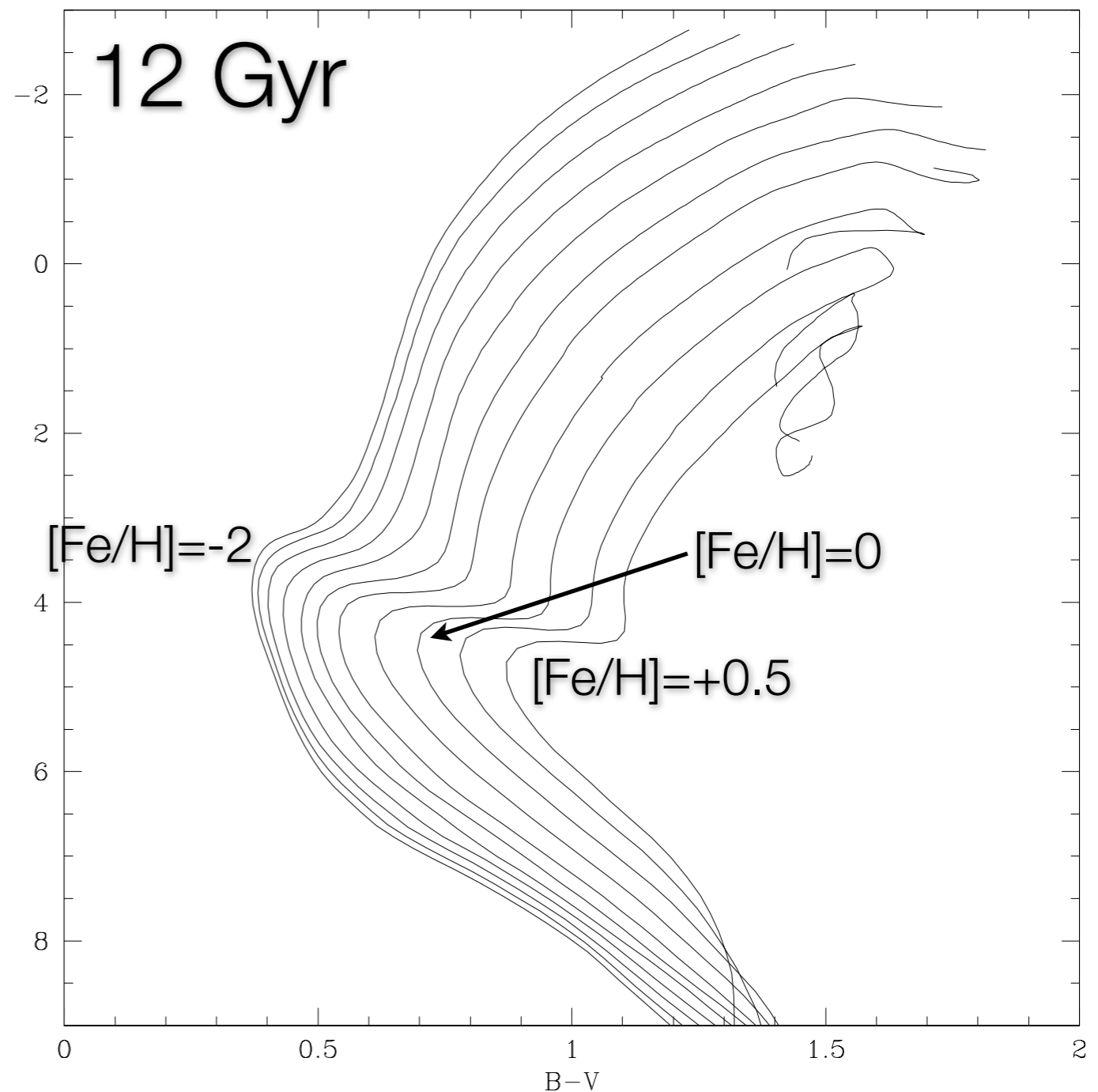
The effect of metallicity

- Stellar temperatures and luminosities also depend on composition at fixed age
- This is an *opacity* effect: more metals mean more absorption, especially in the blue (it's reradiated into the IR), so stars become cooler (redder) and dimmer



The effect of metallicity

- At fixed L , $\Delta[\text{Fe}/\text{H}]=-1.7$ shifts the unevolved MS (“zero-age MS” or ZAMS) to the blue by $\Delta T_{\text{eff}}=+0.06$
 - $\Delta(V-K)=-0.32$ mag
- At fixed color, $\Delta[\text{Fe}/\text{H}]=-1.7$ shifts the ZAMS fainter by ~ 1 magnitude
- low-metallicity ZAMS stars are thus called “subdwarfs”, because they’re fainter at the same mass



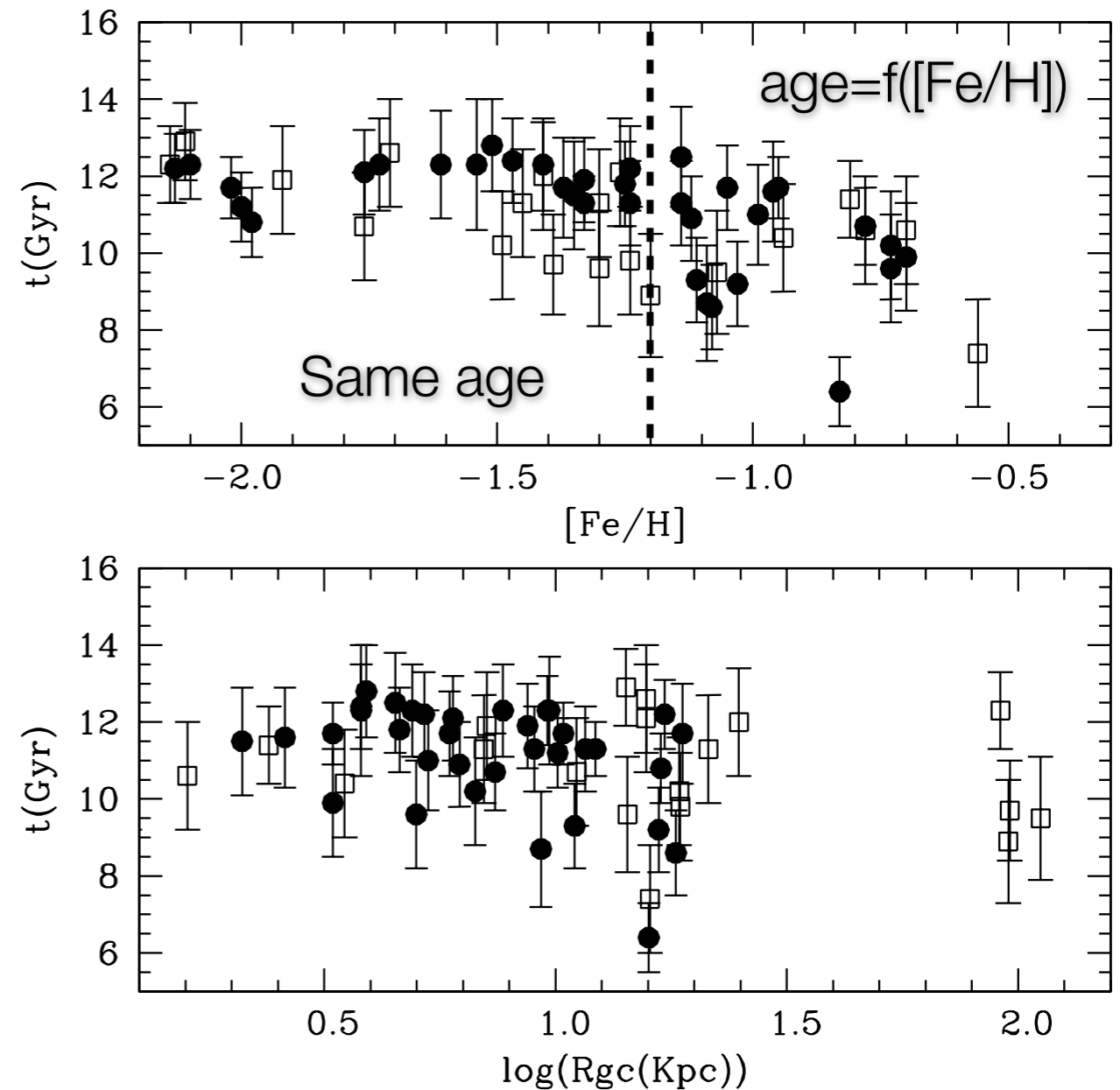
- ✦ Putting all of this together, the absolute V magnitude of the MSTO varies as

$$M_V(\text{MSTO}) = 2.7 \log(t/\text{Gyr}) + 0.3[\text{Fe}/\text{H}] + 1.4$$

- ✦ Uncertainties:
 - ✦ Difficult to measure the location of the MSTO, as the isochrones are \sim vertical
 - ✦ Distance errors cause big age errors:
 - ✦ 10% distance error \rightarrow 0.2 mag error in distance modulus \rightarrow 0.2 mag uncertainty in MSTO magnitude \rightarrow 20% error in age

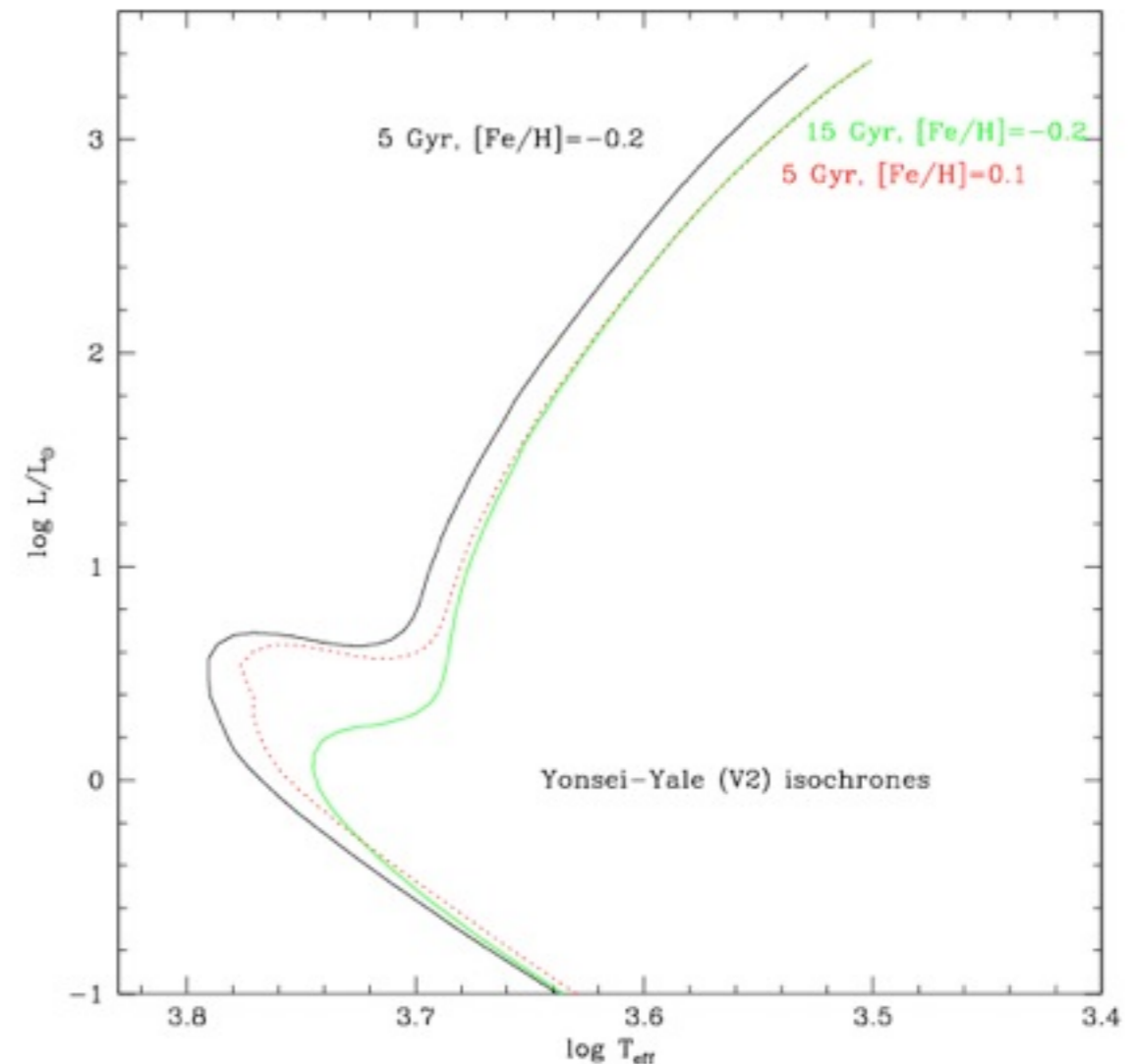


- Globular clusters in the Milky Way are generally very old (>10 Gyr), with only a few younger, metal-rich globular clusters



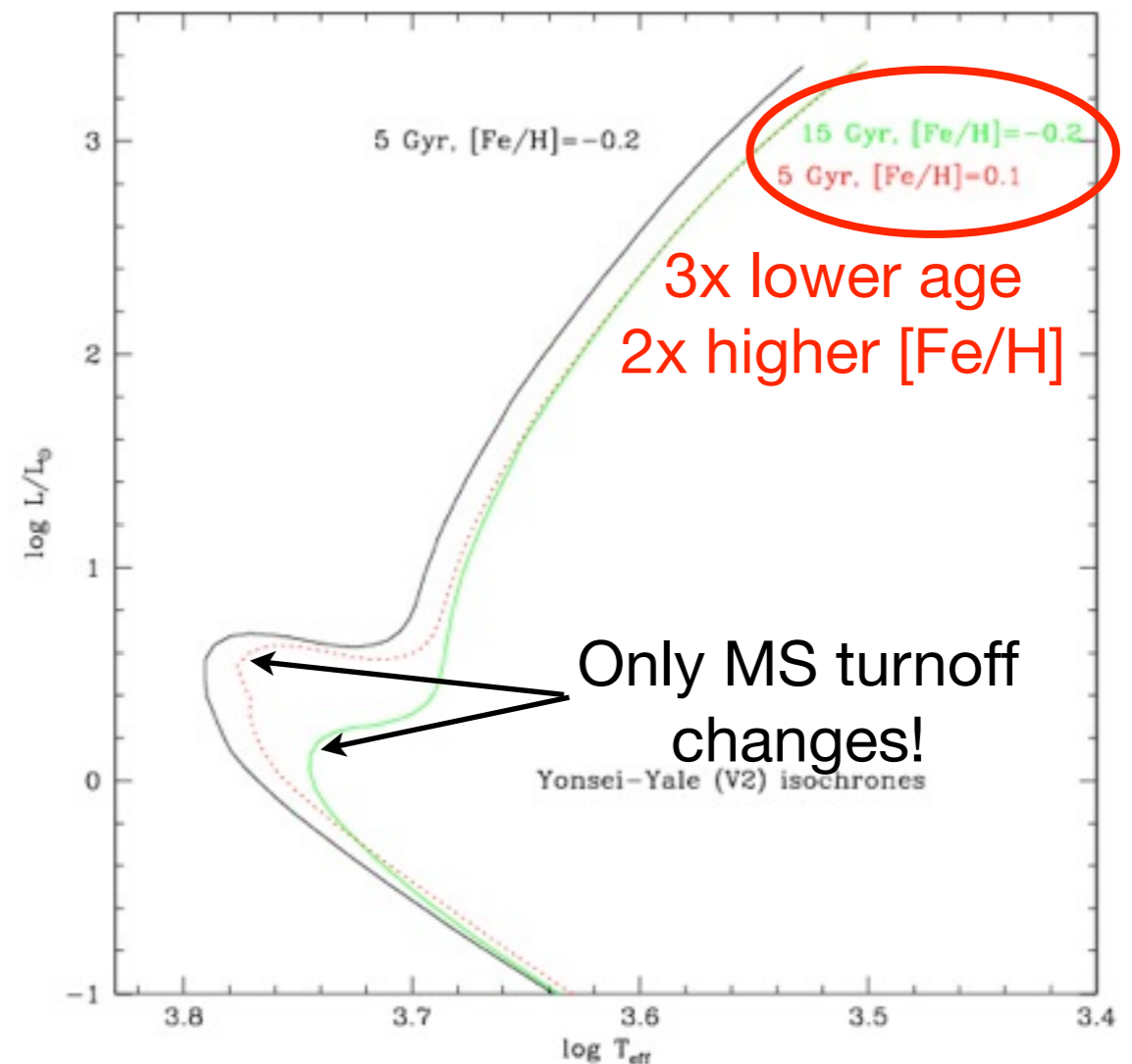
The age–metallicity degeneracy

- Because **both** age and metallicity affect isochrones, it can be *very* difficult to separate their effects on single-burst stellar populations
- Because the *colors* of an *unresolved* stellar population come from both the giant branch and the MS, colors are **not useful** to determine the ages of “old” populations



The age–metallicity degeneracy

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The luminosity function

- At a given point in the galaxy, the number of stars will vary, with both luminous and faint stars
- Consider a number dN of stars with absolute magnitudes in the range $(M, M+dM)$ in the volume element $d^3\mathbf{x}$ around the point \mathbf{x} :

$$dN = \Phi(M, \mathbf{x}) dM d^3\mathbf{x}$$



- Now let's assume that the *mix* of stars of different luminosities is independent of location (not always true...):
- Then we can separate $\Phi(M, \mathbf{x})$ into two functions, $\Phi(M)$ and $\nu(\mathbf{x})$: $dN = [\Phi(M) dM][\nu(\mathbf{x}) d^3\mathbf{x}]$
- We call $\Phi(M)$ the **luminosity function**, the relative fraction of stars of different luminosities, while $\nu(\mathbf{x})$ is the total number density of stars at \mathbf{x} .

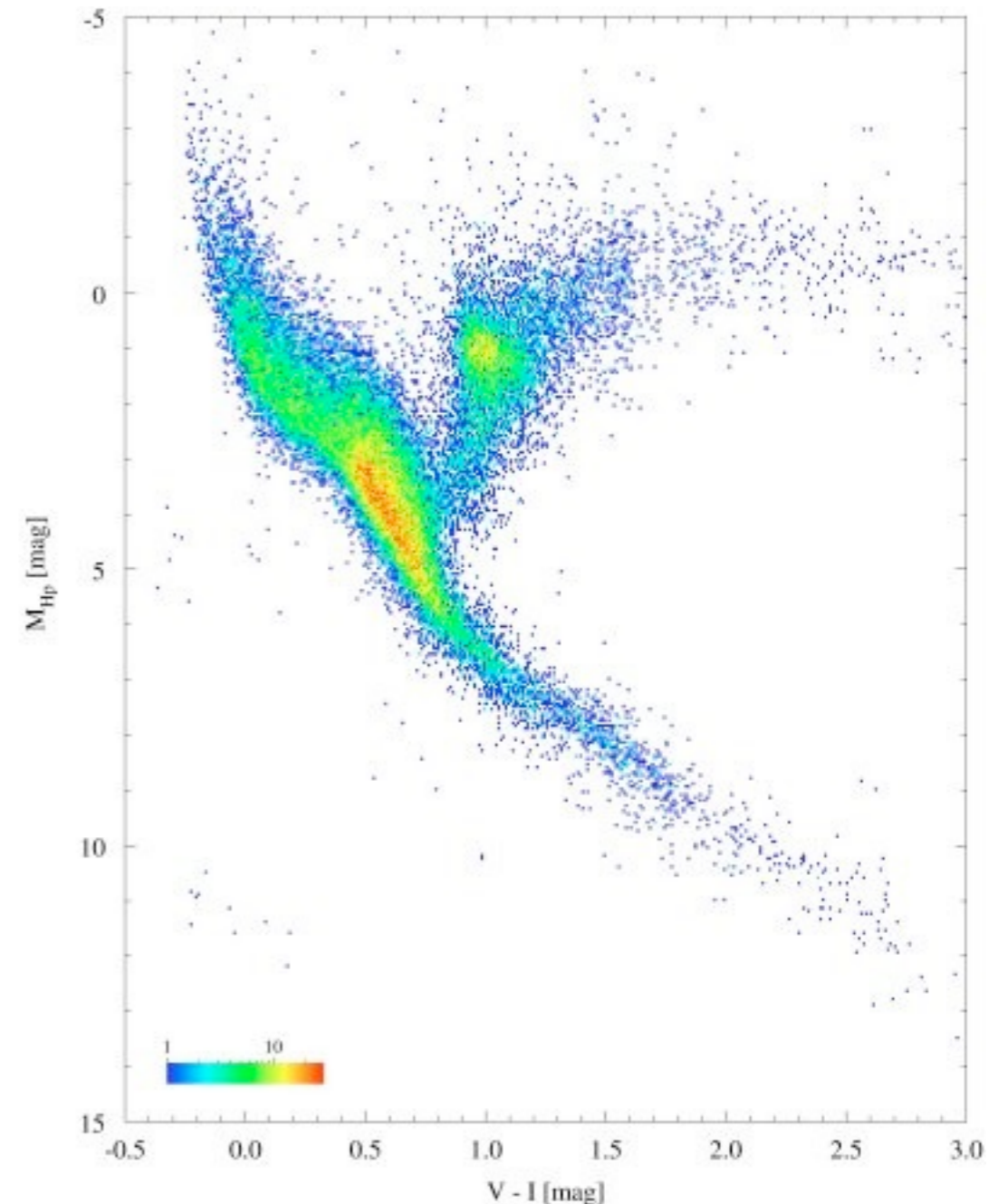


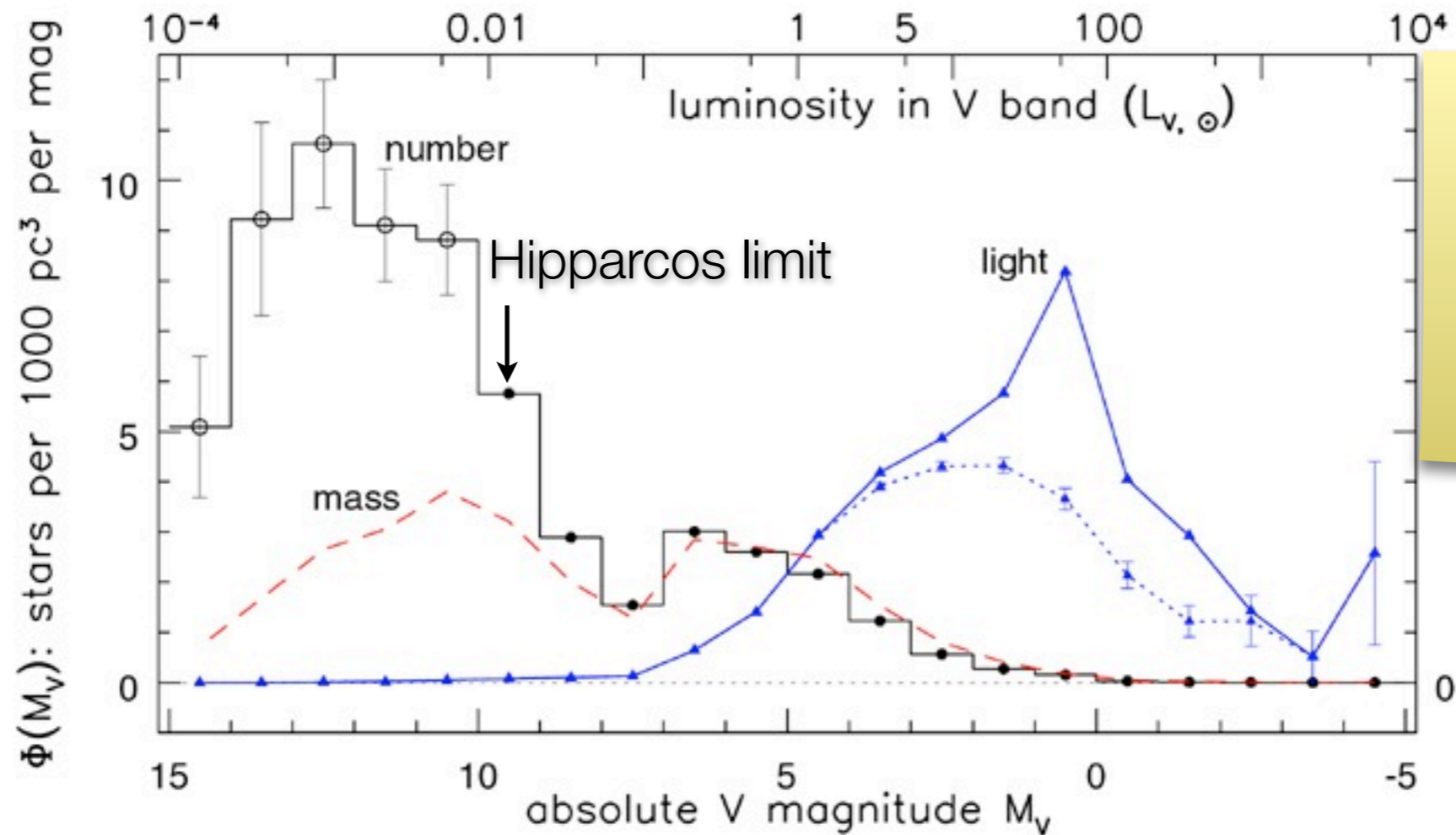
- Note that we will encounter and use the concept of luminosity functions even when we talk about galaxies
- In this case we generalize $\Phi(M)$ to mean the relative fraction of galaxies with different luminosities and $\nu(\mathbf{x})$ to be the number of density of galaxies



- Let's consider again the local stars observed by Hipparcos within 100 pc
- Hipparcos was *targeted* on stars with $m_V < 8$
- So we can construct an LF by writing

$$\Phi(M_V)dM = \frac{dN}{V_{\max}}$$
- where V_{\max} is the volume over which stars with M_V could be seen





- ✦ This is the result (in black). Note the following:
 - ✦ Most of the stars are **faint**: peak is at $M_V \sim 14$
 - ✦ If we weight by the *luminosity* of the stars, nearly all of the light is in **bright** stars: peak is at $M_V \sim 1$
 - ✦ If we weight by the *mass* of the stars, most of the stellar mass is in **low-mass** stars: over range $3 < M_V < 15$



- Another interesting result is the **mass-to-light ratio** of stars in the Solar Neighborhood:
 - the V-band luminosity density is $0.053 L_{\odot} \text{ pc}^{-3}$
 - the mass density is (including white dwarfs) $0.039 M_{\odot} \text{ pc}^{-3}$
 - combining, the mass-to-light ratio in solar units is

$$M/L_V \sim 0.67 M_{\odot}/L_{\odot}$$

- this is a *lower limit* to the total mass per unit luminosity (in a given band), because we haven't included the dark matter – which we must do in a different way...



- Note that what we've defined here is the **present-day luminosity function** (PDLF)
- this is the LF we see after the high-mass stars have evolved away



The initial mass function: IMF

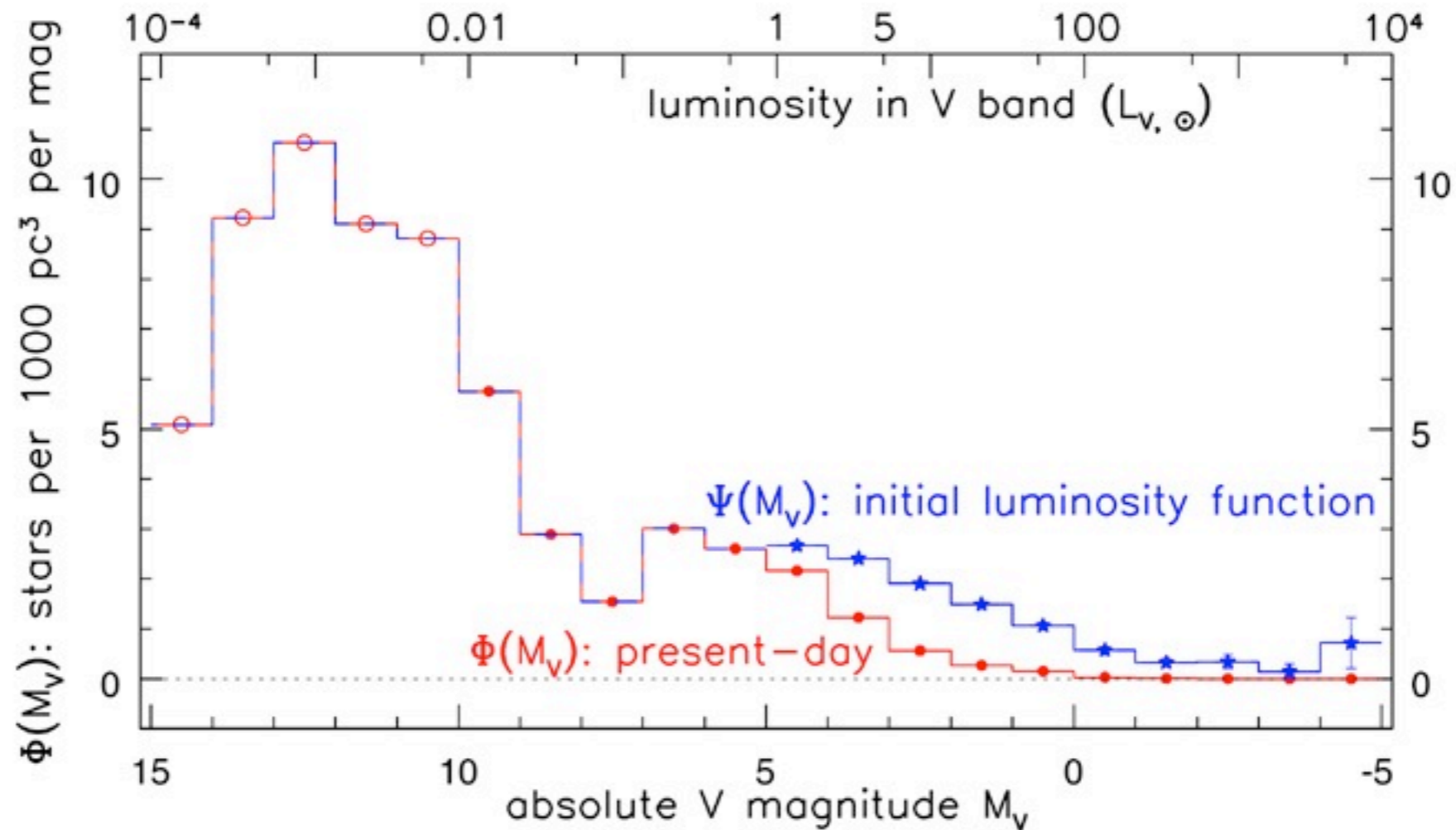
- A critical input for stellar population models crucial to understand the formation and evolution of galaxies is the **initial mass function**
- specifies the distribution of mass in stars *immediately after a star formation event*: the number of stars dN with masses between \mathcal{M} and $\mathcal{M}+d\mathcal{M}$ is

$$dN = N_0 \xi(\mathcal{M}) d\mathcal{M}$$

- We normalize $\xi(\mathcal{M})$ such that N_0 is the total number of solar masses formed in the event:

$$\int \mathcal{M} \xi(\mathcal{M}) d\mathcal{M} = M_{\odot}$$





- To determine $\xi(\mathcal{M})$, we need to determine dN immediately after a star-formation event to get the *initial luminosity function* Φ_0 : remember that $dN = \Phi(M) dM$
 - if all the stars *just* formed in one burst: no correction



- If that's not the case, but the system has been forming stars at a *constant* rate (roughly true for the MW disk), then

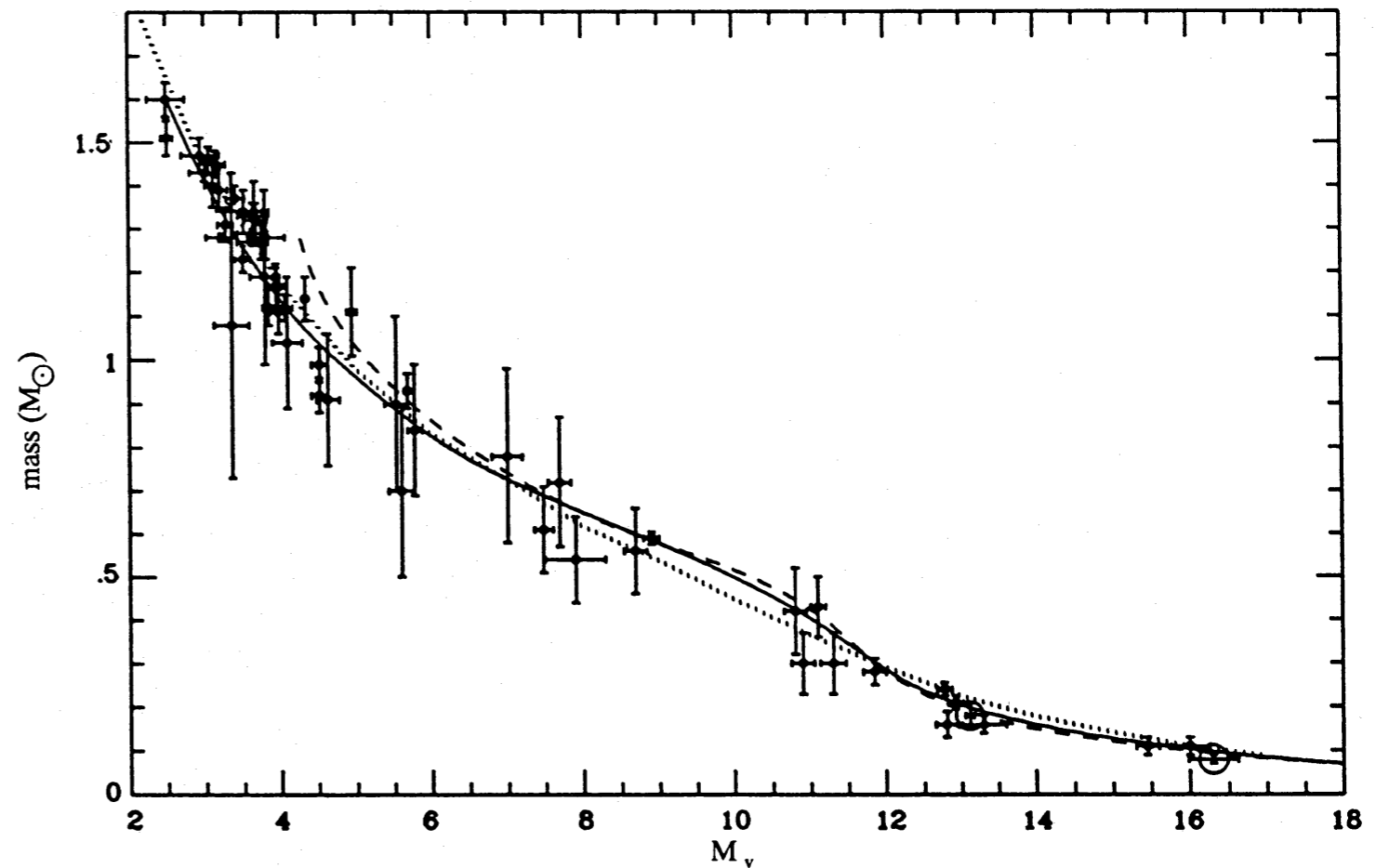
$$\Phi_0(M) = \Phi(M) \times \begin{cases} t/\tau_{\text{MS}} & \text{for } \tau_{\text{MS}}(M) < t \\ 1 & \text{otherwise} \end{cases}$$

- The factor t/τ_{MS} corrects for the fact that we only see stars of magnitude M that formed in the last fraction τ_{MS}/t of the population's life
- Now we can determine $\xi(\mathcal{M})$:

$$\xi(\mathcal{M}) = \frac{dM}{d\mathcal{M}} \Phi_0[M(\mathcal{M})]$$



- The function $M(\mathcal{M})$ specifies the relation between stellar mass and absolute magnitude
- unfortunately, you also need the derivative of this function...



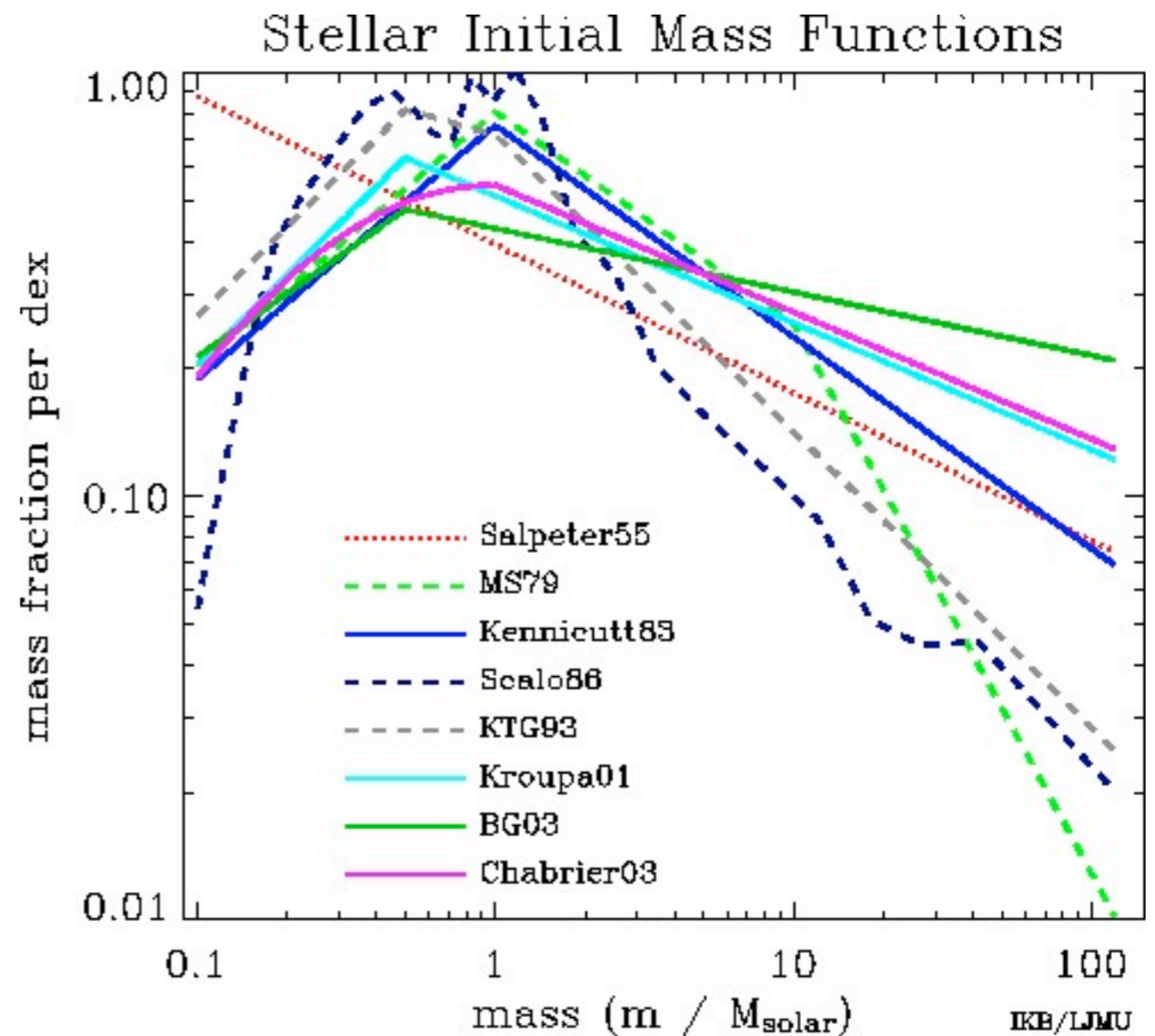
- The simplest (and still commonly used) IMF is that inferred by Salpeter (1955), which is a *power-law*:

$$\xi(\mathcal{M}) \propto \mathcal{M}^{-2.35}$$

- note that you often see the exponent as -1.35 , but this is a different definition of the IMF, in $\log \mathcal{M}$ instead of (linear) \mathcal{M}
- In reality, we know that the low-mass IMF must be *flat* or even *decline* with decreasing mass, at least in the Solar Neighborhood



- This has led to a number of different parameterizations of the IMF...
- Salpeter, Kroupa, and Chabrier are the most commonly-used IMFs today



- ✦ The big question: **is the IMF universal?**
- ✦ Nearly everyone assumes so... but this may not be true....



The evolution of stellar populations

- Stars with masses $<1.5 M_{\odot}$ live for >2.5 Gyr (at solar abundance) and put out most of their energy **after** the MSTO
- for instance, a $1 M_{\odot}$ star puts out $E_{MS} \sim 10.8 L_{\odot}$ Gyr on the MS and $E_{GB} \sim 24 L_{\odot}$ Gyr in the post-MS (RGB, HB, AGB) phases



- If stars of mass M emit a total energy of E_{GB} on the giant branch, then the luminosity of the population will be

$$L \approx \left(E_{\text{GB}} \frac{dN}{d\mathcal{M}} \right)_{\mathcal{M}_{\text{GB}}} \left| \frac{d\mathcal{M}_{\text{GB}}}{dt} \right|$$

this is the energy released on the giant branch per star times the change in the number of stars per unit time

- For these low-mass stars, the MS lifetimes are

$$\frac{\tau_{\text{MS}}}{10 \text{ Gyr}} \approx \left(\frac{\mathcal{M}}{M_{\odot}} \right)^{-3}$$

Note that GB lifetimes are ~10% of MS lifetimes, so the range of masses on the GB is small --- roughly that of the MSTO plus a very small amount

- or

$$\frac{\mathcal{M}_{\text{GB}}}{M_{\odot}} \approx \left(\frac{\tau_{\text{MS}}}{10 \text{ Gyr}} \right)^{-1/3}$$

- and so

$$\frac{d\mathcal{M}}{dt} \approx -\frac{1}{3} \left(\frac{\mathcal{M}_{\text{GB}}}{M_{\odot}} \right)^4 \left(\frac{M_{\odot}}{10 \text{ Gyr}} \right)$$



- Now, if we take the IMF to be a power-law with $\alpha \leq -2.35$ around \mathcal{M}_{GB} , then $dN/d\mathcal{M} \approx K(\mathcal{M}/M_{\odot})^{-\alpha}$

- And the luminosity is then

$$L \approx \frac{K M_{\odot} E_{\text{GB}}(\mathcal{M}_{\text{GB}})}{30 \text{ Gyr}} \left(\frac{\mathcal{M}_{\text{GB}}}{M_{\odot}} \right)^{4-\alpha}$$

- Finally, we can differentiate L to find

$$\frac{d \ln L}{dt} \approx \left[\frac{d \ln E_{\text{GB}}}{d \ln \mathcal{M}_{\text{GB}}} + (4 - \alpha) \right] \frac{d \ln \mathcal{M}_{\text{GB}}}{dt}$$

$$\approx 0.3\alpha - \left[1.3 + 0.3 \frac{d \ln E_{\text{GB}}}{d \ln \mathcal{M}_{\text{GB}}} \right]$$

take the log of the 2nd eq. on the previous page...

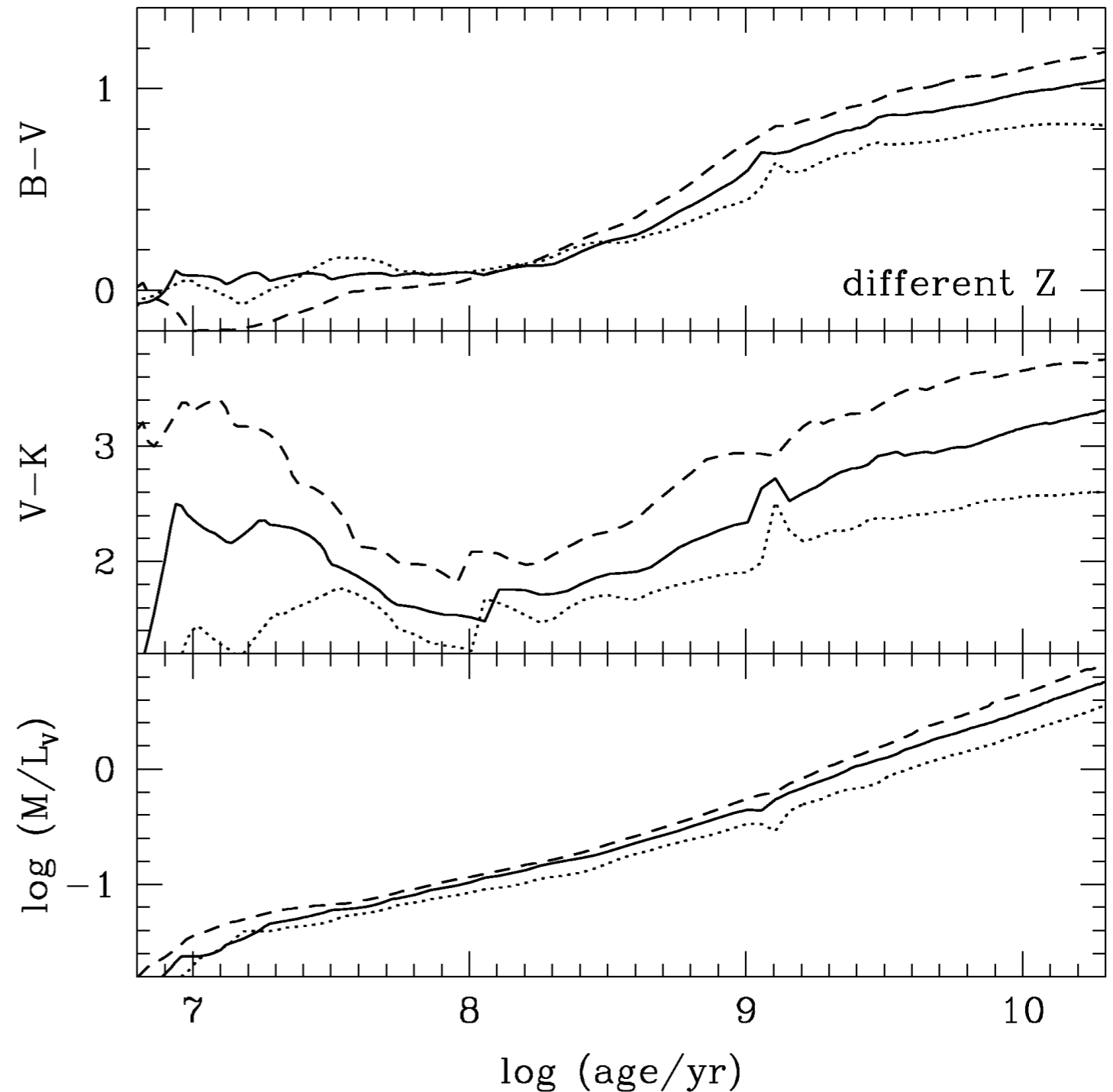
- This is strikingly close to the expression found by Tinsley & Gunn (1976): $(d \ln L/dt) \approx 0.3\alpha - 1.3$



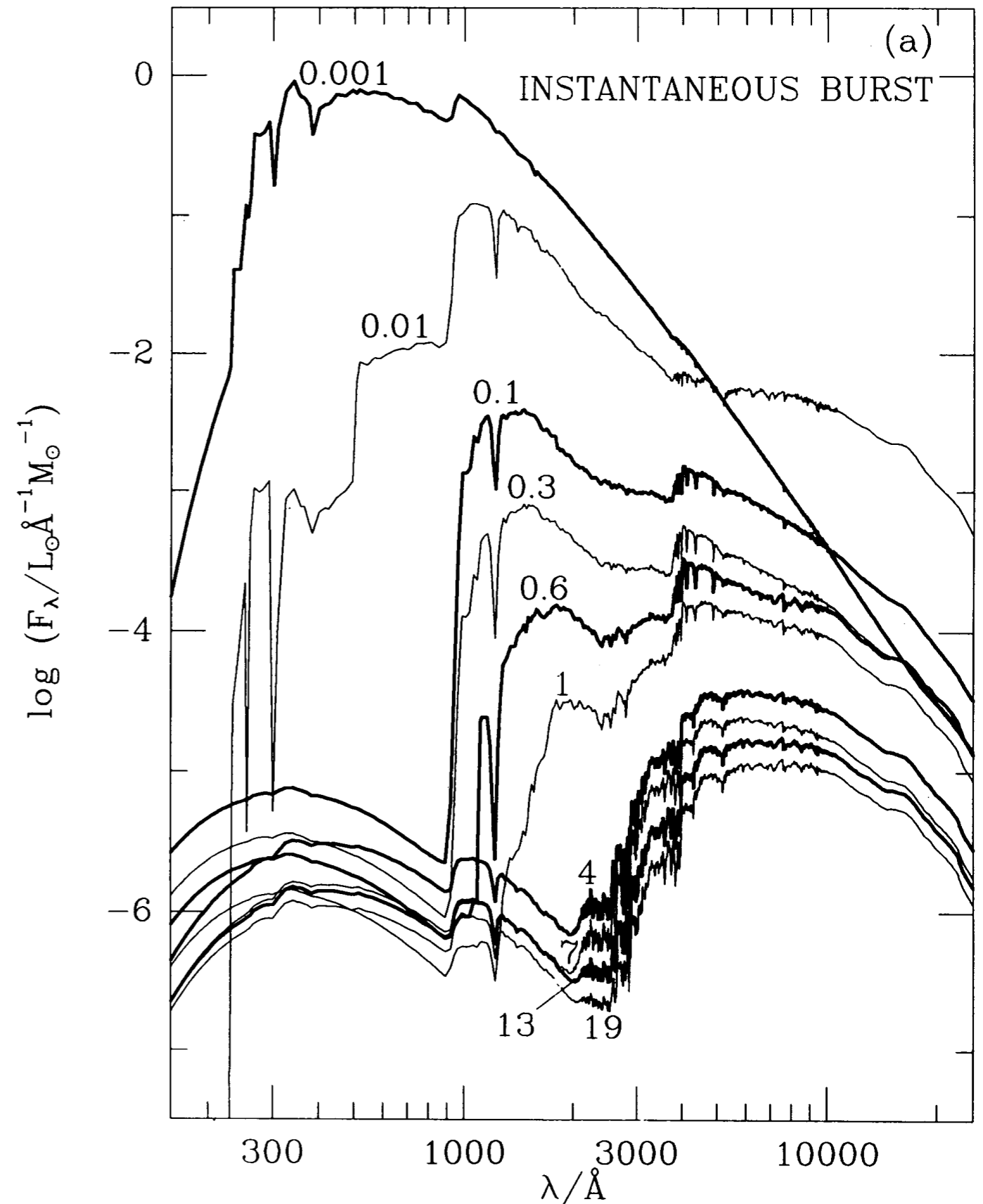
- So as long as $\alpha < 4$, then *the luminosity of a coeval population **decreases with time***
- in other words, in the absence of star formation, galaxies (and star clusters) get fainter with time!

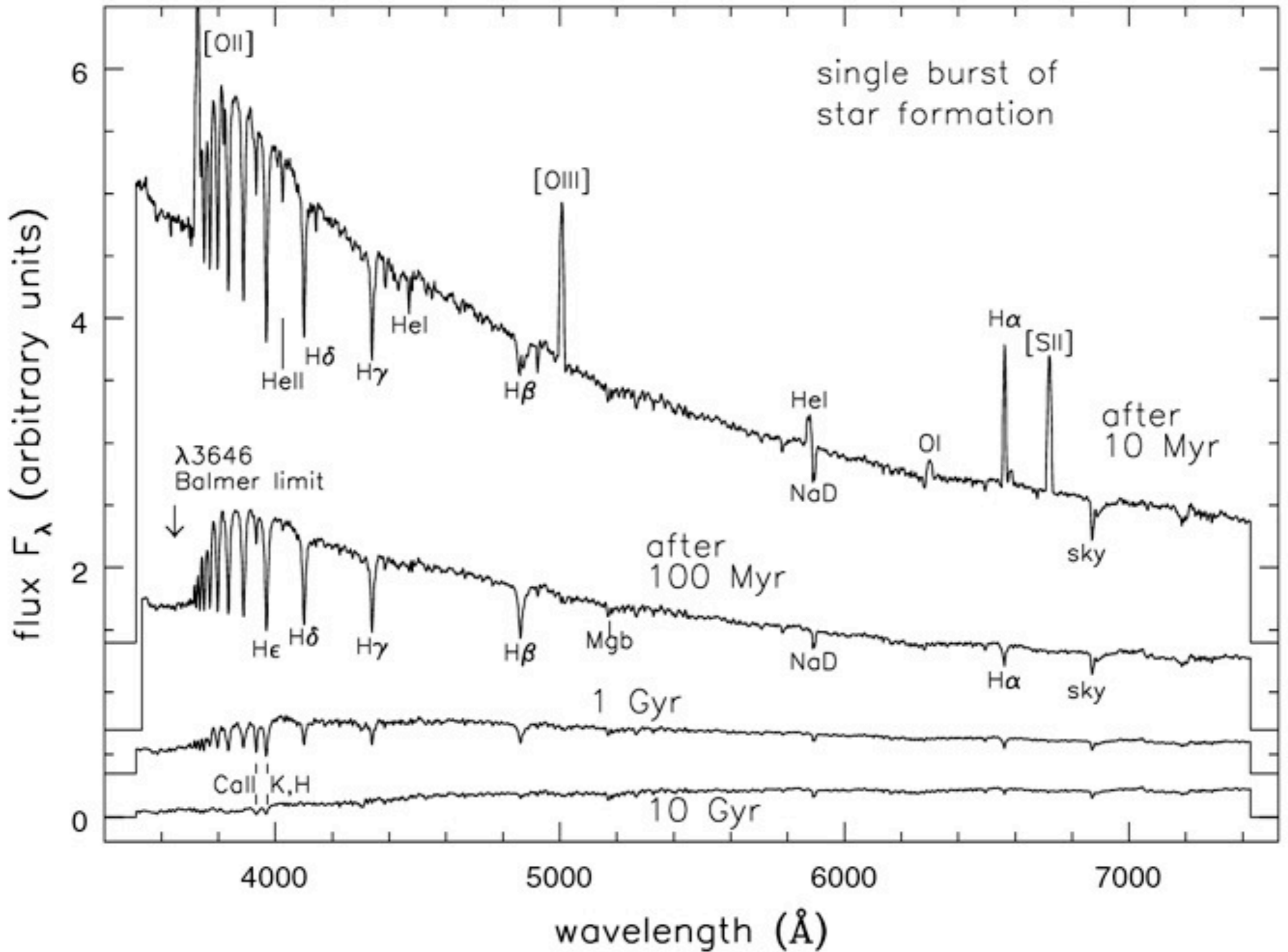


- Stellar population models support this (using Salpeter IMFs)
- note also that populations get redder with age, as expected!

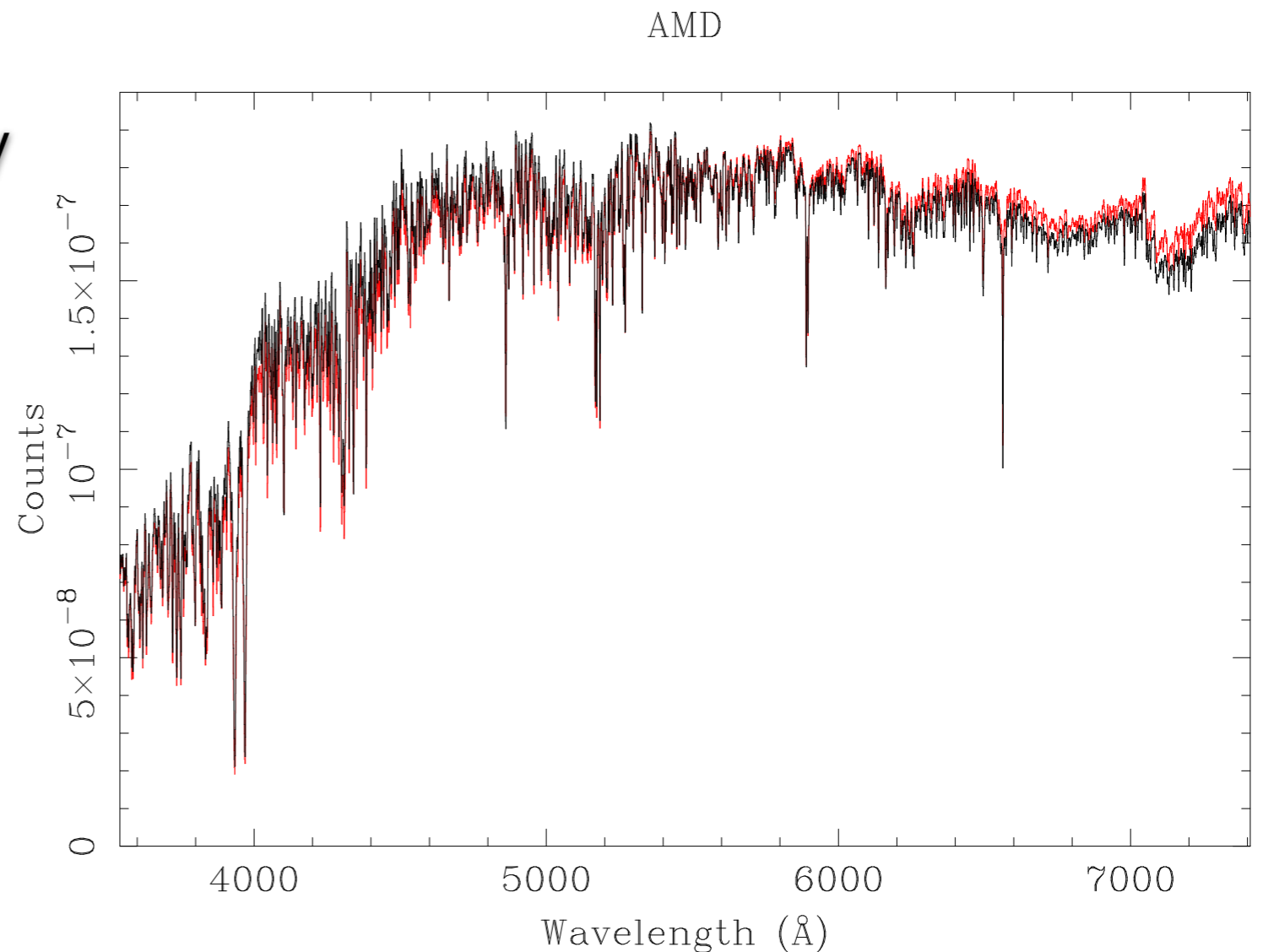


- Spectra do the same thing, of course...





- But watch out for the dreaded age–metallicity degeneracy!
- Two spectra with
 - 5 Gyr, $[Fe/H]=-0.4$
 - 15 Gyr, $[Fe/H]=-0.7$
- Can you tell the difference?



Vazdekis MILES models

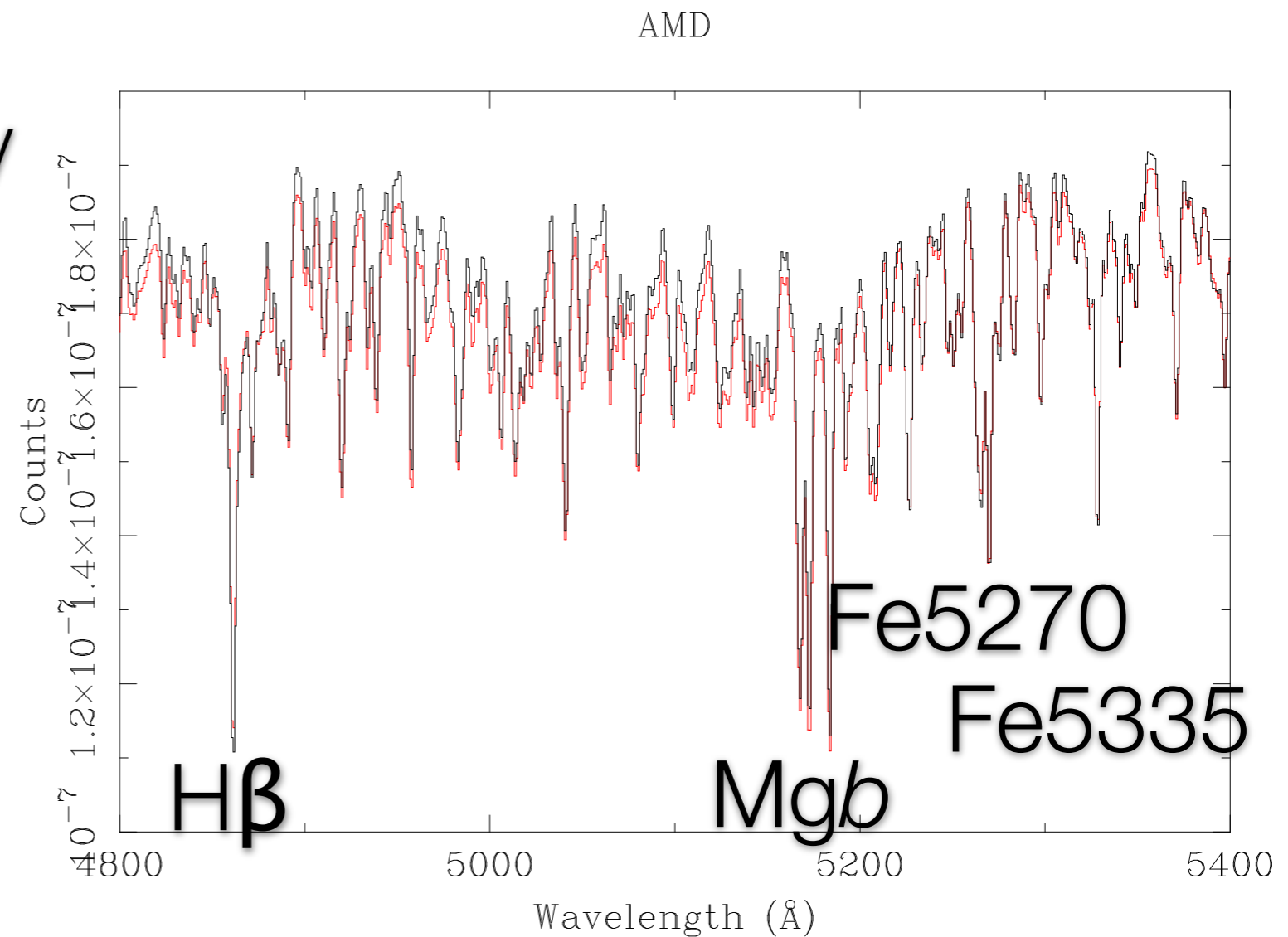


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