

# The Calibration of Monochromatic Far–Infrared Star Formation Rate Indicators.<sup>1</sup>

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## ABSTRACT

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Spitzer data at 24, 70, and 160  $\mu\text{m}$  and ground-based  $\text{H}\alpha$  images are analyzed for a sample of 189 nearby star-forming and starburst galaxies to investigate whether reliable star formation rate (SFR) indicators can be defined using the monochromatic infrared dust emission centered at 70 and 160  $\mu\text{m}$ . We compare recently published recipes for SFR measures using combinations of the 24  $\mu\text{m}$  and observed  $\text{H}\alpha$  luminosities with those using 24  $\mu\text{m}$  luminosity alone. From these comparisons, we derive a reference SFR indicator for use in our analysis. Linear correlations between SFR and the 70  $\mu\text{m}$  and 160  $\mu\text{m}$  luminosity are found for  $L(70) \gtrsim 1.4 \times 10^{42} \text{ erg s}^{-1}$  and  $L(160) \gtrsim 2 \times 10^{42} \text{ erg s}^{-1}$ , corresponding to  $\text{SFR} \gtrsim 0.1\text{--}0.3 \text{ M}_{\odot} \text{ yr}^{-1}$ , and calibrations of SFRs based on  $L(70)$  and  $L(160)$  are proposed. Below those two luminosity limits, the relation between SFR and 70  $\mu\text{m}$  (160  $\mu\text{m}$ ) luminosity is non-linear and SFR calibrations become problematic. A more important limitation is the dispersion of the data around the mean trend, which increases for increasing wavelength. The scatter of the 70  $\mu\text{m}$  (160  $\mu\text{m}$ ) data around the mean is about 25% (factor  $\sim 2$ ) larger than the scatter of the 24  $\mu\text{m}$  data. We interpret this increasing dispersion as an effect of the increasing contribution to the infrared emission of dust heated by stellar populations not associated with the current star formation. Thus, the 70 (160)  $\mu\text{m}$  luminosity can be reliably used to trace SFRs in large galaxy samples, but will be of limited utility for individual objects, with the exception of infrared-dominated galaxies. The non-linear relation between SFR and the 70 and 160  $\mu\text{m}$  emission at faint galaxy luminosities suggests a variety of mechanisms affecting the infrared emission for decreasing luminosity, such as increasing transparency of the interstellar medium, decreasing effective dust temperature, and decreasing filling factor of star forming regions across the galaxy. In all cases, the calibrations hold for galaxies with oxygen abundance higher than roughly  $12 + \text{Log}(\text{O}/\text{H}) \sim 8.1$ . At lower metallicity the infrared luminosity no longer reliably traces the SFR because galaxies are less dusty and more transparent.

*Subject headings:* infrared:galaxies – galaxies: starburst – galaxies: interactions – galaxies: ISM – ISM: structure – stars: formation

## 1. Introduction

The star formation rate (SFR) is one of the principal parameters that needs to be measured in star-forming regions and galaxies, in order to characterize their evolution. Extensive efforts have been made over the past couple of decades to derive SFR indicators from

luminosities at a variety of wavelengths, spanning from the UV, where the recently formed massive stars emit the bulk of their energy, to the infrared, where the dust-reprocessed light from those stars emerges, to the radio, which is mostly a tracer of supernova activity (e.g., Kennicutt 1998; Yun, Reddy & Condon 2001; Kewley et al. 2002; Ranalli, Comastri & Setti 2003; Hirashita, Buat & Inoue 2003; Bell 2003; Kewley, Geller & Jansen 2004; Calzetti et al. 2005; Schmitt et al. 2006; Moustakas, Kennicutt & Tremonti 2006; Alonso–Herrero et al. 2006; Calzetti et al. 2007; Salim et al. 2007; Persic & Rephaeli 2007; Rosa–Gonzalez et al. 2007; Kennicutt et al. 2007; Bigiel et al. 2008; Rieke et al. 2009; Calzetti 2009, for earlier papers, see Kennicutt (1998)).

In recent years, investigations of monochromatic (i.e., based on a single band measurement) SFR indicators based on the infrared emission from galaxies have experienced a new resurgence, thanks to the high–sensitivity and high–angular resolution data provided by the Spitzer Space Telescope, which have yielded both deeper distant galaxy surveys and more accurate information on the relation between dust and stellar emission in nearby galaxies. Deep surveys are often characterized by limited information across the infrared wavelength range and monochromatic SFRs are an important tool for these projects.

The rest–frame mid–infrared emission from dust in galaxies, in particular the emission detected in the 8  $\mu\text{m}$  and 24  $\mu\text{m}$  Spitzer bands, has been analyzed by a number of authors (Roussel et al. 2001; Förster Schreiber et al. 2004; Boselli, Lequeux & Gavazzi 2004; Calzetti et al. 2005; Wu et al. 2005; Alonso–Herrero et al. 2006; Perez–Gonzalez et al. 2006; Relaño et al. 2007; Calzetti et al. 2007; Zhu et al. 2008; Rieke et al. 2009; Salim et al. 2009), and a general correlation (but also a number of caveats) between mid–IR infrared emission and SFR has been found.

Over the next few years, new facilities, both from space (e.g., the Herschel Space Telescope) and from the ground (ALMA and the Large Millimeter Telescope, to mention just two) will open new windows of sensitivity at even longer wavelengths than those explored by Spitzer, and will, in turn, provide even more powerful tools for probing the evolution of the rate at which galaxies have assembled their gas and dust components. Deep surveys will be able to probe the dust emission from galaxies at rest–frame infrared wavelengths that are close to the peak emission, and to the bulk of the infrared energy budget. Herschel will, for instance, enable us to probe the peak dust emission from galaxies ( $\sim 60\text{--}150 \mu\text{m}$ ) up to redshift  $z \sim 2$ , while the Rayleigh–Jeans tail of the dust emission will be the dominion of sub–millimeter and millimeter facilities.

Exploring the viability, and limitations, of using the monochromatic emission close to the infrared peak and at longer wavelengths as SFR indicators is thus timely for providing a reference for those future surveys. In this paper we investigate the use of the 70  $\mu\text{m}$  and

160  $\mu\text{m}$  emission, the two longest wavelength Spitzer bands, from nearby galaxies as SFR diagnostics.

The present paper is organized as follows: Section 2 introduces the sample of local star-forming galaxies and the data used for the present analysis; Section 3 describes the quantities used in this work; Section 4 compares a variety of existing SFR indicators in the optical/infrared, to derive ‘reference’ SFRs for our galaxies, which are then compared with the Spitzer 70  $\mu\text{m}$  and 160  $\mu\text{m}$  emission from the same galaxies and with expectations from models in Sections 5 and 6, respectively. A discussion of the results and the conclusions are given in Section 7. Throughout the paper, we adopt a value of the Hubble constant  $H_0=70 \text{ km s}^{-1} \text{ Mpc}^{-1}$ .

## 2. Sample Description

We combine data from two Spitzer Legacy surveys, LVL (Local Volume Legacy, Dale et al. 2009; Lee et al. 2009a) and SINGS (Spitzer Infrared Nearby Galaxies Survey, Kennicutt et al. 2003), with the local starburst galaxy sample of the Spitzer MIPS GTO programs (Engelbracht et al. 2008) and the Luminous Infrared Galaxies (LIRGs) sample of Alonso–Herrero et al. (2006), in order to span close to 5 orders of magnitude in SFR surface density (SFSD=SFR/area), and thus cover as wide as possible a range that may characterize star-forming galaxies at all redshifts.

Galaxies with MIPS detections at 24, 70, and 160  $\mu\text{m}$  and  $\text{H}\alpha$  measurements are preferentially selected, though, in order to broaden the sample to low-luminosity and low-metallicity objects, galaxies with MIPS upper limits are included if  $\text{H}\alpha$  ( $\lambda 0.6563 \mu\text{m}$ ) or  $\text{P}\alpha$  ( $\lambda 1.876 \mu\text{m}$ ) data exist for them. In general, when both MIPS and  $\text{H}\alpha$  measurements are present, extinction corrected SFRs will be derived by combining the observed 24  $\mu\text{m}$  and  $\text{H}\alpha$  luminosities (see section 4.2); for the few MIPS 24  $\mu\text{m}$  upper limits in our sample, we will rely on the extinction corrected (via the  $\text{H}\alpha/\text{P}\alpha$  ratio)  $\text{P}\alpha$  luminosity, or will assume, for low metallicity objects, that the SFR is practically unobscured by dust and use the observed  $\text{H}\alpha$  luminosity. We also require that measurements of oxygen abundance be available for each galaxy, either from the spectroscopy of Moustakas et al. (2009) or from the literature (Marble et al. 2010; Engelbracht et al. 2008), as the infrared luminosity is sensitive to the metal abundance in galaxies (Cannon et al. 2005; Engelbracht et al. 2005; Cannon et al. 2006a,b; Walter et al. 2007; Calzetti et al. 2007; Relaño et al. 2007), and this is one of the parameters investigated in the present paper. Finally, images in the emission lines of  $\text{H}\alpha$  or  $\text{H}\alpha$ + $[\text{NII}]$  need to be available, to measure the extent of the emitting region in the warm ionized gas; we use the line emitting area to normalize all luminosities, in order to remove

any dependence on distance from our results, and to ensure that galaxies are not scaled by mass (or by global luminosity). In this paper, we use the definition of luminosity surface density ( $LSD = \text{luminosity}/\text{area}$ ), and infrared luminosity surface density (IRSD), to remove any dependence of luminosity on the galaxy distance or size.

The selection criteria and observation strategy for the 258 galaxies in the LVL sample are described in Dale et al. (2009) and Lee et al. (2009a); IRAC and MIPS total fluxes for each galaxy are reported in Dale et al. (2009), while the foreground–galactic–extinction–corrected  $H\alpha + [NII]$  total fluxes and  $[NII]/H\alpha$  ratios and other general information on the galaxies, including distances, are reported in Kennicutt et al. (2008). The main criterion for the LVL sample is that galaxies need to be within the local 11 Mpc volume. Oxygen abundances from the literature are currently available for 108 of the LVL galaxies (Marble et al. 2010); the additional requirement to have narrow–band optical images available for measuring emission region sizes further reduces the sample to 80 galaxies<sup>1</sup>. Of these, 11 are starbursts already included in the MIPS–GTO sample (see below); we thus consider our final LVL sample as consisting of 69 galaxies.

The SINGS sample contains 75 galaxies closer than about 30 Mpc, representative of a wide range of properties in terms of morphology, luminosity, dust temperature, etc. (Kennicutt et al. 2003). Of those 75 galaxies, 33 are in common with the LVL sample. The number of SINGS galaxies included in our sample is further reduced by excluding non–star–forming, early type galaxies, mostly Sy2–dominated ellipticals or S0’s (NGC584, NGC855, NGC1266, NGC1316, NGC1404, NGC4552, NGC5195), and one galaxy, NGC5033, for which optical narrow–band data are not available. We thus have a final sample of 67 SINGS galaxies, 31 of which in common with the LVL sample. Total MIPS fluxes for the SINGS galaxies are listed in Dale et al. (2007), and the  $H\alpha + [NII]$  fluxes plus the  $[NII]/H\alpha$  ratios are listed in Kennicutt et al. (2009). Oxygen abundances are from Moustakas et al. (2009).

In the LVL and SINGS samples we retain four galaxies, M81DwA, HolmbergIX, UGC6900, and UGC9128, which are undetected or marginally detected in the MIPS bands, but have  $H\alpha$  measurements; these galaxies are metal poor, and consequently, dust–poor, thus likely to have most of their star formation light emerge directly at UV/optical wavelengths, unobscured by dust.

The areas used to normalize luminosities and calculate LSDs are derived, for each galaxy, from the area occupied by the ionized gas emission as traced by the  $H\alpha$ . For the LVL and SINGS samples, ionized gas emission sizes are measured directly from our images, and are defined as the semi–major axis of the ellipse that includes 2/3 of the total  $H\alpha$  emission flux.

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<sup>1</sup>The final LVL  $H\alpha$  sample will include 174 galaxies.

Given the uniform nature of the LVL and SINGS imaging strategies, this implies that the sizes obtained for those galaxies are internally consistent. The ‘2/3 ionized gas emission’ radius typically includes all or most of the high surface brightness emitting regions in a galaxy. This definition for the emitting area is, therefore, especially useful for our analysis: it ensures some level of uniformity even when including galaxies for which the  $H\alpha$  image depths vary from object to object. This is the case for some of the galaxies described in the next samples (see below), where the only  $H\alpha$  images available are from the literature.

The MIPS–GTO sample of starbursts is described in Engelbracht et al. (2008): the galaxies cover a wide range in metal content, from very metal–poor starbursts, like IZw18 and SBS0335–052, to metal–rich examples like IC342. Of the 65 (we consider SBS0335–052E and SBS0335–052W a single galaxy in this paper) galaxies listed in Engelbracht et al. (2008), we only keep the 55 that satisfy one of the following two criteria: 1. they are detected in all three MIPS bands (we exclude NGC1614 and NGC3256 which are part of the LIRGs sample, see below); or 2. they are detected and measured in either  $H\alpha$  or  $P\alpha$  (Calzetti et al. 2007). The latter criterion ensures inclusion in the sample of metal–poor galaxies, like IZw18, SBS0335–052, UGCA292, and HS0822+3542, which are marginally or un–detected in MIPS because their dust content is low. Ten of the starbursts in the MIPS–GTO sample have images in the  $P\alpha$  ( $\lambda 1.876 \mu\text{m}$ ) line (Calzetti et al. 2007), which we use to derive SFRs, after extinction correction using the  $H\alpha$  emission, and to measure the size of the line emitting region when  $H\alpha$  images are not available. For much of our analysis, we give emphasis to starbursts with *both* infrared broad–band and optical/nearIR narrow–band detections, since our ‘reference’ SFR indicators generally require information on both (see next section). The distances of the 55 starbursts are typically within 100 Mpc, except for one case (Tol2138–405, located at 246 Mpc). Total MIPS flux values, and literature values for distances, oxygen abundances, and  $H\alpha$  fluxes are listed in Engelbracht et al. (2008). We supplement, where possible,  $H\alpha$  measurements from the literature, when not present in Engelbracht et al. (2008), and we end up with line emission measurements for a total of 45 galaxies.

In order to obtain the emitting areas for the starbursts, we have searched the literature for images of the ionized gas emission (either  $H\alpha$  or  $P\alpha$ ); we were able to retrieve such information for a total of 44 starburst galaxies. For the 34 galaxies with  $H\alpha$  images available through NED<sup>2</sup> in FITS format and/or with  $P\alpha$  images, we derive the ionized gas emission size in the same fashion as for the LVL and SINGS galaxies, with the proviso that non–uniform depths from image to image will reflect into larger uncertainties for the sizes (next

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<sup>2</sup> NED is the NASA/IPAC Extragalactic Database which is operated by the Jet Propulsion Laboratory, California Institute of Technology, under contract with the National Aeronautics and Space Administration.

section). Finally, sizes derived from published hard-copy images only (10 galaxies) have larger uncertainties still, as discussed in section 3. Overall, the values for the sizes of the ionized gas emitting region of the starbursts are the least homogeneous in our sample; variations in sizes of up to 30% may be expected from differences in the depths and characteristics of the published images.

Ionized gas emission sizes for the 160 total LVL (67), SINGS (69, 31 of which are in common with LVL), and starburst galaxies (55) are listed in Table 1, together with other derived quantities relevant to the present analysis.

Our sample of 160 star-forming and starburst galaxies is augmented with 29 LIRGs from Alonso-Herrero et al. (2000, 2001, 2002, 2006), to expand our SFR surface density range at the high end, covering in total almost 5 orders of magnitude with the combined sample. Selection criteria for the sample are presented in Alonso-Herrero et al. (2006), together with measurements of the  $P\alpha$  flux, extinction corrections, as well as information on the physical extent of the star forming area for 24 LIRGs. For a few of the LIRGs with extended ionized gas emission, size information is taken from Hattori et al. (2004). To the 24 LIRGs, we add five additional LIRGs (NGC1614, NGC3256, NGC3690, NGC5653, and Zw049.057) with data from Alonso-Herrero et al. (2000, 2001, 2002). Infrared measurements at 25, 60, and 100  $\mu\text{m}$  from IRAS and distances (which we rescale to our adopted value of the Hubble constant) for each galaxy are from Sanders et al. (2003) and Surace, Sanders & Mazzarella (2004). For two of the LIRGs, NGC1614 and NGC3256, MIPS measurements are available from Engelbracht et al. (2008). Although individual oxygen abundance measurements are not available for the majority of the LIRGs, their metallicities are characteristic of high-metallicity galaxies (Alonso-Herrero et al. 2006). Basic parameters for the 29 LIRGs are listed in Table 2.

The final sample consists of 189 nearby star-forming galaxies, which we divide into a ‘high-metallicity’ sample, consisting of 142 objects with oxygen abundances  $12+\text{Log}(\text{O}/\text{H})\gtrsim 8.1$  (113 normal star-forming and starburst and 29 LIRGs), and a ‘low-metallicity’ sample of 47 objects with oxygen abundance  $12+\text{Log}(\text{O}/\text{H})<8.1$  (Table 1). We do not consider the sample large enough to be reasonably divided into more than two metallicity bins. The separating value in metallicity is chosen after Engelbracht et al. (2005), Jackson et al. (2006), Draine et al. (2007), and Engelbracht et al. (2008), where it is found that galaxies roughly below this metallicity value tend to be underluminous in their PAH emission (as measured in the Spitzer-IRAC 8  $\mu\text{m}$  band). We do not expect the striking trend of the PAH bands to be also present in the thermal-dust-emission-dominated longer wavelength bands. However, we do expect that, as the metal and dust content of galaxies decreases, the galaxies become increasingly transparent, i.e., their dust opacity decreases, and the infrared dust emission

will be a progressively less accurate tracer of the current SFR. We thus select the value  $12 + \text{Log}(\text{O}/\text{H}) \sim 8.1$  as a convenient ‘transition’ value for dust content in our galaxies. In section 7 (Discussion and Conclusions), we also discuss the implications of further dividing the high-metallicity sample into two subsamples, in order to investigate the role of metal abundance in driving some of the observed scatter in the data.

The range of infrared luminosity covered by our sample is shown in Figure 1, together with the range of 70-to-160  $\mu\text{m}$  luminosity ratio, the latter being a rough proxy for effective dust temperature (e.g. Engelbracht et al. 2008). The combination of low-metallicity and high-metallicity star forming and starburst galaxies plus the addition of the LIRGs enables us to explore almost six orders of magnitude in total infrared luminosity, from  $\sim 3 \times 10^{39}$   $\text{erg s}^{-1}$  to  $\sim 3 \times 10^{45}$   $\text{erg s}^{-1}$  (the derivation of fluxes in ‘MIPS-equivalent’ bands for the LIRGs is discussed in the next section). We also explore one order of magnitude range in the 70-to-160  $\mu\text{m}$  luminosity ratio, from about 1/3 to roughly 6 times in  $L(70)/L(160)$ . The LIRGs occupy a relatively small range in luminosity ratio, with values in the range  $L(70)/L(160) \sim 2-5$ , as expected for these warm (relatively to many galaxies) systems.

### 3. Luminosities, Sizes, and Uncertainties

We derive total infrared luminosities,  $L(\text{TIR})$ , in the range 3–1100  $\mu\text{m}$  for the galaxies in the LVL, SINGS, and MIPS-GTO samples using equation 4 in Dale & Helou (2002). Most of the galaxies are also bright enough to have IRAS detections, at least at 25, 60, and 100  $\mu\text{m}$ , but we concentrate on the MIPS data alone. All infrared fluxes are used in this work as directly measured (different from the approach of, e.g., Engelbracht et al. 2008, who use color-corrected fluxes).

For most of the LIRGs, only IRAS infrared data are available, typically with detections at 25, 60, and 100  $\mu\text{m}$ . Exceptions are NGC1614 and NGC3256, which have published MIPS measurements (Engelbracht et al. 2008), and NGC2369, ESO320-G030, and Zw049.057, for which Rieke et al. (2009) published infrared SEDs. Since our reference luminosities are those in the MIPS bands, we need to interpolate the IRAS band fluxes of all the other LIRGs to derive MIPS-equivalent ones. The difference between the luminosity at MIPS 24  $\mu\text{m}$  and at IRAS 25  $\mu\text{m}$  is small (see, also Kennicutt et al. 2009), and we adopt the best fit:

$$\text{Log}[L(24)] = \text{Log}[L(25)] - 0.028, \quad (1)$$

where  $L(\lambda) = 4 \pi d^2 (c/\lambda) f_\nu(\lambda)$ , in units of  $\text{erg s}^{-1}$ , is the monochromatic luminosity at wavelength  $\lambda$ ,  $f_\nu(\lambda)$  is the flux density per unit frequency, and  $d$  is the galaxy’s distance. The difference between the 60  $\mu\text{m}$  IRAS band and the 70  $\mu\text{m}$  MIPS band is larger than



between the 25  $\mu\text{m}$  and 24  $\mu\text{m}$  bands, and we use a polynomial interpolation between the 60 and 100  $\mu\text{m}$  IRAS bands to derive a 70  $\mu\text{m}$ -equivalent flux density:

$$f_\nu(70) = f_\nu(60)(a_0 + a_1x + a_2x^2 + a_3x^3 + a_4x^4 + a_5x^5), \quad (2)$$

with  $f_\nu(\lambda)$  the flux density in Jy,  $x=f_\nu(60)/f_\nu(100)$ , and the vector  $a_n=[1.5078, 0.9356, -3.8900, 4.0576, -1.7617, 0.2712]$ . Tests using starburst galaxies from the MIPS–GTO sample with IRAS 60/100 colors similar to the LIRGs ( $0.4 \leq f_\nu(60)/f_\nu(100) \leq 1.3$ ) show that the above interpolation reproduces the observed  $L(70)$  with a typical uncertainty of less than 10% (i.e., less than 0.04 dex in logarithmic scale). We also compare our interpolated values with MIPS–equivalent fluxes derived from the template SEDs of Rieke et al. (2009), by matching LIRGs to the templates according to TIR luminosity and SED shape (using the IRAS fluxes). We typically find that the values derived from the templates differ from those derived from equations 1 and 2 by  $\sim 1$ –5% at 24  $\mu\text{m}$  and less than 12% at 70  $\mu\text{m}$ .

We use the SED templates also to extrapolate fluxes at 160  $\mu\text{m}$  for the LIRGs with only IRAS measurements. Nine of the 29 LIRGs have long–wavelength measurements at 850  $\mu\text{m}$  from SCUBA on the JCMT (Dunne et al. 2000) and one, NGC3690/IC694, at 1.2 mm from the IRAM–30 m telescope (Braine & Dumke 1998). Interpolations between these long wavelengths and the IRAS 100  $\mu\text{m}$  band to recover the MIPS 160  $\mu\text{m}$  emission are complicated by the presence of a cool,  $T \approx 20$  K, dust component in the IR emission of LIRGs (e.g. Dunne & Eales 2001). The MIPS 160  $\mu\text{m}$  band is close to the emission peak of this cool dust component, thus any interpolation will be subject to the uncertainty of the relative, and unknown, ratio of the cool and warm dust emission components. A simple linear interpolation between 100  $\mu\text{m}$  and 450  $\mu\text{m}$  (using the roughly constant 450–to–850 ratio of 7.9 found by Dunne & Eales 2001) provides estimated 160  $\mu\text{m}$  fluxes that are between 30% and 75% lower than those estimated from the template SEDs, in agreement with the presence of the 20 K emitting dust component. To avoid a significant underestimate in the 160  $\mu\text{m}$  flux, we adopt the results from the extrapolation of the SED templates as our ‘bona fide’ 160  $\mu\text{m}$  fluxes for the LIRGs. The resulting 70–to–160  $\mu\text{m}$  luminosity ratio for the LIRGs has median value  $L(70)/L(160) \sim 2.7$ , to be compared with the median ratio  $L(70)/L(160) \sim 3.5$  of the five galaxies for which direct MIPS measurements are available. The difference is only  $\sim 30\%$  (0.11 dex), implying that our procedure overestimates  $L(160)$  by that amount at most. This is well within the typical uncertainty for our measurements.

For the wavelength–integrated infrared luminosity  $L(\text{TIR})$  of the LIRGs, we compare the results of equation 4 of Dale & Helou (2002), using the MIPS–equivalent fluxes obtained with the SED templates and equations (1) and (2) with the results of their equation 5, using the IRAS bands. The two equations are known to provide different  $L(\text{TIR})$  values for the

more quiescently star forming galaxies and for galaxies with cooler dust (Dale & Helou 2002), because the MIPS 160  $\mu\text{m}$  band traces cool dust more accurately than the IRAS 100  $\mu\text{m}$  band (the two being the longest wavelength bands from Spitzer and IRAS, respectively). The offset is dependent on the 60/100 color (see, e.g. Kennicutt et al. 2009), and  $L(\text{TIR})$  values obtained with the two methods for the starburst galaxies from the sample of Engelbracht et al. (2008) provide an empirical correction:

$$\text{Log}[L(\text{TIR})_{\text{MIPS}}] = \text{Log}[L(\text{TIR})_{\text{IRAS}}] - 0.35\text{Log}[f_{\nu}(60)/f_{\nu}(100)] - 0.15, \quad (3)$$

which gives consistent  $L(\text{TIR})$  values with a scatter typically less than 25%. A list of ‘MIPS–equivalent’ quantities for the LIRGs is given in Table 2.

We finally apply equations (1)–(3) and the SED templates of Rieke et al. (2009) to recover, from the IRAS measurements, the 24, 70, and 160  $\mu\text{m}$  fluxes, and the TIR luminosity of NGC3034<sup>3</sup>, a starburst galaxy included in both the SINGS and LVL samples, but which is saturated in all three MIPS bands.

Galaxy–integrated emission line data include  $\text{H}\alpha + [\text{NII}]$  for the LVL (Kennicutt et al. 2008) and SINGS (Kennicutt et al. 2009) samples; those papers also report the  $[\text{NII}]/\text{H}\alpha$  ratios appropriate for each galaxy. Typical uncertainties on the final integrated  $\text{H}\alpha$  fluxes are about 10%–15%. Similar data are also available for 39 of the MIPS–GTO galaxies, mostly from the literature (Engelbracht et al. 2008), although measurements for 11 galaxies are from the LVL data (Kennicutt et al. 2008), and, for 10 galaxies,  $\text{P}\alpha$  fluxes are also available (Calzetti et al. 2007). Larger uncertainties, of order 20%, are expected for the  $\text{H}\alpha$  fluxes of the starbursts, because of the non–uniform nature of the data. Finally, the only ionized gas data readily available for the LIRGs are the  $\text{P}\alpha$  data presented in Alonso–Herrero et al. (2006), together with estimates of the dust extinction values.

Uncertainties for the line–emitting region sizes range from about 10% for the galaxies for which we perform a direct measure on the images (mostly the LVL and SINGS samples), to  $\sim 15\%$ – $20\%$  for the 34 MIPS–GTO galaxies with available FITS images, to about 20%–30% for the starburst galaxies for which images were available only in hard–copy published form. For the LIRGs, the region size information is from Hattori et al. (2004) and Alonso–Herrero et al. (2006), with an estimated uncertainty of about 20%. For 11 of the MIPS–GTO starbursts, line–emission images are not available, either from archives (e.g., NED) or from the literature. We observe, however, that the size of the star forming region decreases relative to the size of the stellar emission, defined by  $R_{25}$ , as the infrared luminosity surface density ( $\text{IRSD} = L(\text{TIR})/\text{area}$ ) increases (Figure 2). This had already been

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<sup>3</sup>We recover the following flux densities for NGC3034:  $f_{\nu}(24) = 287.4 \text{ Jy}$ ,  $f_{\nu}(70) = 1622. \text{ Jy}$ ,  $f_{\nu}(160) = 639.6 \text{ Jy}$ ; the IRAS 25, 60, and 100  $\mu\text{m}$  data used to derive the ‘MIPS–equivalent’ fluxes are from Sanders et al. (2003).

reported by Wang & Helou (1992) and Lehnert & Heckman (1996), as an effect of increasing compactness of the region for increasing SFSD. Indeed, the infrared luminosity traces the SFR fairly accurately at high luminosities (Hunter et al. 1986; Devereux & Young 1990; Lehnert & Heckman 1996), and only at low luminosities this correlation breaks, i.e., in a regime where the dust column density is low and most of the star formation light is unprocessed by dust; a similar trend to that of Figure 2 is, in fact, observed when the IRSD is replaced by the SFSD in the correlation with  $\text{area}_{25}/\text{area}_{H\alpha}$ . We use the median value measured for starbursts ( $\text{area}_{25}/\text{area}_{H\alpha} \sim 3$ ) to ascribe a star-forming size to the 11 MIPS–GTO starbursts for which only  $R_{25}$  is available (Table 1).

In addition to normalizing luminosities by the emission area, we also normalize by total stellar mass in the galaxy as a way to remove distance dependencies. The  $3.6 \mu\text{m}$  emission from the galaxies is used here as a proxy for the stellar mass, since, at these long wavelengths, the light is dominated by the photospheric emission from low-mass stars; hence the measurement is insensitive to the details of the galaxy’s star formation history, and dust attenuation is small (see, e.g., Nishiyama et al. 2009, for a discussion on the dust extinction curve at these wavelengths). Furthermore, the contribution from dust emission is also negligible, even if the  $3.3 \mu\text{m}$  PAH emission band is included in the Spitzer–IRAC  $3.6 \mu\text{m}$  band (Pahre et al. 2004; Calzetti et al. 2005; Dale et al. 2007; Engelbracht et al. 2008; Leroy et al. 2008). The list of  $3.6 \mu\text{m}$  fluxes for each galaxy is given in the same publications where the MIPS data are listed. For the LIRGs, we use the  $K_S$ -band measurements (retrieved from NED), which, after extinction-correction using the starburst curve of Calzetti et al. (2000), we extrapolate to  $3.6 \mu\text{m}$ -equivalent fluxes using the starburst models from Starburst99 (2007 Update; Leitherer et al. 1999), with ages in the range 20–100 Myr; we note that at these long wavelengths the spectral slope is basically insensitive to age, up to about 10 Gyr. The inferred ‘ $3.6 \mu\text{m}$ -equivalent’ luminosities for the LIRGs correspond to extinction-corrected values, which are the appropriate values to use in our analysis. In the presence of significant extinction at  $3.6 \mu\text{m}$ , the observed luminosities are expected to be lower than those listed in Table 2.

#### 4. Reference Star Formation Rates

When testing the suitability of a luminosity (monochromatic and non-monochromatic) as a SFR indicator, the first task is to decide what the ‘reference’ SFR tracer is going to be for the selected sample. In this work, we give preference to ‘reference’ SFR tracers which incorporate infrared luminosities, as the goal is to test the  $70 \mu\text{m}$  luminosity as a SFR indicator, and remaining close to this wavelength allows us to avoid biases due to large

differences in the dust optical depth probed.

In addition to the historical SFR tracers based on the bolometric infrared emission (Hunter et al. 1986; Lonsdale Persson & Helou 1987; Rowan–Robinson & Crawford 1989; Devereux & Young 1990; Sauvage & Thuan 1992; Lehnert & Heckman 1996; Kennicutt 1998), a number of calibrations have appeared in the literature in recent years based on the Spitzer MIPS 24  $\mu\text{m}$  monochromatic band emission, or, equivalently, on the IRAS 25  $\mu\text{m}$  emission (e.g., Calzetti et al. 2005; Wu et al. 2005; Alonso–Herrero et al. 2006; Perez–Gonzalez et al. 2006; Calzetti et al. 2007; Zhu et al. 2008; Rieke et al. 2009; Kennicutt et al. 2009). Composite indicators based on a combination of the  $\text{H}\alpha$  optical recombination line emission or UV stellar continuum and one infrared band emission, at 8  $\mu\text{m}$  or 24  $\mu\text{m}$ , have also been calibrated (Calzetti et al. 2007; Kennicutt et al. 2007; Zhu et al. 2008; Leroy et al. 2008; Kennicutt et al. 2009), since the infrared bands capture the SFR from regions obscured by dust, while the optical recombination line (or UV stellar continuum) captures the SFR from regions that are relatively unobscured by dust. The composite indicators are indeed the most appropriate to use in the case of low–metallicity galaxies, where a large fraction of the star formation emerges virtually unaffected by dust re–processing.

We intentionally avoid SFR indicators based on the total infrared emission,  $L(\text{TIR})$ , because the 70  $\mu\text{m}$  wavelength region is close to the peak emission for many galaxies, and  $L(70)$  is thus a significant contributor to  $L(\text{TIR})$ ; this causes a degeneracy between  $L(70)$  and  $L(\text{TIR})$ . Furthermore, the dust contributing to the long–wavelength infrared emission can be heated by evolved stellar populations, as well as by the currently star–forming ones; the ‘extra’ contribution to the infrared emission will produce an overestimate of the true SFR (Lonsdale Persson & Helou 1987; Sauvage & Thuan 1992; Buat & Xu 1996). Although the magnitude of this contribution is uncertain and its impact on the infrared emission not unambiguously established (see, e.g., Kewley et al. 2002), its potential dependence on the relative ratio of evolved/star–forming stellar populations in the galaxy is sufficient reason for not using  $L(\text{TIR})$  as a reference SFR in the present work.

#### 4.1. The Infrared Spectral Energy Distribution of Star–forming Galaxies

Most findings indicate that the correlation between  $L(24)$  and SFR is non–linear, in the sense that the infrared emission at 24  $\mu\text{m}$  is overluminous in more actively star–forming systems relative to a linear scaling between  $L(24)$  and SFR. Two non–exclusive interpretations are possible for this result. The first interpretation attributes the overluminosity at the bright end of the 24  $\mu\text{m}$  emission to the higher mean dust temperatures of more active systems (Calzetti et al. 2007); in this case, the infrared spectral energy distribution (SED)

becomes ‘bluer’ and  $L(24)/L(\text{TIR})$  increases at higher SFRs. The second interpretation suggests that, as the dust optical depth increases with increasing SFR, standard dust extinction correction methods based, e.g., on the Balmer decrement or other hydrogen line ratios, are insufficient to recover the intrinsic values of optical or near-infrared indicators ( $H\alpha$ ,  $P\alpha$ , etc.), thus artificially leading to the interpretation that  $L(24)$  is overluminous relative to those indicators (Rieke et al. 2009).

In order to investigate the nature of the  $L(24)$  overluminosity, we plot the ratio  $L(24)/L(\text{TIR})$  as a function of the infrared luminosity surface density (Figure 3, left). A general increase of the  $24\ \mu\text{m}$  luminosity as a function of IRSD is observed for normal and starburst galaxies, both at low and high metallicities; however, a significant break in the trend is observed for the LIRGs in our sample. These galaxies show significantly ‘cooler’ infrared SEDs (lower  $L(24)/L(\text{TIR})$ , by roughly 0.4–0.5 dex) than one would infer from their luminosity surface density if they followed the same trend as normal star forming and starburst galaxies. Low metallicity galaxies mark a trend similar to that of the high metallicity ones, but at lower average IRSD, as expected if the low metallicity galaxies have lower dust content and lower global infrared emission at fixed  $L(24)/L(\text{TIR})$  (Cannon et al. 2005, 2006a,b; Walter et al. 2007). Although some of the low metallicity galaxies are hotter than the ‘hottest’ high metallicity galaxies, we do not consider this significant enough in our sample, in light of the broad scatter in properties shown by our low metallicity galaxies.

The LIRGs still show a deviation, albeit less extreme than in the previous case, from the trend of the other star-forming and starburst galaxies when the ratio  $L(24)/L(\text{TIR})$  is plotted as a function of  $L(\text{TIR})/L(3.6)$ , i.e., the infrared luminosity per unit stellar mass (Figure 3, right). Thus, the observed deviation is not due to a bias in our choice of the emitting region, but is still present when a different normalization to the infrared luminosity is used. An overestimate of  $L(\text{TIR})$  for the LIRGs may potentially produce the observed deviation, but we would need such an overestimate to be in the range 3–100 $\times$ , in order to reconcile the LIRGs with the other starburst galaxies in Figure 3; this is unlikely to be the case, as our typical uncertainty on  $L(\text{TIR})$  is around 25% (see discussion in section 3).

We concentrate the rest of the discussion on the high-metallicity sample, to avoid the impact of low dust content on the observed IR luminosity. Figure 4 (left) shows the comparison between data and expectations from models for  $L(24)/L(\text{TIR})$  as a function of the infrared luminosity surface density. For clarity, in the left panel of Figure 4 we also show the data after they are binned in 1 dex intervals along the IRSD axis (for two binning schemes shifted by 0.5 dex relative to each other). The models are the same as described in Calzetti et al. (2007), which we summarize briefly here. Stellar populations SEDs are from the Starburst99 models, for a stellar initial mass function (IMF) given by a broken

power law as described in Kroupa (2001) (see, also, Chabrier 2003). The dust extinction and geometry are modeled with the empirical prescription of Calzetti (2001), which is accurate for starburst galaxies, but will lose accuracy both at the low end and high end of the SFSD range (Calzetti, Kinney & Storchi–Bergmann 1994; Meurer, Heckman & Calzetti 1999; Calzetti et al. 2000; Buat et al. 2002; Goldader et al. 2002; Bell 2003; Buat et al. 2005; Calzetti et al. 2005; Seibert et al. 2005). However, we expect, to first approximation, that this simple prescription will at least guide in evaluating trends. The infrared SED is modeled according to the prescription of Draine & Li (2007), according to which the fraction of infrared light emerging in each band is a function of the starlight intensity. The starlight intensity in the Draine & Li (2007) models is related to the SFSD of the stellar population models: larger SFSDs correspond to higher starlight intensity, which in turn produces warmer infrared emission, i.e., higher  $L(24)/L(\text{TIR})$  (Calzetti et al. 2007).

Calzetti et al. (2007) model sub-kpc regions in galaxies, thus comparisons with whole galaxies need to be performed with some care. The fiducial models from Calzetti et al. (2007), adopt as default a 100 Myr constant star formation population attenuated by the starburst curve of Calzetti (2001). This requires differential extinction between stars and ionized gas, with  $E(B-V)_{\text{star}} \sim 0.44 E(B-V)_{\text{gas}}$ ; furthermore, the models relate the intensity of the SFSD to the opacity of the medium, in the sense that more actively star-forming systems also tend to be more dust-obscured and stronger infrared emitters (Wang & Heckman 1996; Heckman et al. 1998; Hopkins et al. 2001; Calzetti 2001). While this may be true on average, whole galaxies may encompass a wider range of dust emission characteristics. For instance, in sub-kpc regions or HII regions larger SFRs (and higher IR emission) imply higher stellar field intensities, which produce hotter infrared SEDs (Draine & Li 2007). In a galaxy, higher IR emission can also be accomplished with larger filling factors of the infrared emitting regions; conversely, lower IRSDs do not necessarily need to be accompanied by an increased transparency of the interstellar medium. Whole galaxies offer a wider range of physical scenarios than HII regions also in terms of the stellar population mix.

Indeed, we can span a wider range in the parameter space shown in Figure 4 if we add additional simple scenarios to our ‘default’ model. We include, as examples, the cases of: (1) constant and large ( $A_V > 10$  mag) dust opacity, independent of the SFSD; (2) 10% filling factor by area of the SFR relative to the total area considered (e.g., disk galaxies, where star formation is concentrated along the spiral arms, and there is little star formation in the interarm regions); and (3) a variation in the underlying stellar population, where the fiducial 100 Myr old constant star formation population of Calzetti et al. (2007) is accompanied by both 1 Gyr and 10 Gyr constant star formation models and by stellar population models with exponentially decreasing star formation and e-folding times of  $\tau = 2$  Gyr and  $\tau = 5$  Gyr. We finally add a model in which the decreasing star formation population with 2 Gyr e-folding

time is homogeneously mixed with dust; here the extinction curve is adopted to be the mean Milky Way one, with stars and ionized gas affected by the same dust column densities (Calzetti 2001), i.e., no differential extinction is introduced. These scenarios account for much of the observed dispersion of the star-forming and starburst galaxies in Figure 4 (left). For instance, galaxies that have very low  $L(24)/L(TIR)$  ratios for  $IRSD > 10^{41} \text{ erg s}^{-1} \text{ kpc}^{-2}$  are generally of morphological type Sab or earlier; their infrared SEDs (see, also, Figure 5) are best described by a model of decreasing star formation with the stellar population mixed with the dust (cyan dashed lines in both Figures 4 and 5). In these types of galaxies, in fact, the bolometric emission tends to be dominated by evolved stellar populations; these populations are likely responsible for a non-negligible portion of the infrared luminosity and tend to heat the dust to colder temperatures than actively star-forming populations (e.g., Helou 1986).

The LIRGs tend to have their infrared emission dominated by the current burst of star formation (Scoville et al. 2000), thus by relatively young stellar populations with ages  $\lesssim 100$  Myr. Yet, their  $L(24)/L(TIR)$  ratios fall outside of the range spanned by the appropriate models of Figure 4 (Left). This ratio is lower than expected for the measured infrared emission even when the latter is normalized to the galaxy’s stellar mass (Figure 4, right). Our models (see above) go through the locus occupied by the LIRGs for constant star-forming populations of  $\approx 100$  Myr–1 Gyr age. However, this apparent agreement is misleading, since LIRGs have evolved underlying galaxies, and the stellar populations contributing to the  $3.6 \mu\text{m}$  emission need to be at least a few Gyr old (Scoville et al. 2000), requiring a shift by at least a factor of 2 along the horizontal axis of Figure 4, right-hand-side panel.

Supporting evidence that the LIRGs in our sample have cool SEDs, comparable to those of less luminous star-forming/starburst galaxies is provided by the dependence of the  $L(70)/L(TIR)$  and the  $L(160)/L(TIR)$  ratios on the IRSD (Figure 5). The models of Draine & Li (2007), combined with the prescription of Calzetti et al. (2007) to relate the SFSD to the starlight intensity, predict that by the time the IRSD of LIRGs is reached, the peak of the IR SED should have moved out of the MIPS  $70 \mu\text{m}$  band, towards shorter wavelengths, thus following first an increase in the  $L(70)/L(TIR)$  ratio from low-luminosities, to a peak around  $\Sigma_{TIR} \approx 10^{42} \text{ erg s}^{-1} \text{ kpc}^{-2}$ , to be followed by a steady decrease. Similarly, the same models predict a steady decrease of the contribution of the  $160 \mu\text{m}$  emission to the total far-infrared luminosity as the peak emission moves to shorter wavelengths at higher luminosity. The data are, instead, consistent with a roughly constant  $L(70)/L(TIR)$  over, at least, 4 orders of magnitude from  $\Sigma_{TIR} \sim 10^{41} \text{ erg s}^{-1} \text{ kpc}^{-2}$  to  $\sim 10^{45} \text{ erg s}^{-1} \text{ kpc}^{-2}$ , and possibly over the full range covered by our data. The trend and dispersion of the data at low IRSD ( $\Sigma_{TIR} < 10^{42} \text{ erg s}^{-1} \text{ kpc}^{-2}$ ) can be, at least partially, accounted for by the model scenarios discussed above. The high luminosity trend, on the other hand, cannot

be accounted for by any of our simple models. A similar argument can be applied to the  $L(160)/L(\text{TIR})$  ratio (Figure 5, right), which, after a decrease reasonably consistent with the models for increasing luminosity up to  $\Sigma_{\text{TIR}} \sim 10^{42.5} \text{erg s}^{-1} \text{kpc}^{-2}$ , flattens out for the highest values of the IRSD, indicating the presence of an extra contribution from cool dust to the infrared SED. When combined with the findings for  $L(24)/L(\text{TIR})$  (Figure 4), we infer that, for the LIRGs, the energy absorbed by the dust itself in the mid-IR (self-absorption; Rieke et al. 2009) is re-emitted at longer infrared wavelengths, thus producing the cool IR SEDs observed by Dunne & Eales (2001). For the purpose of this work, we will modify the models of Draine & Li (2007), by imposing that  $L(70)/L(\text{TIR}) = \text{constant} = 0.5$  and  $L(160)/L(\text{TIR}) \sim \text{constant} = 0.21$  for  $\Sigma_{\text{TIR}} > 10^{42.5} \text{erg s}^{-1} \text{kpc}^{-2}$  (indicated by the black lines in both panels of Figure 5).

The overall conclusion from the above discussion is that normal star-forming and starburst galaxies tend to have ‘hotter’ dust, i.e., increasing  $L(24)/L(\text{TIR})$  ratios, for increasing IRSD, in agreement with the predictions of Draine & Li (2007) coupled with a direct correlation between SFR and stellar field intensity (Calzetti et al. 2007). However, the trend ‘flattens’ around an  $\text{IRSD} \sim 0.3 - 1 \times 10^{43} \text{erg s}^{-1} \text{kpc}^{-2}$  (Figure 4, left). This roughly corresponds to  $L(\text{TIR}) \approx 1 - 3 \times 10^{44} \text{erg s}^{-1} \sim 3 - 8 \times 10^{10} L_{\odot}$  or  $L(24) \approx 2 - 6 \times 10^{43} \text{erg s}^{-1} \sim 0.6 - 2 \times 10^{10} L_{\odot}$ . The conversion between luminosity/area and luminosity should, however, be taken with caution, as it carries a large scatter, almost an order of magnitude. Galaxies with IRSDs or luminosities above these values, which in our sample include mostly LIRGs, show some evidence for flattening or decreasing  $L(24)/L(\text{TIR})$  and flattening  $L(70)/L(\text{TIR})$  and  $L(160)/L(\text{TIR})$  values with increasing IRSD, corresponding to ‘cooler’ dust emission than expected from the intensity of the infrared emission (Dunne et al. 2000; Dunne & Eales 2001). This trend is likely a manifestation of self-absorption by dust even at  $24 \mu\text{m}$  (Rieke et al. 2009, see the Discussion section for a quantification of this effect). An inflection in the  $L(24)/L(\text{TIR})$  trend above  $L(\text{TIR}) \sim 10^{11} L_{\odot}$  has indeed been previously observed by Rieke et al. (2009, see their Figure 8). As suggested by these authors, presence of self-absorption even at  $24 \mu\text{m}$  in LIRGs will hamper the usefulness of standard SFR indicators based on optical-infrared hydrogen recombination lines at these high luminosity values (see below).

## 4.2. Comparisons Between Calibrations

Several calibrations of SFRs using the  $24 \mu\text{m}$  emission of galaxies (or star-forming regions in galaxies) have been recently published, using both simple proportionality between SFR and  $L(24)$  and non-linear relations. The calibrations are based on a variety of galaxy and/or region samples, and stellar IMFs. For sake of comparison among the different cali-



brations, we report below most of the published ones, all of them converted to a common IMF and luminosity scale. In this section, we will use luminosity, rather than our preferred (because distance- and size-independent) luminosity/area, since most published calibrations are derived from luminosities. We choose the Kroupa (Kroupa 2001) IMF, which has a slope  $-2.3$  for stellar masses in the range  $0.5\text{--}100 M_{\odot}$  and  $-1.3$  for stellar masses in the range  $0.1\text{--}0.5 M_{\odot}$ . For this IMF, Starburst99 stellar population models with solar metallicity and constant SFR give the following relation between number of ionizing photons and SFR:

$$SFR(M_{\odot} yr^{-1}) = 7.41 \times 10^{-54} N_{ion}(s^{-1}), \quad (4)$$

which corresponds to<sup>4</sup> :

$$SFR(H\alpha)(M_{\odot} yr^{-1}) = 5.45 \times 10^{-42} L(H\alpha)(erg s^{-1}), \quad (5)$$

where the  $H\alpha$  luminosity is extinction-corrected. For reference, adopting a Salpeter (1955) IMF in the stellar mass range  $0.1\text{--}100 M_{\odot}$  would increase the calibration coefficient in equation 5 by a factor 1.51.

Linear relations between SFR and  $L(24)$  have been published by Wu et al. (2005); Zhu et al. (2008); Rieke et al. (2009). Reported on a common scale, the three relations<sup>5</sup> are:

$$SFR(24)_{W05,l} = 2.75 \times 10^{-43} L(24) \quad 1 \times 10^{42} \lesssim L(24) \lesssim 1 \times 10^{44}, \quad (6)$$

$$SFR(24)_{Z08,l} = 2.46 \times 10^{-43} L(24) \quad 4 \times 10^{41} \lesssim L(24) \lesssim 2 \times 10^{44}, \quad (7)$$

$$\begin{aligned} SFR(24)_{R09} &= 2.04 \times 10^{-43} L(24), & 4 \times 10^{42} &\leq L(24) \leq 5 \times 10^{43}, \\ &= 2.04 \times 10^{-43} L(24) \times [2.03 \times 10^{-44} L(24)]^{0.048} & L(24) > 5 \times 10^{43}, \end{aligned} \quad (8)$$

where the  $24 \mu\text{m}$  luminosities are in units of  $\text{erg s}^{-1}$ . The linear (with a small non-linear correction at the high luminosities for the derivation of Rieke et al. 2009) relations are averages

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<sup>4</sup>The calibration coefficient in equation 5 is about 3% higher than what reported in Calzetti et al. (2007, their equation 6). The present calibration is based on a  $>1$  Gyr age constant star formation stellar population, while the calibration in Calzetti et al. (2007) is based on a 100 Myr age constant star formation population. This difference is, however, smaller than the  $\sim 12\%$  uncertainty to be expected for variations in the physical conditions of the emitting regions, i.e., variations in electron temperature between 5,000 K and 15,000 K, consistent with the variations in metallicity typical of our galaxy sample (Osterbrock & Ferland 2006).

<sup>5</sup>The luminosity range of applicability for the calibrations by Wu et al. (2005) and Zhu et al. (2008) is not explicitly provided by the authors. We report ranges as derived from the luminosity range of the main samples from those authors.

of the calibrations derived in each paper. The galaxy samples and the reference calibrators are different in the three papers: Wu et al. (2005) use both radio and the extinction-corrected  $H\alpha$  luminosity of the galaxies in common between the SDSS and the Spitzer First Look Survey sample; Zhu et al. (2008) use extinction-corrected UV and  $H\alpha$  luminosities and FIR emission of the galaxies in the SDSS and the Spitzer SWIRE galaxy sample (Lonsdale et al. 2003); Rieke et al. (2009) use both  $P\alpha$  and FIR emission on templates of LIRGs and ULIRGs with data of various origins. Despite these differences, the three calibrations are remarkably close, within  $\sim 10\%$ – $35\%$  of each other, consistent with the calibration uncertainties of 0.1–0.15 dex quoted by the authors.

Linear calibrations are derived under the assumption that the  $24\ \mu\text{m}$  luminosity increases proportionally to the SFR. The alternative scenario discussed in the previous section predicts the  $24\ \mu\text{m}$  luminosity to increase proportionally faster than the SFR, as a result of the increasing mean dust temperature, and corresponding shift to bluer wavelengths of the infrared SED. This trend may simply imply, for whole galaxies, a progressive shift for the dominant population heating the dust from evolved stars to young stellar populations; this argument is expected to hold on average, rather than on a galaxy-by-galaxy basis. As seen in section 4.1, the effective dust temperature appears to reach a plateau as a function of increasing TIR luminosity around  $L(24) \sim \text{a few} \times 10^{43}\ \text{erg s}^{-1}$ .

Direct fits of  $L(24)$  as a function of a variety of reference SFR indicators do show such non-linear behavior (Wu et al. 2005; Alonso-Herrero et al. 2006; Perez-Gonzalez et al. 2006; Calzetti et al. 2007; Zhu et al. 2008), and the relation between SFR and  $L(24)$  requires the mediation of a power law with less-than-unity values. Two of the published calibrations involving whole galaxies are from Wu et al. (2005); Zhu et al. (2008), which expressed in terms of our default IMF become:

$$SFR(24)_{W05} = 1.07 \times 10^{-38} L(24)^{0.893} \qquad 1 \times 10^{42} \lesssim L(24) \lesssim 1 \times 10^{44}, \quad (9)$$

$$SFR(24)_{Z08} = 8.10 \times 10^{-37} L(24)^{0.848} \qquad 4 \times 10^{41} \lesssim L(24) \lesssim 2 \times 10^{44}, \quad (10)$$

where the units are the same as the previous linear equations. Calibrations involving either HII-emitting regions in galaxies (Perez-Gonzalez et al. 2006; Calzetti et al. 2007; Relaño et al. 2007) or combinations of HII regions and LIRGs (Alonso-Herrero et al. 2006) (called ‘HII-regions’ calibrations from now on) are given by<sup>6</sup> :

$$SFR(24)_{P06} = 9.01 \times 10^{-34} L(24)^{0.768} \qquad 1 \times 10^{38} \lesssim L(24) \lesssim 3 \times 10^{41}, \quad (11)$$

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<sup>6</sup>As in the case of Wu et al. (2005) and Zhu et al. (2008), most authors do not explicitly provide the luminosity range of validity of their SFR calibrations. As before, we report limits derived from the luminosity range of the main samples from each paper.

$$SFR(24)_{A06} = 5.83 \times 10^{-38} L(24)^{0.871} \quad 1 \times 10^{40} \lesssim L(24) \lesssim 3 \times 10^{44}, \quad (12)$$

$$SFR(24)_{C07} = 1.31 \times 10^{-38} L(24)^{0.885} \quad 1 \times 10^{40} \lesssim L(24) \lesssim 3 \times 10^{44}, \quad (13)$$

$$SFR(24)_{Re07} = 5.66 \times 10^{-36} L(24)^{0.826} \quad 1 \times 10^{38} \lesssim L(24) \lesssim 3 \times 10^{44}. \quad (14)$$

Typical uncertainties for the calibrations of equations 9–14 are 0.01–0.02 in the exponent ( $\sim 0.06$  for Wu et al. 2005)) and 6%–15% in the calibration constant. Within these uncertainties, and within the typical scatter of 0.2–0.3 dex of galaxy and HII regions samples, all calibrations mostly agree with each other, when limited to the luminosity ranges where they were derived. Deviations from one to the other are around a factor  $\sim 2$  (Figure 6, left panel). However, customary use does not carry the uncertainties for the SFR calibrations; this can produce larger discrepancies than a factor  $\sim 2$  when the calibrations are extrapolated beyond their range of validity, as exemplified by the left panel of Figure 6.

Galaxy-wide SFR(24) non-linear calibrations (Wu et al. 2005; Zhu et al. 2008) are marginally, 25%–60%, albeit systematically, higher than HII regions SFR(24) calibrations at the high luminosity end (Figure 6, left panel). The same is observed for galaxy-wide linear calibrations (equations 6–8, and Figure 6, right panel), with deviations around a factor  $\sim 2$  relative to HII regions non-linear calibrations, which is, again, systematic, albeit marginal within the uncertainties. In order to account for this difference, we can speculate that, in whole galaxies, the contrast between the contributions of recent star formation and old stellar populations in the heating of dust is lower than in their HII regions. This possibly leads to a lower average dust temperature (and lower, on average,  $L(24)/L(TIR)$ ) and, consequently, to a higher calibration constant for SFR(24). However, this speculation cannot be easily verified at this time; it will have to await for observations with the Herschel Space Telescope, in order to obtain long infrared wavelength images of sufficient high angular resolution to isolate HII-emission-dominated regions in external galaxies.

The comparison of the linear SFR(24) calibrations (equations 6–8, all of them for whole galaxies) with the non-linear ones (Figure 6, right, compared with the non-linear calibration of Relaño et al. 2007) shows, again, general agreement within the existing uncertainties, when considered in the luminosity range for which those calibrations were derived. However, extending the linear calibrations to the lowest luminosity values of our high-metallicity sample show significant discrepancies from the non-linear calibrations; specifically, the linear calibrations predict SFRs that are a factor  $\sim 2$ –4 lower than predicted by the non-linear calibrations. This is expected in light of the fact that galaxies become increasingly transparent for decreasing infrared luminosity, and a larger fraction of the light from recent star formation emerges directly at UV and optical wavelengths. Thus, the linear SFR(24) calibrations will underestimate the true SFR below  $L(24) \approx 3 \times 10^{42}$  erg s<sup>-1</sup> (Figure 6, right) .

Non-linear calibrations may provide a better alternative for extending the SFR(24) to low luminosity values, because their functional form accounts for two physical effects: increasing transparency of the medium and lower effective dust temperature towards lower SFSDs (Draine & Li 2007). However, they are also potentially problematic. The first problem is that for sufficiently low dust content, even the non-linear SFR(24) will end up underestimating the true SFR (Calzetti et al. 2007). The second problem is that low global SFRs can be obtained both for less active, metal rich (but gas-poor) galaxies and for active, low-metallicity, low-luminosity galaxies. In the latter case, as the metal content of the system decreases, the effective dust temperature increases (e.g. Hunter et al. 1989; Walter et al. 2007; Engelbracht et al. 2008, and Figure 4), thus complicating any trend. A third problem is that non-linear calibrations between SFR and luminosity are ‘object-dependent’, in the sense that their validity relies on the presence of the physical conditions that lead to the non-linear relation itself. Thus, calibrations obtained for whole galaxies will in general not be applicable to regions within those galaxies, and viceversa.

Indeed, the decreasing dust opacity of galaxies for decreasing SFR (Wang & Heckman 1996; Calzetti et al. 2007) calls into question the usefulness of an infrared-based SFR indicator at the low luminosity end of our sample, especially in the low-metallicity regime. For many of the galaxies in our sample star formation appears to ‘shine’ unimpeded by dust absorption. This is evidenced in Figure 7, where the ratio of  $H\alpha$  to the  $24\ \mu\text{m}$  luminosity (rescaled according to the proportionality factor of Kennicutt et al. 2009, between the two luminosities) is plotted as a function of the SFR as defined in Kennicutt et al. (2009) for the normal star forming and starburst galaxies in our sample. For most of these galaxies, the amount of star formation emerging at optical wavelengths (as traced by  $H\alpha$ ) is larger than the star formation emerging in the infrared (as traced by the  $24\ \mu\text{m}$  emission). For this reason, composite SFR indicators, involving the combination of an optical and an infrared tracer of SFR, have been proposed (Calzetti et al. 2007; Kennicutt et al. 2007; Zhu et al. 2008; Kennicutt et al. 2009). For whole galaxies, Kennicutt et al. (2009) propose (see, also Zhu et al. 2008):

$$SFR_{mix,K09} = 5.45 \times 10^{-42} [L(H\alpha)_{obs} + 0.020L(24)] \quad 3 \times 10^{38} \lesssim L(24) \lesssim 3 \times 10^{44}, \quad (15)$$

where the calibration is expressed in terms of our selected IMF. For HII-dominated regions and starburst galaxies, Calzetti et al. (2007) propose:

$$SFR_{mix,C07} = 5.45 \times 10^{-42} [L(H\alpha)_{obs} + 0.031L(24)] \quad 3 \times 10^{38} \lesssim L(24) \lesssim 3 \times 10^{44}. \quad (16)$$

In both equations,  $L(H\alpha)_{obs}$  is the observed (non-extinction-corrected) luminosity at  $H\alpha$ . The difference between the two equations has been suggested to be due to a difference in the underlying stellar population heating the dust: continuous star formation no older

than about 100 Myr for starburst galaxies (and instantaneous bursts only a few Myr old for HII regions) and much longer lived star formation, of–order 10 Gyr, for whole galaxies (Kennicutt et al. 2009).

Our sample spans about 6 orders of magnitude in SFRs, from  $\sim 10^{-4} M_{\odot} \text{ yr}^{-1}$  all the way up to LIRGs producing stars at a pace close to a hundred  $M_{\odot} \text{ yr}^{-1}$ . The stellar populations dominating the dust heating will thus transition from those characteristic of normal star forming spirals or irregulars (with timescales for star formation of–order 10 Gyr) to those characteristic of young starbursts (timescales of order 100 Myr or less). In particular, we expect that as we approach the luminosity of the LIRGs, and consequently high dust opacities in the galaxies, the inferred SFRs will need to coincide with those derived from SFR(24). In Figure 8, we compare both equations 15 and 16 among themselves and with the SFR(24) calibration of Rieke et al. (2009), for the galaxies in our sample (excluding LIRGs). The difference between equations 15 and 16 is typically less than  $\sim 15\text{--}20\%$  up to  $L(24) \sim 4 \times 10^{42} \text{ erg s}^{-1}$ , and becomes larger for luminosities above that value. However, this is also approaching the luminosity range where  $\text{SFR}_{mix}$  needs to reach, as asymptotic value, SFR(24), as galaxies become more and more dominated by infrared emission. We elect to use equation 15 to derive SFRs for galaxies with luminosity  $L(24) \lesssim 4 \times 10^{42} \text{ erg s}^{-1}$ , and equation 16 for higher luminosity; this ‘transition’ luminosity corresponds to  $\text{SFR} \sim 1 M_{\odot} \text{ yr}^{-1}$ , which is a transition point where SFRs are equally contributed by the  $\text{H}\alpha$  and  $24 \mu\text{m}$  luminosities (Figure 7). In our sample, that transition SFR corresponds to  $\Sigma_{SFR} \approx 0.05 M_{\odot} \text{ yr}^{-1} \text{ kpc}^{-2}$ , albeit with a large uncertainty, since the relation between SFR and  $\Sigma_{SFR}$  has over one order of magnitude scatter. The exact value of the transition point is not crucial, as the difference between SFR values (equations 15 and 16) around that point are  $\sim 25\%$  ( $\sim 0.1$  dex in logarithmic scale), and physical systems are expected to transition from star formation timescales of many Gyrs down to a few hundred Myrs in a smoother way than what we describe here.

Equation 16 is almost identical to the linear part of Equation 8 for  $L(\text{H}\alpha)_{obs} \ll 0.031 L(24)$ . The closeness of the two equations becomes even more obvious if we recall that we measure a mean ratio  $L(\text{H}\alpha)/L(24) = 0.031$  for the metal–rich regions/galaxies in the sample of Calzetti et al. (2007), rather than 0.038, as measured by Rieke et al. (2009) in the same luminosity range  $10^7 L_{\odot} \leq L(24) \leq 10^{10} L_{\odot}$ . The difference is small, about 20%. The ratio  $L(\text{H}\alpha)/L(24) = 0.031$ , when combined with the median ratio  $L(24)/L(\text{TIR}) \sim 0.16$  (see Figure 3, and also Rieke et al. 2009) for the LIRGs, yields a mean value for  $L(\text{H}\alpha)/L(\text{TIR}) \sim 0.005$ , consistent with expectations for a 100 Myr old constant star formation population (from the models of Leitherer et al. 1999, where we assume  $L(\text{TIR}) \sim L(\text{bol})$ , the bolometric luminosity). It is worth noticing that the ratio  $L(\text{H}\alpha)/L(24) = 0.038$  determined by Rieke et al. (2009) yields values for  $L(\text{H}\alpha)/L(\text{TIR}) \sim 0.006$ , consistent with a 30 Myr old constant star formation population. Both calibrations are, thus, consistent with infrared emission from

galaxies whose emission is dominated by the young stellar population component, as expected for starbursts and LIRGs (e.g., Scoville et al. 2000, and references therein). We use our own measured value of the  $L(\text{H}\alpha)/L(24)$  mean ratio, and we replace it in equation 8, to derive SFRs in the high luminosity range of our sample.

The lowest luminosity LIRG in our sample has  $L(24)\sim 4\times 10^{43}$  erg  $\text{s}^{-1}$ , and the relative contributions of  $L(24)$  and  $L(\text{H}\alpha)_{obs}$  to the SFR now weight in favor of  $L(24)$ , with  $L(\text{H}\alpha)_{obs}/(0.031 L(24))<0.14$ . Thus, for our purposes, we can ignore the contribution of  $L(\text{H}\alpha)_{obs}$  to the SFR at the high luminosity end, and we will use the re-normalized equation 8 to measure SFRs for the LIRGs in our sample. Of the two expressions in equation 8, we only use the non-linear part, in agreement with the prescription of Rieke et al. (2009) for our range of luminosities. The non-linear correction in equation 8 attempts to correct for the presence of self-absorption in the LIRGs and brighter galaxies (Rieke et al. 2009). Derivations of  $\text{SFR}(24)$  for whole galaxies by other authors do not include such non-linear term although their luminosity range often includes LIRGs (equations 6, 7, 9, 10, and 12); the presence of a non-linear term can remain unnoticed if optical or infrared hydrogen recombination lines ( $\text{H}\alpha$ ,  $\text{P}\alpha$ ) are used as reference SFR indicators, since, even when extinction corrected, the luminosity of those lines will be underestimated in the presence of high dust opacity (as typical of LIRGs or brighter galaxies:  $L(24)>a\text{ few }10^{43}$  erg  $\text{s}^{-1}$ ). We note that in our sample the highest luminosity LIRG ( $L(24)\sim 8.3\times 10^{44}$  erg  $\text{s}^{-1}$ ) shows a difference between the SFRs derived from the re-normalized equation 8 and its linear version, equation 16, by less than 15%; for 27 of 29 LIRGs, this difference is less than 10%.

In summary, we derive SFRs for our sample of galaxies as follows:

$$\begin{aligned}
 \text{SFR} &= 5.45 \times 10^{-42} [L(\text{H}\alpha)_{obs} + 0.020L(24)], & L(24) < 4 \times 10^{42}, \\
 &= 5.45 \times 10^{-42} [L(\text{H}\alpha)_{obs} + 0.031L(24)], & 4 \times 10^{42} \leq L(24) < 5 \times 10^{43}, \\
 &= 1.70 \times 10^{-43} L(24) \times [2.03 \times 10^{-44} L(24)]^{0.048} & L(24) \geq 5 \times 10^{43}, \quad (17)
 \end{aligned}$$

where the SFR is in units of  $M_{\odot} \text{ yr}^{-1}$ , and luminosities are in units of erg  $\text{s}^{-1}$ . The transition regions between the three regimes are not sharp, and especially the one between equation 15 and 16 may begin at lower luminosity than the one used here. However, differences between the average SFRs at the transition points are about or less than 0.1 dex, which is smaller than both our typical error bar and the spread in the data points.

Equation 17 attempts to capture a general behavior of galaxies in regard to their dust opacity. Below  $L(24)\sim 4.0\times 10^{42}$  erg  $\text{s}^{-1}$  ( $L(\text{TIR})\approx 10^{10} L_{\odot}$ ), the young, star-forming regions in galaxies become increasingly transparent at optical and UV wavelengths for decreasing infrared luminosity (Buat et al. 2007, e.g.), and optical emission lines are useful tools for measuring their dust opacity (Kennicutt et al. 2009). At higher infrared lu-

minosities the star-forming regions become on average increasingly opaque, and, above  $L(24) \sim 5 \times 10^{43}$  ( $L(\text{TIR}) \approx 10^{11} L_{\odot}$ ), their infrared emission becomes the most reliable SFR indicator (Rieke et al. 2009, e.g.). In between  $L(\text{TIR}) \approx 10^{10} L_{\odot}$  and  $L(\text{TIR}) \approx 10^{11} L_{\odot}$  is a transition regime, where the SFR traced by optical and infrared line emission (e.g.,  $P\alpha$ ) is still non-negligible (Calzetti et al. 2007), although it becomes vanishingly small as the upper luminosity limit is reached. This argument should be interpreted as ‘mean behavior’, as galaxies generally contain large numbers of star-forming regions, with different levels of dust opacity. For instance, many galaxies with  $L(24) > 5 \times 10^{43}$  still show measurable emission at blue wavelengths, specifically  $H\beta(\lambda 4861 \text{ \AA})$  (Kennicutt et al. 2009).

For the 10 starburst galaxies with  $P\alpha$  data, we use this emission line, after correction for dust extinction, to measure the SFRs, adopting a ratio  $H\alpha/P\alpha = 7.82$  for the high-metallicity galaxies and  $H\alpha/P\alpha = 8.73$  for the low-metallicity ones (Calzetti et al. 2007).

## 5. Calibration of SFR(70)

In order to establish whether the  $70 \mu\text{m}$  luminosity can be used as a SFR indicator, and to check its limitations, we first concentrate on the 142 galaxies that constitute the high-metallicity sample. We will discuss the effect of metallicity on our calibrations later in this section.

The  $70 \mu\text{m}$  LSD tightly correlates with the SFSD (Figure 9) over the full 5 orders of magnitude spanned by our high-metallicity sample. The best fit through the data points is:

$$\text{Log}[\Sigma_{70}(\text{erg s}^{-1} \text{ kpc}^{-2})] = (1.089 \pm 0.013)\text{Log}[\Sigma_{SFR}(M_{\odot} \text{ yr}^{-1} \text{ kpc}^{-2})] + (43.303 \pm 0.029), \quad (18)$$

with a  $1 \sigma$  dispersion of the data around the best fit line of  $(-0.26, +0.22)$  dex. Although the best fit is non-linear, we need to recall that data with lower SFSD often correspond to galaxies with star formation scattered across the disk. These systems generally tend to be less dust opaque than systems with large and concentrated SFRs (Figures 4 and 5, and, e.g., Wang & Heckman 1996; Calzetti et al. 2007). Thus, the  $70 \mu\text{m}$  emission becomes ‘underluminous’ because the dust emission does not fully trace the light from star formation. Furthermore, the peak of the infrared emission shifts to longer wavelengths for decreasing IRSD values (Figure 5, and Draine et al. 2007), exacerbating the trend. If we constrain the fitting range to increasingly large SFSDs, the slope in equation 18 tends asymptotically to unity; it is about  $2 \sigma$  away from unity for  $\text{Log}(\Sigma_{SFR}) > -1$ , and less than  $1 \sigma$  away from unity for  $\text{Log}(\Sigma_{SFR}) > -0.6$ .

An independent test is to use the luminosities or SFRs normalized by the  $3.6 \mu\text{m}$  lumi-

osity,  $L(3.6)$  (Figure 10), since this proxy for stellar mass is not dominated by the current star formation. The best fit line through the data is:

$$\text{Log}[L(70)/L(3.6)] = (1.05 \pm 0.02)\text{Log}[SFR/L(3.6)] + (45.30 \pm 0.02), \quad (19)$$

where SFR is from equation 17, and the units of  $SFR/L(3.6)$  are  $M_{\odot} \text{ yr}^{-1} (\text{erg s}^{-1})^{-1}$ . The slope in equation 19 is also slightly above unity, only  $2.5 \sigma$  away from it, and quickly converges to unity for a fitting range that excludes the faintest galaxies, in agreement with our conclusions above.

The simple stellar population and dust emission models briefly presented in section 4.1 can further help clarify some of the observed trends. At the high luminosity end, the models have slope roughly unity by construction (Figure 5), because in this regime TIR is a good proxy for SFR; thus the trend  $\Sigma_{70} - \Sigma_{SFR}$  does not provide independent information. For  $SFSD \gtrsim 0.5 M_{\odot} \text{ yr}^{-1} \text{ kpc}^{-2}$ , the opacity of the system is high and most of the stellar light is reprocessed by dust into the infrared. Thus, the details of the extinction model are unimportant, and the normalization factor between the SFSD and the total infrared emission is provided by the stellar population. The mean trend of the  $70 \mu\text{m}$  surface brightness is well represented by a model where a 100 Myr constant star formation population heats the dust (Figure 11), which is how we describe LIRGs (equation 17).

As the SFSD decreases, the average opacity of the galaxy also tends to decrease (Wang & Heckman 1996; Heckman et al. 1998; Hopkins et al. 2001; Calzetti 2001). The opacity models of Calzetti et al. (2007, constructed for HII regions), however, under-predict the observed  $\Sigma_{70}$  by a factor  $\approx 2-3$ , for  $\Sigma_{SFR} \lesssim 0.01 M_{\odot} \text{ yr}^{-1} \text{ kpc}^{-2}$  (Figure 11). All the other models, the extreme case of constant and large dust opacity for the entire SFSD range, the model with a filling factor of the emitting regions of about 10% of the normalizing area, and the stellar population models with exponentially decreasing star formation, mark a range that brackets the spread of the datapoints. The mean trend of the data at low SFSDs is better represented by the model with a 10% filling factor of the emitting regions, for constant star formation. This is perhaps not surprising, since star formation in low surface brightness galaxies tends to display a morphology of isolated knots distributed over a large, non-contiguous area, while our definition of ‘ionized gas emitting area’ (section 2) measures a single size per galaxy, irrespective of the distribution and filling factor of the emitting regions. In the case of exponentially decreasing star formation, the mean trend of the data requires higher filling factors, around 30%–50% for the model with e-folding time  $\tau = 2 \text{ Gyr}$ , but still around 15% for the 5 Gyr e-folding time model.

Our solutions are not unique, and different combinations of dust opacity models and filling factors can provide similar trends. For instance, the extreme case of a foreground dust geometry for the galaxy, with a Small Magellanic Cloud extinction curve (Bouchet et al.



1985) and common dust column density to both stars and ionized gas gives higher values of the infrared emission at low luminosities than the starburst attenuation curve (Figure 11, right). Thus, the low luminosity data points can be accounted for by a larger filling factor of the emitting regions,  $\sim 30\%$ , about a factor 3 larger than the case of the starburst attenuation curve with constant star formation. The conclusion is still that a low filling factor is needed for the emitting regions, between 10% and 30%, for constant star formation, and around 15%–50% for exponentially decreasing star formation, in order to account for the observed trend between the 70  $\mu\text{m}$  emission and the SFSD for  $\Sigma_{SFR} \lesssim 0.01 M_{\odot} \text{ yr}^{-1} \text{ kpc}^{-2}$ .

The overall trend of the high–metallicity datapoints is consistent with both increasing dust opacity and a smooth transition from a low filling factor to a high filling factor as the SFSD increases, with the transition point located around  $\Sigma_{SFR} \sim 0.01 M_{\odot} \text{ yr}^{-1} \text{ kpc}^{-2}$ .

Including the low–metallicity data in the  $\Sigma_{70}$ – $\Sigma_{SFR}$  plot provides further insights into the effect of dust opacity, as low metallicity galaxies tend to be more transparent than higher metallicity ones (Cannon et al. 2005; Walter et al. 2007). These data are located typically below the trend marked by the high metallicity data (Figure 12), as expected for decreasing dust amounts in the ISM. The inclusion of the model line for a 1/10 solar metallicity ISM from Calzetti et al. (2007) shows that metallicity has a major effect on the location of galaxies on the plot. The model line for low–metallicity systems is located between 4 times (at high luminosity) and an order of magnitude (at low luminosity) below the analogous model for solar metallicity systems. There are a few galaxies that are located even below our low–metallicity model line (Figure 12). Two galaxies, Tol 65 and SBS 0335–052, are more than  $3 \sigma$  below the low–metallicity curve, while IZw18 is about  $2 \sigma$  below the curve. These three galaxies, with metallicities between 7.20 and 7.45 (Nagao, Maiolino & Marconi 2006), are among the most metal–poor in our sample, and well below 1/10th solar. Furthermore, Tol 65 and SBS 0335–052 have high effective dust temperature and their peak infrared emission is located shortward of 70  $\mu\text{m}$  (Engelbracht et al. 2008). Thus, it is not surprising that the three galaxies lie below our 1/10th solar model line.

Although metallicity has an important effect on the applicability of the 70  $\mu\text{m}$  emission as a SFR indicator, luminous infrared galaxies tend to be, on average, metal and dust rich. Indeed, most of the low metallicity galaxies in our sample cluster at the low–luminosity and low  $\Sigma_{SFR}$  end of the range (Figure 12), an effect of typically ‘hotter’ IR SEDs relative to the high metallicity galaxies (Dale et al. 2005). From the discussion above, we approximate the trend of the high metallicity data (Figure 13) with:

$$\Sigma_{SFR} (M_{\odot} \text{ yr}^{-1} \text{ kpc}^{-2}) = \frac{\Sigma_{70} (\text{erg s}^{-1} \text{ kpc}^{-2})}{1.7 \times 10^{43}}, \quad (20)$$

for  $7 \times 10^{40} \text{ erg s}^{-1} \text{ kpc}^{-2} \lesssim \Sigma_{70} \lesssim 5 \times 10^{44} \text{ erg s}^{-1} \text{ kpc}^{-2}$ , and:

$$\Sigma_{SFR} (M_{\odot} \text{ yr}^{-1} \text{ kpc}^{-2}) = \left( \frac{\Sigma_{70} (\text{erg s}^{-1} \text{ kpc}^{-2})}{3.3 \times 10^{44}} \right)^{0.65}, \quad (21)$$

for  $6 \times 10^{38} \text{ erg s}^{-1} \text{ kpc}^{-2} \lesssim \Sigma_{70} \lesssim 7 \times 10^{40} \text{ erg s}^{-1} \text{ kpc}^{-2}$ . The non-linear relation between the 70  $\mu\text{m}$  LSD and the SFSD (equation 21) attempts to take into account the increasing transparency of the medium for decreasing SFSDs, and is a rough approximation to a more complicated trend (Figure 13). The important caveat about equation 21 is that the calibration is highly dependent on the type of objects from which it is derived (whole galaxies in this case), and should not be applied to different physical systems (e.g., HII regions or complexes within galaxies) or metal-poor systems, without proper checks. Nevertheless, the two power laws, equations 20 and 21, provide a reasonable description of the datapoints in our sample with  $12 + \text{Log}(\text{O}/\text{H}) \gtrsim 8.1$ , which show a relatively symmetric dispersion around those mean trends, with  $1 \sigma \sim 0.2 \text{ dex}$  (Figure 14), or about 60% in linear scale. The linear relation at the high luminosity end offers the opportunity to derive directly a star formation rate:

$$SFR(70) (M_{\odot} \text{ yr}^{-1}) = \frac{L(70) (\text{erg s}^{-1})}{1.7 \times 10^{43}}, \quad (22)$$

applicable to luminosities  $L(70) \gtrsim 1.4 \times 10^{42} \text{ erg s}^{-1} \sim 3.7 \times 10^8 L_{\odot}$ . Again, the relation between luminosity and luminosity/area has at least one order of magnitude dispersion in our sample. Adopting the SFR–L(H $\alpha$ ) calibration of Kennicutt (1998) would change the denominator in equation 22 to  $1.2 \times 10^{43}$ . The change is due to a combination of adopted stellar IMF (Kennicutt 1998, adopts a Salpeter IMF in the range 0.1–100  $M_{\odot}$ ) and small ( $\sim 10\%$ ) variations in the stellar population models.

## 6. Can we also derive SFR(160)?

The 160  $\mu\text{m}$  luminosity surface density correlates, like  $\Sigma_{70}$ , with SFSD, with linear best-fit (Figure 15):

$$\text{Log}[\Sigma_{160} (\text{erg s}^{-1} \text{ kpc}^{-2})] = (0.954 \pm 0.012) \text{Log}[\Sigma_{SFR} (M_{\odot} \text{ yr}^{-1} \text{ kpc}^{-2})] + (42.910 \pm 0.028). \quad (23)$$

The shallower-than-unity slope is consistent with our interpretation that galaxies of decreasing luminosity have ‘cooler’ infrared SEDs (Figure 5). This result is partially counteracted by a sharp decline in 160  $\mu\text{m}$  luminosity for  $\text{Log}(\Sigma_{SFR}) < -3.2$ , indicating the regime where the galaxies become measurably transparent. Like in the case of the 70  $\mu\text{m}$  LSD, the best fit line tends asymptotically to a slope of unity for galaxies of increasing SFSD; in the case of the 160  $\mu\text{m}$  LSD, the best fit has slope of unity within  $1 \sigma$  already for  $\text{Log}(\Sigma_{SFR}) > -2$ .

The dispersion of the data about the best fit line,  $(-0.42, +0.34)$  dex, is larger for  $\Sigma_{160}$  than for  $\Sigma_{70}$  by a factor  $\sim 1.3$ – $1.5$  (Figures 9 and 14, left). A similar ratio,  $\sim 1.2$ – $1.4$ , between the  $1\sigma$  dispersions holds when quantities normalized to the  $3.6\ \mu\text{m}$  luminosity are considered (Figure 16). Furthermore, the dispersion on the positive side, i.e., for large  $160\ \mu\text{m}$  emission relative to the SFSD, is dominated by galaxies at the low SFSD end,  $\text{Log}(\Sigma_{SFR}) < -2$  (Figure 15, right). These two facts, when taken together, suggest that the  $160\ \mu\text{m}$  luminosity is affected by emission from dust heated by stellar populations unrelated to those powering the current star formation. This effect becomes more and more evident as total luminosities decrease. Indeed, the same models with 10% area filling factor of the emitting regions that satisfactorily account for the trend of the data in the  $\Sigma_{70}$ – $\Sigma_{SFR}$  plane provide an unsatisfactory description of the data in the  $\Sigma_{160}$ – $\Sigma_{SFR}$  plane for  $\Sigma_{SFR} < 0.01$ , for both the 100 Myr and the 10 Gyr constant star formation population: the models underpredict, by roughly 0.5 dex or more, the median locus of the data for most of the data in this SFSD range (Figure 15, left). A similar result is obtained when normalizing the  $160\ \mu\text{m}$  luminosity and the SFR to the  $3.6\ \mu\text{m}$  luminosity, instead of the area (Figure 16). In this case, the best fit to the data has a shallower-than-unity slope:

$$\text{Log}[L(160)/L(3.6)] = (0.87 \pm 0.01)\text{Log}[SFR/L(3.6)] + (37.48 \pm 0.02), \quad (24)$$

(SFR is from equation 17, and the units of  $SFR/L(3.6)$  are  $\text{M}_\odot \text{yr}^{-1} (\text{erg s}^{-1})^{-1}$ ) and the slope approaches unity only for increasingly luminous galaxies. Equation 24 supports, again, the presence of two effects for low luminosity galaxies: ‘cooler’ infrared SEDs for lower luminosity galaxies and a proportionally increasing contribution, for decreasing luminosity, to the  $160\ \mu\text{m}$  emission from dust heated by stellar populations that are unassociated with the current star formation.

All of the above strengthens the consideration made in section 5 that a fraction of the dust emission detected in the  $70\ \mu\text{m}$  band towards low SFSDs is due to stellar populations unrelated to the current star formation, and this effect is, expectedly, larger in the  $160\ \mu\text{m}$  showing up as both a larger dispersion of the data around the mean trends and an excess deviation of the data from the model predictions at the low luminosity end.

This conclusion is not new: many authors have noted over the past  $\sim 2$  decades the presence of a contribution to the infrared emission of galaxies from stellar populations not associated to the current star formation (e.g. Lonsdale Persson & Helou 1987; Sauvage & Thuan 1992; Buat & Xu 1996; Walterbos & Greenawalt 1996). The presence of such contribution limits the usefulness, especially at low galaxy luminosity, of SFR calibrations based on monochromatic IR emission long-ward of the SED peak.

In summary, although defining a SFR calibration based on the  $160\ \mu\text{m}$  emission may

be tempting:

$$SFR(160) (M_{\odot} yr^{-1}) = \frac{L(160) (erg s^{-1})}{7.0 \times 10^{42}}, \quad (25)$$

with applicability for  $L(160) > 1 \times 10^{42} \text{ erg s}^{-1} \sim 2.6 \times 10^8 L_{\odot}$ , the presence of a large dispersion,  $1 \sigma_{160} \sim 2 \times (1 \sigma_{70})$ , and the likelihood that a portion of the 160  $\mu\text{m}$  emission (and a large fraction of it at the low luminosity end) is due to dust heating by non–star–forming populations prevents such definition from having the same confidence level as the calibration of SFR(70). As will be seen in the next section, a safer range of applicability of equation 25 is for  $L(160) > 2 \times 10^{42} \text{ erg s}^{-1} \sim 5.2 \times 10^8 L_{\odot}$ . With the Kennicutt (1998) calibration for SFR, the numerical constant in equation 25 changes to  $4.8 \times 10^{42}$ .

## 7. Discussion and Conclusions

Monochromatic SFR indicators have the undeniable advantage of immediate application, especially when observational data cover only a limited range of wavelengths, as is often the case for distant galaxies. Within this framework, Spitzer observations of nearby galaxies are offering a unique opportunity to test the applicability and limitations of a number of SFR indicators based on the emission from the dust heated by stars. The main goal of this paper has been to investigate whether long wavelength infrared fluxes at 70 and 160  $\mu\text{m}$  can be reasonably calibrated as SFR indicators.

Previous works (Roussel et al. 2001; Boselli, Lequeux & Gavazzi 2004; Förster Schreiber et al. 2004; Calzetti et al. 2005; Dale et al. 2005; Wu et al. 2005; Alonso–Herrero et al. 2006; Perez–Gonzalez et al. 2006; Calzetti et al. 2007; Dale et al. 2007; Relaño et al. 2007; Zhu et al. 2008; Rieke et al. 2009) have concentrated on the short–wavelength infrared, in the range  $\sim 5\text{--}40 \mu\text{m}$ , commonly referred to as the mid–IR. Although this wavelength range is below the peak infrared emission and only accounts for a few to a few tens of percent of the total IR emission, the dust heated by the hot, massive stars in a recent star formation episode can have high effective temperatures and may preferentially emit at the shorter infrared wavelengths. This general picture translates into a correlation between SFRs and luminosities in the mid–IR, such as L(8) and L(24), centered on the Spitzer 8  $\mu\text{m}$  and 24  $\mu\text{m}$  bands. However, the correlation is also non–linear, and careful analysis shows that other contributors, in addition to the SFR, determine the short infrared wavelength luminosity, including transiently heated dust by single UV or optical photons in the general radiation field of a galaxy (Haas, Klaas & Bianchi 2002; Boselli, Lequeux & Gavazzi 2004; Bendo et al. 2008), variations in the dust effective temperature (for L(24), see Calzetti et al. 2007), and, especially for L(8), gas metallicity (Engelbracht et al. 2005; Rosenberg et al. 2006; Wu et al. 2006; Madden et al. 2006; Jackson et al. 2006; Draine et al. 2007).

### 7.1. Establishing a reference SFR indicator

In the present paper, we first establish a “reference” SFR calibration against which we can compare the 70 and 160  $\mu\text{m}$  luminosities and LSDs. The adopted reference SFR is given in equation 17. We choose this particular three-part formulation after comparing a variety of published 24  $\mu\text{m}$  and composite  $\text{H}\alpha$  plus 24  $\mu\text{m}$  calibrations.

Above a luminosity  $L(24)\sim 5\times 10^{43}$   $\text{erg s}^{-1}$ , the contribution of direct (unabsorbed by dust) massive stars emission to SFR measurements is less than  $\sim 15\%$  (Figure 8). For  $L(24)\gtrsim 5\times 10^{43}$   $\text{erg s}^{-1}$ , the 24  $\mu\text{m}$  emission is a reliable SFR indicator *for whole galaxies*. However, in this luminosity regime, the 24  $\mu\text{m}$  dust emission is also self-absorbed (Rieke et al. 2009), and a non-linear correction to the calibration needs to be applied (third term of equation 17, re-calibrated from the original calibration of Rieke et al. 2009).

Below a luminosity  $L(24)\sim 5\times 10^{43}$   $\text{erg s}^{-1}$ , the galaxies become transparent in the mid-infrared, and, as the luminosity further decreases, they become transparent at UV and optical wavelengths (e.g. Bell 2003; Buat et al. 2007). At  $L(24)\sim 4\times 10^{42}$   $\text{erg s}^{-1}$ , the contribution to the SFR from emission at optical ( $\text{H}\alpha$ ) and IR (24  $\mu\text{m}$ ) wavelengths is roughly equal (Figure 7). Thus, below  $L(24)\sim 5\times 10^{43}$   $\text{erg s}^{-1}$ , the 24  $\mu\text{m}$  luminosity becomes increasingly insufficient, by itself, to fully characterize the SFR of a whole galaxy, as the portion of the star formation unabsorbed by dust becomes an increasingly significant contribution to the total SFR. Additionally, a correlation between the effective temperature and the luminosity/area of the thermal dust emission is present in HII regions for decreasing total SFRs (Calzetti et al. 2007), and in whole galaxies as well (Chanial et al. 2007, and Figure 4). This correlation is further seen for the galaxies in our sample by plotting the ratio  $L(70)/L(24)$  as a function of  $\Sigma_{24}$  (Figure 17, left); the star-forming and starburst galaxies show a decreasing IR ratio for increasing LSD, as is expected for increasing effective dust temperature. This trend is less tight, albeit still present to some extent, when the  $L(70)/L(24)$  ratio is plotted as a function of luminosity (Figure 17, right), rather than luminosity/area. The use of total luminosity introduces a scatter because galaxies are weighted not only by their intrinsic level of activity but also by their size (a low-activity, low IR/area, but large, galaxy can have the same total luminosity as a high-activity, but small, galaxy).

The relation between the effective temperature and the luminosity/area is present in galaxies up to an IRSD  $\Sigma_{TIR}\sim 0.3\text{--}1\times 10^{43}$   $\text{erg s}^{-1}$   $\text{kpc}^{-2}$  ( $L(\text{TIR})\approx 1\text{--}3\times 10^{44}$   $\text{erg s}^{-1}\sim 3\text{--}8\times 10^{10}$   $L_{\odot}$  or  $L(24)\approx 2\text{--}6\times 10^{43}$   $\text{erg s}^{-1}\sim 0.6\text{--}2\times 10^{10}$   $L_{\odot}$ ), independent of the metal content of the galaxy (Figures 3, left, and 17). This can be further seen in Figure 18, where we plot  $L(24)/L(\text{TIR})$  as a function of IRSD only for those galaxies in our sample that have oxygen abundance  $12+\text{Log}(\text{O}/\text{H})>8.5$ . There are 66 normal star-forming and starburst galaxies above this oxygen abundance value (or about 1/2 of the normal star-forming and

starburst galaxies in our high-metallicity sample), and the 29 LIRGS. A similar trend to that of Figure 4 of  $L(24)/L(\text{TIR})$  as a function of IRSD is observed for this restricted, higher-metallicity, sample, implying that the effect is not primarily driven by the abundance of dust in the galaxy.

The combination of the two effects discussed above, increasing transparency of the galaxy and decreasing effective dust temperature for decreasing luminosity/area, may account for why different published calibrations for  $\text{SFR}(24)$  diverge when extrapolated towards the low luminosity end of our sample (Figure 6).

A more robust approach appears to involve measuring SFRs by combining the  $24\ \mu\text{m}$  emission and the observed  $\text{H}\alpha$  emission, the first probing the dust-absorbed star formation and the second the unabsorbed one (Kennicutt et al. 2007; Calzetti et al. 2007; Kennicutt et al. 2009). Even in this case, though, a transition in the proportionality factor between  $L(24)$  and  $L(\text{H}\alpha)_{\text{obs}}$  needs to be implemented for estimating accurate SFRs (equation 17, first and second terms). The transition point is marked by the condition  $L(\text{H}\alpha)_{\text{obs}}/(0.020 L(24)) \approx 1$ , which corresponds, for decreasing luminosity, to the transition from galaxies dominated by dust-obscured star formation to those dominated by unobscured star formation (e.g., Buat et al. 2007). The transition corresponds to a change of roughly 50% in the fraction of  $L(24)$  to be added to  $L(\text{H}\alpha)_{\text{obs}}$ , from 0.031 to 0.020 for  $L(24) \lesssim 4 \times 10^{42}\ \text{erg s}^{-1}$  (first and second relation in equation 17), or  $\text{SFR} \sim 1\ M_{\odot}\ \text{yr}^{-1}$ . Our results thus indicate that around this SFR, the stellar population that dominates the heating of the dust transitions from being a relatively young one (constant star formation over the past  $\sim 100$  Myr), typical of starburst events, to a much more evolved one (constant star formation over the past  $\sim 10$  Gyr, or exponentially decreasing star formation), typical of widespread, but isolated, star formation amid an otherwise quiescent galaxy. In the latter case, the stellar populations not associated with the current star formation events contribute, in a proportionally larger fraction than the starburst case, to the heating of the  $24\ \mu\text{m}$ -emitting dust, whose emission contribution to the SFR calibration needs then to be accordingly removed, on average, by reducing  $L(24)$  via a smaller multiplicative factor (Kennicutt et al. 2009) than the starburst or HII region case (Calzetti et al. 2007).

## 7.2. The Emission at 70 and 160 $\mu\text{m}$ as SFR Indicators

Our reference SFR calibration (equation 17) enables the testing of other IR monochromatic luminosities as potential SFR indicators, over a wide range of luminosity surface densities. In this work, we have investigated the applicability of  $L(70)$  and  $L(160)$ , i.e., the long wavelength Spitzer bands, as such indicators. These two wavelength regions are attrac-

tive because each contains a larger fraction of the total infrared emission than the 24  $\mu\text{m}$  emission.  $L(70)$  is 50% of  $L(\text{TIR})$  for infrared LSD  $\Sigma_{\text{TIR}} \gtrsim 10^{41} \text{ erg s}^{-1} \text{ kpc}^{-2}$ , and  $L(160)$  is about 20% or larger fraction of  $L(\text{TIR})$  for the entire luminosity range investigated in this work (Figure 5), while  $L(24)$  only contains  $\sim 10\text{--}20\%$  of the total infrared emission (Figure 4). At the high luminosity LIRGs regime, both  $L(70)$  and  $L(160)$  represent a larger fraction of  $L(\text{TIR})$  than expected from models (e.g., Draine et al. 2007). This indicates a ‘cooler’ (larger effective temperature) IR SED for those systems than predicted from models, an effect already previously observed (Dunne & Eales 2001), and in line with the re-emission of the energy absorbed in the mid-IR wavelength range. Galaxies in the LIRG luminosity regime ( $L(\text{TIR}) \gtrsim 1\text{--}3 \times 10^{44} \text{ erg s}^{-1}$  or  $\Sigma_{\text{TIR}} \gtrsim 0.3\text{--}1 \times 10^{43} \text{ erg s}^{-1} \text{ kpc}^{-2}$ ) have their energy output dominated by a star formation event that is evolving in a dusty region of sufficient optical depth that even light emerging in the mid-IR gets absorbed, and that light is re-emitted at longer infrared wavelengths, thus producing an overall cool IR SED, specifically cooler than expected from models where the emission is emerging from a transparent (at IR wavelengths) medium. This requires that the dust attenuation at 24  $\mu\text{m}$  be substantial,  $\approx 1$  mag, which translates to  $A_V \sim 15\text{--}50$  mag, depending on the extinction curve adopted (Draine & Li 2007; Flaherty et al. 2007). Such large extinctions are not uncommon in bright infrared galaxies. Genzel et al. (1998) determines that Ultraluminous InfraRed Galaxies (ULIRGs) have extinction values in the range  $A_{V,\text{screen}} = 5\text{--}50$  mag or  $A_{V,\text{mixed}} = 50\text{--}1000$  mag. Although LIRGs are typically an order of magnitude less luminous than ULIRGs, dust extinctions are still expected to be large in these systems.

The 70  $\mu\text{m}$  emission from galaxies correlates linearly with the SFR for luminosities  $L(70) \gtrsim 1.4 \times 10^{42} \text{ erg s}^{-1}$ , and a calibration is given in equation 22. Similarly,  $L(160)$  is linearly correlated with the SFR for luminosities  $L(160) > 10^{42} \text{ erg s}^{-1}$  (equation 25).

### 7.3. Scatter and the Impact of Metallicity

Both calibrations, however, have significant scatter,  $\sigma_{70} \sim 0.2$  dex and  $\sigma_{160} \sim 0.4$  dex, around the mean trend, and larger than the scatter,  $\sim 0.12\text{--}0.16$  dex, of calibrations based on  $L(24)$  or a mix of  $L(24)$  and  $L(\text{H}\alpha)_{\text{obs}}$  (Rieke et al. 2009; Kennicutt et al. 2009). Furthermore, the scatter increases for increasing wavelength, and appears to decrease for increasing luminosity (Figures 9 and 15). To further investigate these trends, we subdivide the high-metallicity sample into two sub-samples of oxygen abundance higher and lower than  $12 + \text{Log}(\text{O}/\text{H}) = 8.5$ . This separating value is chosen after Figure 5 of Engelbracht et al. (2008), which shows that galaxies above a metallicity of  $12 + \text{Log}(\text{O}/\text{H}) = 8.5$  tend to have a roughly constant  $f_\nu(70)/f_\nu(160)$  flux ratio (albeit with a large scatter), while galaxies below

that metallicity value tend to have an increasing ratio for decreasing oxygen abundance. Figure 19 shows the high-metallicity data from our sample divided into the two sub-samples. The scatter of the data for both  $\Sigma_{70}$  and  $\Sigma_{160}$  at constant  $\Sigma_{SFR}$  decreases when the sample is restricted to the highest metallicity data ( $12+\text{Log}(\text{O}/\text{H})>8.5$ ); specifically, faint  $\Sigma_{70}$  and  $\Sigma_{160}$  galaxies at fixed  $\Sigma_{SFR}$  drop from the sub-sample. This is in the expected direction, as seen earlier in this paper when analyzing the impact of low-metallicity galaxies (Figure 12).

The results of Figure 19 should be considered preliminary at this stage, since the oxygen abundances for a number of high metallicity galaxies are uncertain (sometimes by about 0.2 dex). Nevertheless, they confirm and support the results of section 5 for SFR(70), in the sense that as the galaxy metallicity increases and its dispersion in the sample decreases,  $\Sigma_{70}$  tends to approach a 1-to-1 relation with  $\Sigma_{SFR}$  (Figure 19, left). Conversely, in the same metallicity range,  $\Sigma_{160}$  tends to show a flattening in its relation with  $\Sigma_{SFR}$  and the dispersion of the datapoints around the mean trend remains large (Figure 19, right). As already discussed in section 6, this is in agreement with the expectation that the emission at long infrared wavelengths receives an increasing contribution (and an increasing scatter) from dust heated by evolved stellar populations unassociated with the current star formation event (Lonsdale Persson & Helou 1987). Figure 19 (right) further stresses that the calibration of SFR(160), equation 25, should only be used for  $L(160)\gtrsim 2\times 10^{42}$  erg s<sup>-1</sup>, equivalent to  $\Sigma(160)\gtrsim 10^{41}$  erg s<sup>-1</sup> kpc<sup>-2</sup>.

The calibrations just quoted apply to galaxies with mean gas-phase metallicity  $12+\text{Log}(\text{O}/\text{H})\gtrsim 8.1$ . The scatter becomes even larger, and strongly asymmetric, when galaxies of lower oxygen abundance value are included (Figure 12). Indeed, the general expectation is that lower metallicity galaxies will in general contain less dust and be more transparent than high metallicity ones. In this case, the infrared will be a poor tracer of the current SFR, as most of the light produced by recently formed stars will emerge unabsorbed by dust. Thus, calibrations of SFR(70) and SFR(160) need to be applied with an awareness of these limitations.

A hybrid approach, which combines, e.g., L(70) with  $L(\text{H}\alpha)_{obs}$ , could likely provide a ‘remedy’ for the dependence of far-infrared SFR indicators on metallicity, similarly to what has been already verified for L(24) and L(8), where the H $\alpha$  observed luminosity accounts for the unabsorbed portion of the SFR (Calzetti et al. 2007; Kennicutt et al. 2007; Zhu et al. 2008; Kennicutt et al. 2009). However, our current data do not enable us to test such hybrid indicator. We already use, in our analysis, the combination of L(24) and  $L(\text{H}\alpha)_{obs}$  to trace SFR in galaxies in an unbiased fashion; this indicator is dominated by the H $\alpha$  luminosity at low SFSDs, as would be any hybrid combination of L(70) and  $L(\text{H}\alpha)_{obs}$ , thus resulting in a degeneracy for any test we may attempt.



#### 7.4. Limits of Validity for the Calibrations

For the high-metallicity galaxies, the presence of a contribution to the long wavelength infrared emission from dust heated by populations unassociated with the current star formation becomes more evident for decreasing galaxy luminosity (Figures 11 and 15), but it also competes with at least three more effects: (1) the effective dust temperature as measured from the IR SED tends to decrease (Figures 4, 5, and 17); (2) the overall interstellar medium becomes progressively more transparent, thus a lower fraction of the SFR is traced in the infrared; (3) the area filling factor of star forming regions within galaxies decreases down to  $\sim 10\%$ – $50\%$  (depending on the dust extinction and stellar population models adopted), contributing to a change in the proportionality between  $\Sigma_{70}$  or  $\Sigma_{160}$  with  $\Sigma_{SFR}$  relative to that of more luminous galaxies. These four effects all contribute, at various levels, to change the relation between  $\Sigma_{70}$  (or  $\Sigma_{160}$ ) with  $\Sigma_{SFR}$  from linear to non-linear (e.g., equation 21), thus complicating any SFR calibration for luminosities below  $L(70) \sim 1.4 \times 10^{42}$  erg s $^{-1}$  ( $L(160) \sim 2 \times 10^{42}$  erg s $^{-1}$ ). These luminosity values correspond to  $SFR \sim 0.1$ – $0.3$  M $_{\odot}$  yr $^{-1}$ , thus they are not extremely restrictive especially in the context of current studies of distant galaxy populations.

#### 7.5. Additional Uncertainties

So far, we have not considered another potential source of heating for the dust associated with the diffuse interstellar medium of galaxies: photons leaking out of HII regions. About 20%–50% of the integrated H $\alpha$  luminosity in galaxies is associated with the diffuse interstellar medium (Reynolds 1990; Ferguson et al. 1996; Wang, Heckman & Lehnert 1997; Martin 1997), and those photons will produce dust heating. Since photons leaking out of HII regions are associated with the current star formation, in our analysis we have assumed that the diffuse ionized gas follows the same scaling relations as the star-forming regions, albeit possibly with different absolute values (e.g., the diffuse photons may heat the dust to typically lower temperatures than those spatially associated with HII regions). This assumption is justified by the findings of Wang, Heckman & Lehnert (1998), according to which the fraction of diffuse-to-total ionized gas remains relatively constant in galaxies of increasing SFR, from normal star-forming to starbursts, once the diffuse gas surface brightness is normalized by the mean star formation rate per unit area of the galaxy.

The infrared emission from the galaxies in our sample is dominated by dust heated by current star formation, with, in some cases, a non-negligible contribution from evolved stellar populations. Central AGNs do not constitute a dominant heating source in the present sample, and their impact on our calibrations could not be investigated. Thus, the

application of equations 22 (or 25) to galaxy populations first requires the evaluation of any contamination from dust heating by AGNs.

Finally, all our derivations are predicated on the assumption that a universal stellar IMF can be applied to all galaxies. While this is generally considered a reasonable assumption, recent investigations using the Sloan Digital Sky Survey galaxy sample (Hoverstern & Glazebrook 2008), a complete sample of nearby galaxies (Lee et al. 2009b), and HI-selected galaxies (Meurer et al. 2009) suggest that the galaxy-integrated IMF may become increasingly bottom-heavy in low luminosity and/or low surface brightness galaxies. Changes in the IMF will clearly impact the calibration of SFR indicators across the full electromagnetic wavelength range, and establishing whether such variations are present and what physical parameters may be causing them is an avenue of future investigation.

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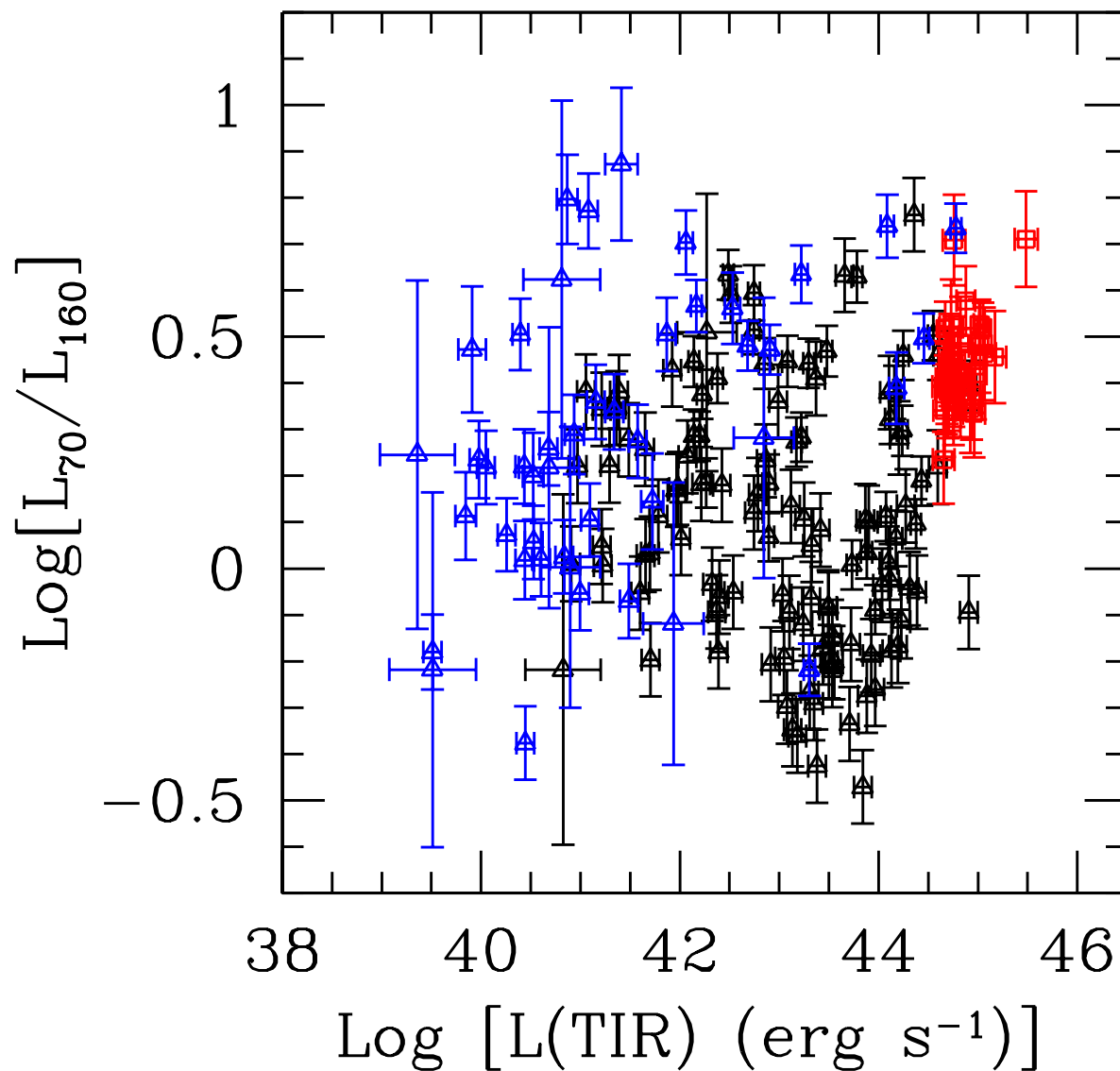


Fig. 1.— The logarithm of the ratio of 70-to-160  $\mu\text{m}$  luminosity for the 189 galaxies in our sample, as a function of the total infrared luminosity TIR. The sample covers almost 6 orders of magnitude in infrared luminosity. Symbol colors identify the star-forming and starburst high-metallicity (black symbols) and low-metallicity (blue symbols) galaxies, and the LIRGs (red).

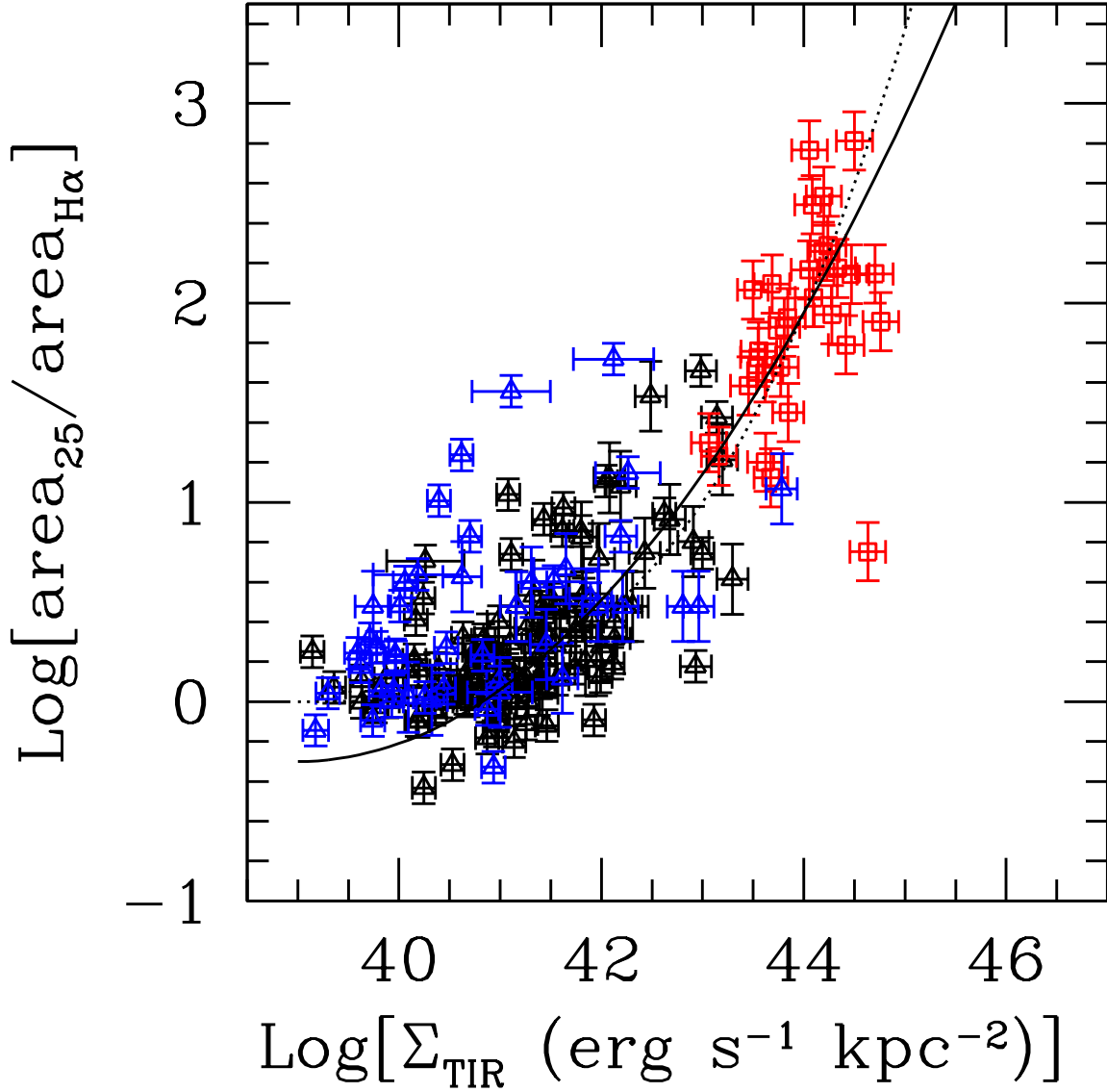


Fig. 2.— The ratio of the area defined by the  $R_{25}$  radius to the ionized gas ( $\text{H}\alpha$ ) emitting area as a function of the infrared luminosity surface density (IRSD:  $\Sigma_{\text{TIR}} = L(\text{TIR})/\text{area}$ ) for our sample galaxies, where the area used to normalize the infrared luminosity is  $\text{area}_{\text{H}\alpha}$ . Symbol colors are as in Figure 1. Two indicative mean trends are shown for the high metallicity datapoints; the trends are not fits to the data.

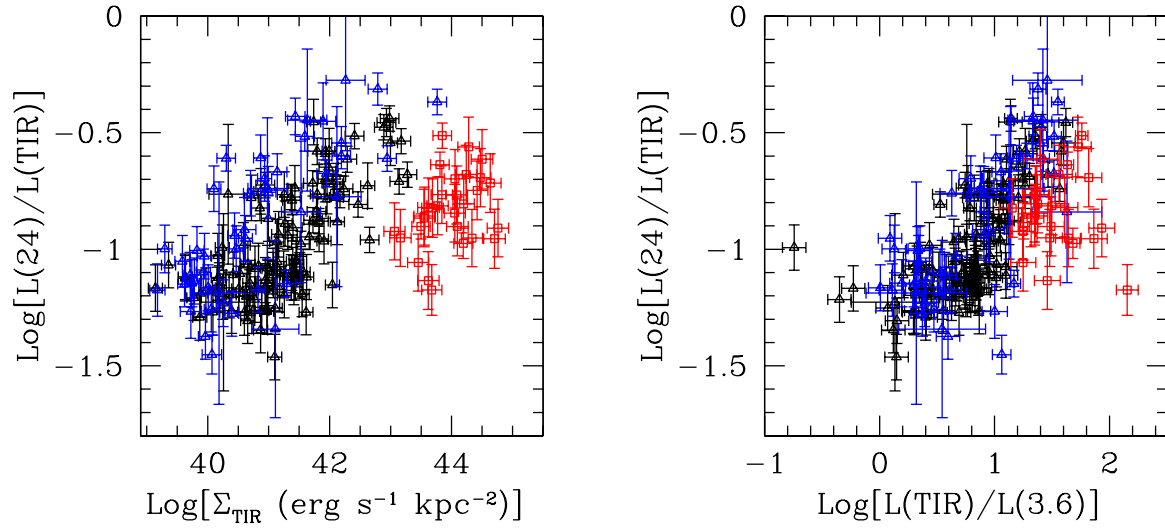


Fig. 3.— The  $L(24)/L(\text{TIR})$  ratio as a function of the IRSD (left) and of the infrared luminosity per unit stellar luminosity at  $3.6 \mu\text{m}$  (a proxy for mass in stars, right) for the galaxies in our sample. The color scheme for the datapoints is as in Figure 1.

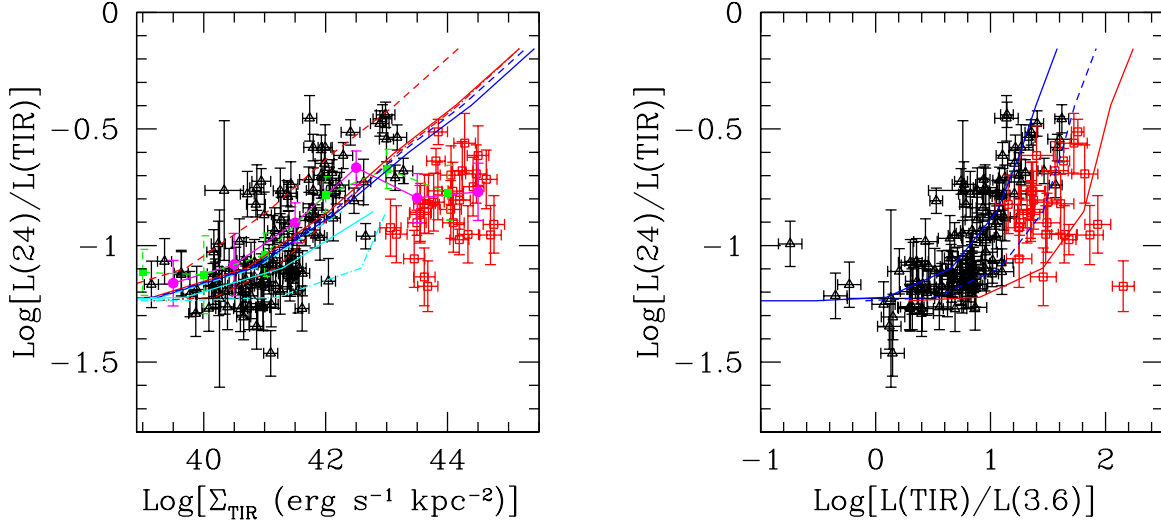


Fig. 4.— The same as Figure 3 for the high-metallicity galaxies, only. (Left) Mean values of the data binned in 1 dex intervals are shown in magenta and in green (with  $1\sigma$  error bars), corresponding to binning intervals shifted by 0.5 dex between the magenta and green points. The expected trends for a 100 Myr old, a 1 Gyr, and a 10 Gyr constant star formation population reddened by increasing amounts of foreground dust for larger SFRD, with a starburst attenuation curve (default dust model, see Calzetti et al. 2007), are marked by the red continuous, the blue dash, and the blue continuous lines, respectively. The case of exponentially decreasing star formation with e-folding times of  $\tau = 5$  Gyr (dashed cyan line) and 2 Gyr (continuous cyan line) are shown for the default dust model. Variations on the default dust model are also shown: a 100 Myr constant star formation population attenuated by constant, large dust opacity ( $A_V > 10$  mag) independent of TIR luminosity or SFRD (red dot-dash line); a 100 Myr constant star formation population attenuated by our default dust model, but with a 10% filling factor for the IR-emitting regions within the galaxy (red dashed line); and a  $\tau = 2$  Gyr exponentially decreasing star formation model, with the stellar population homogeneously mixed with the dust (cyan dot-dash line). (Right) Model lines for a 100 Myr old (red continuous line), a 1 Gyr old (blue dashed line) and a 10 Gyr old (blue continuous line) constant star formation population are overplotted on the data for  $L(24)/L(\text{TIR})$  versus the infrared emission per unit stellar luminosity measured at  $3.6 \mu\text{m}$ .

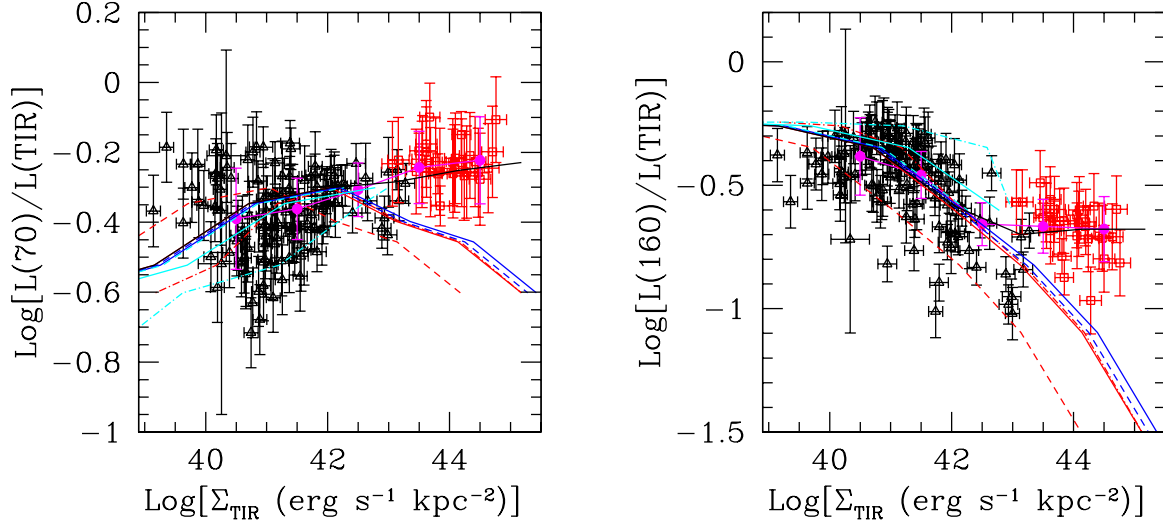


Fig. 5.— The 70  $\mu\text{m}$ -to-TIR and the 160  $\mu\text{m}$ -to-TIR ratios as a function of the IRSD for the metal-rich galaxies in our sample. As in Figure 4, the LIRGs are marked with red symbols. Mean values of the data binned in intervals of 1 dex each are shown in magenta. The model lines are the same as in Figure 4 (left). According to those models, the peak of the infrared emission moves *in and out* of the 70  $\mu\text{m}$  band, i.e., the IR SED becomes of sufficiently high effective temperature that the peak IR emission moves to shorter wavelengths than the MIPS 70  $\mu\text{m}$  band, for the range of luminosity surface densities (LSDs) spanned by our sample (Draine & Li 2007, see their Figure 15). For the same models, the 160  $\mu\text{m}$  emission becomes progressively fainter, as the IRSD increases. Modified models for both  $L(70)/L(\text{TIR})$  and  $L(160)/L(\text{TIR})$  that account for the observed properties of the luminous galaxies, and the LIRGs in particular, require those luminosity ratios to remain approximately constant beyond IRSD  $\Sigma_{\text{TIR}} \approx 10^{42-42.5} \text{ erg s}^{-1} \text{ kpc}^{-2}$  (black line).

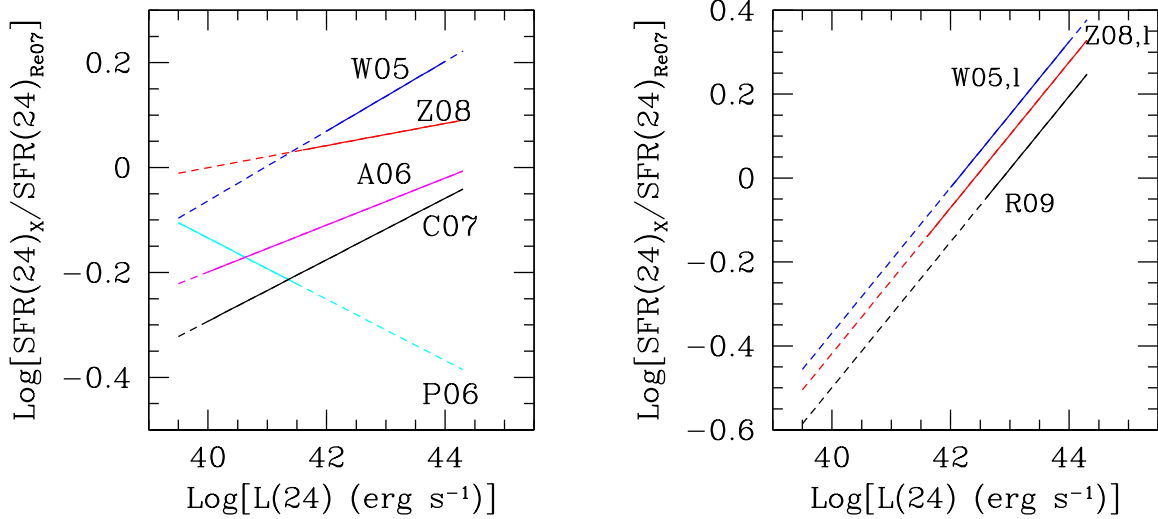


Fig. 6.— The ratio of SFR calibrations using the 24  $\mu\text{m}$  luminosity, as a function of the observed  $L(24)$  (in  $\text{erg s}^{-1}$ ), for the luminosity range spanned by our high-metallicity galaxies sample (excluding LIRGs),  $39.5 \lesssim \text{Log}[L(24)] \lesssim 44.3$ . **(Left)**. The ratio of the non-linear calibrations of Wu et al. (2005, W05, blue line), Alonso-Herrero et al. (2006, A06, magenta line), Perez-Gonzalez et al. (2006, P06, cyan line), Calzetti et al. (2007, C07, black line), Zhu et al. (2008, Z08, red line) to the non-linear calibration of Relaño et al. (2007). The continuous lines show the luminosity range used by the authors to derive their calibrations, while the dashed lines extend those calibrations to the entire luminosity range spanned by our high-metallicity sample. The non-linear calibration of Relaño et al. (2007) is used as reference because it has been derived for the entire luminosity range displayed in our Figure (see equation 14). **(Right)**. The ratio of the linear calibrations of Wu et al. (2005, Wu05,1, blue line), Zhu et al. (2008, Z08,1, red line), and Rieke et al. (2009, R09, black line) to the non-linear calibration of Relaño et al. (2007). As in the left panel, continuous and dashed lines indicate range of derivation and extrapolation of the calibrations, respectively.



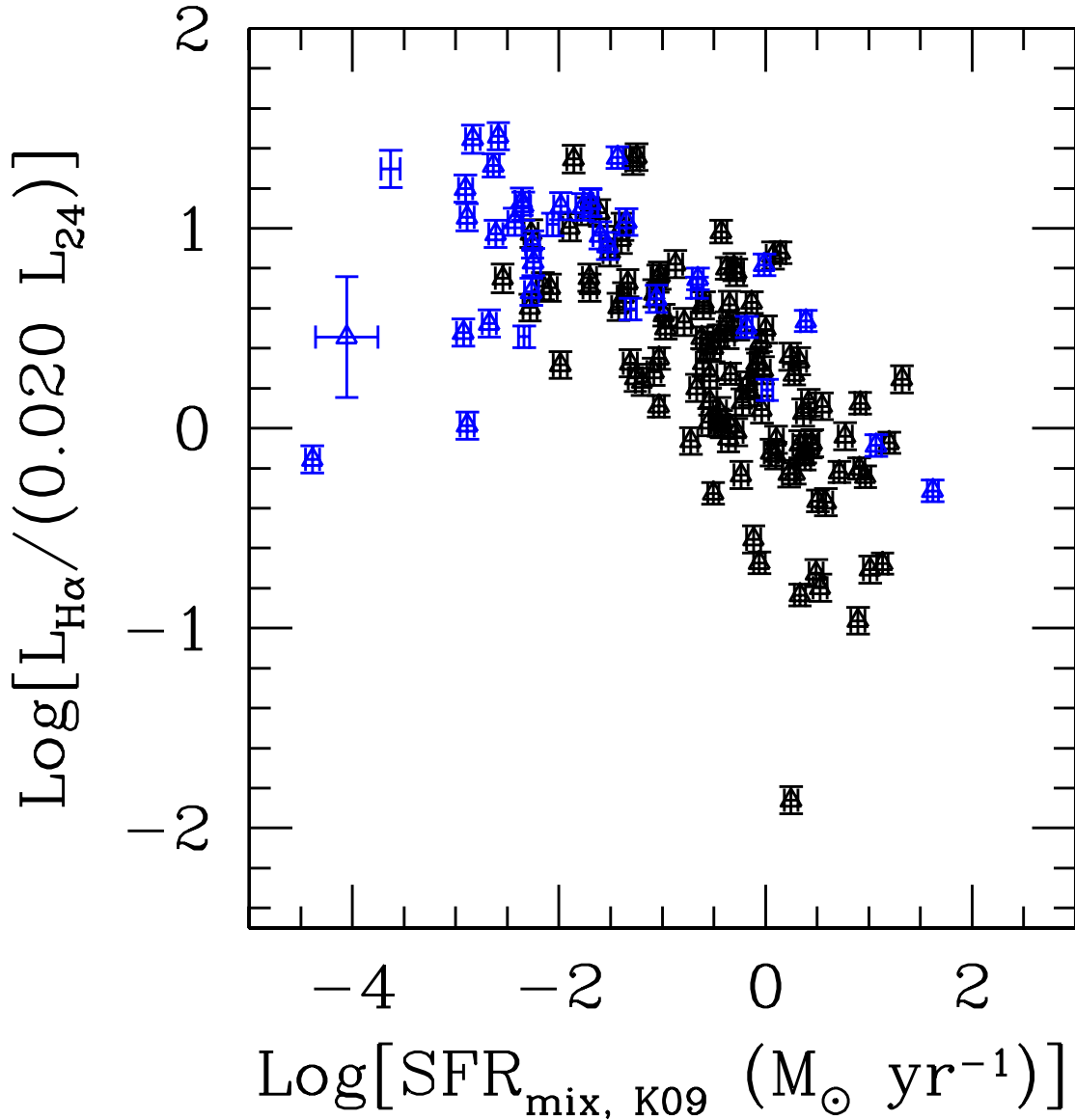


Fig. 7.— The ratio between the H $\alpha$  luminosity and the scaled 24  $\mu\text{m}$  luminosity for the galaxies in our sample (excluding LIRGs) as a function of SFR (equation 15). Blue points mark low-metallicity galaxies, and black points the high-metallicity ones. The scaling factor for the 24  $\mu\text{m}$  luminosity used on the vertical axis of this figure is from equation 15 (Kennicutt et al. 2009). This luminosity ratio measures the fraction of the star formation that emerges from the galaxy unattenuated by dust. Not surprisingly, there is a correlation between this luminosity ratio and SFR, in the sense that more actively star-forming galaxies are fainter in the optical relative to the infrared emission (see, also, Calzetti et al. 2007). In our normal star-forming and starburst sample, most of the galaxies tend to be relatively transparent, meaning that a large fraction of the light from recent star formation emerges unabsorbed by dust.

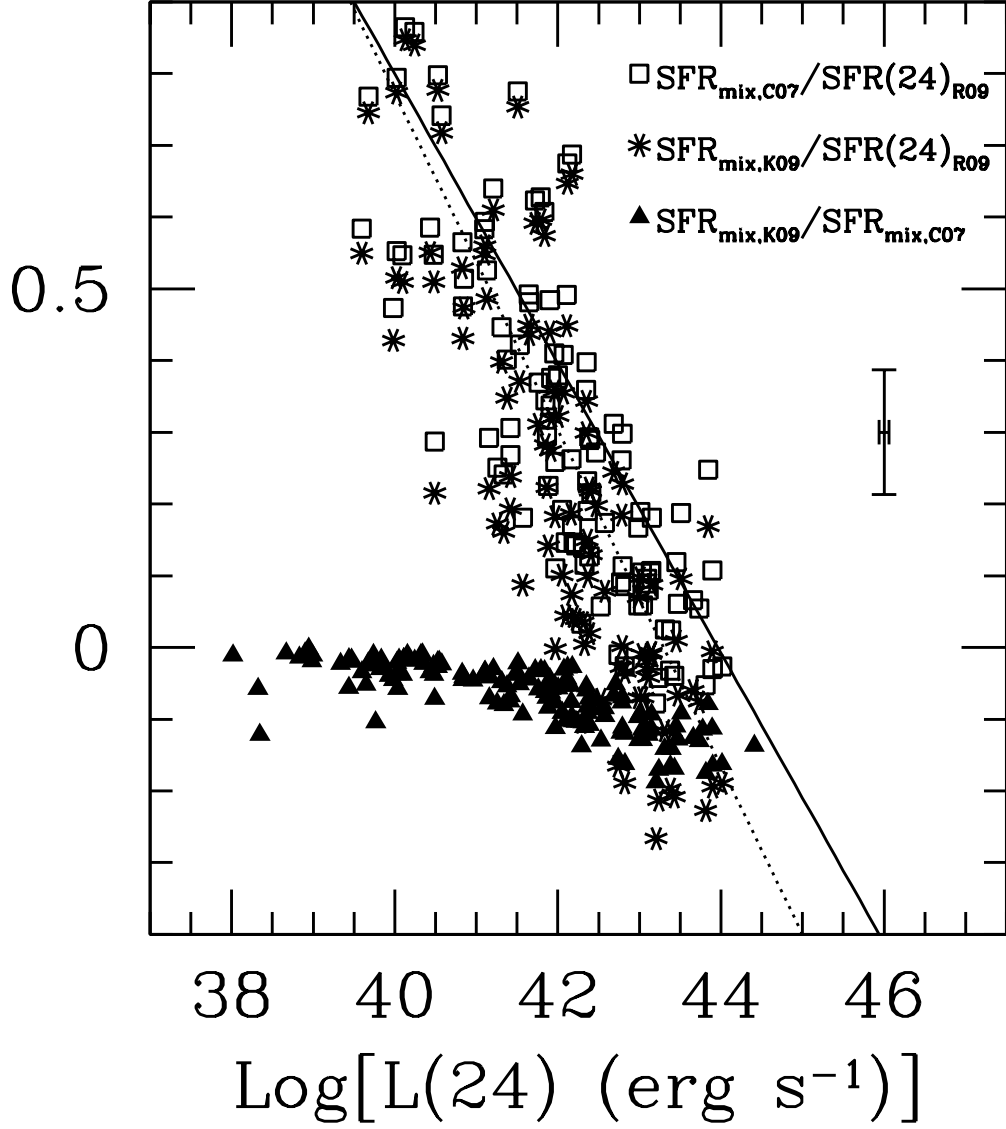


Fig. 8.— The logarithm of the ratio of SFR calibrations:  $\text{SFR}_{\text{mix,K09}}/\text{SFR}_{\text{mix,C07}}$  (filled triangles, equations 15 and 16),  $\text{SFR}_{\text{mix,K09}}/\text{SFR}(24)_{\text{R09}}$  (asterisks, equations 15 and 8), and  $\text{SFR}_{\text{mix,C07}}/\text{SFR}(24)_{\text{R09}}$  (empty squares, equations 16 and 8) as a function of the  $24 \mu\text{m}$  luminosity (LIRGs are excluded). Fits to the data from the latter two ratios are shown. The calibration  $\text{SFR}(24)_{\text{R09}}$  is extended below its limit of validity,  $\text{Log}[L(24)] \sim 42.6$ , for illustrative purposes only. A typical  $1\sigma$  errorbar is shown as well.

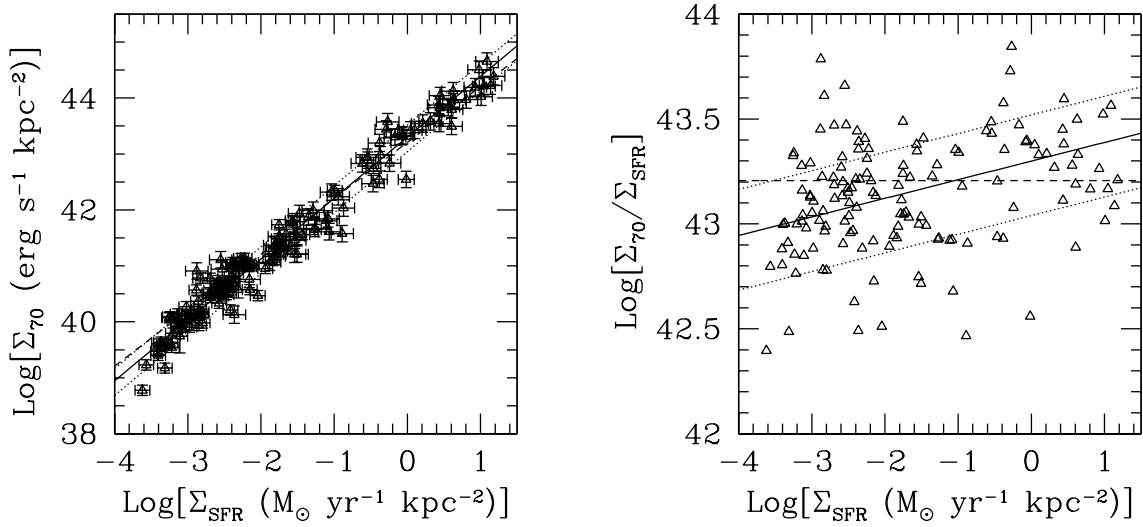


Fig. 9.— **(Left)** The 70  $\mu\text{m}$  LSD,  $\Sigma_{70}=\text{L}(70)/\text{area}$ , as a function of the SFSD,  $\Sigma_{SFR}$ , in log-log scale. In both cases, the area used is the  $\text{H}\alpha$  emitting area (Figure 2). Only the high-metallicity galaxies are reported (142 datapoints, with their  $1\sigma$  errorbars), together with the best fit (continuous line) and the  $1\sigma$  envelope to the data (dotted lines, which correspond to the range  $-0.26,+0.22$  dex). The  $1\sigma$  envelope is calculated as the standard deviation of the logarithm of the ratio relative to the mean trend. The best fit with unity slope is shown as a dashed line. **(Right)** The same data and best fit lines are shown now for the ratio  $\Sigma_{70}/\Sigma_{SFR}$  as a function of the SFSD. Errorbars are omitted for clarity. This plot better shows the deviation of the datapoints from a one-to-one relation.

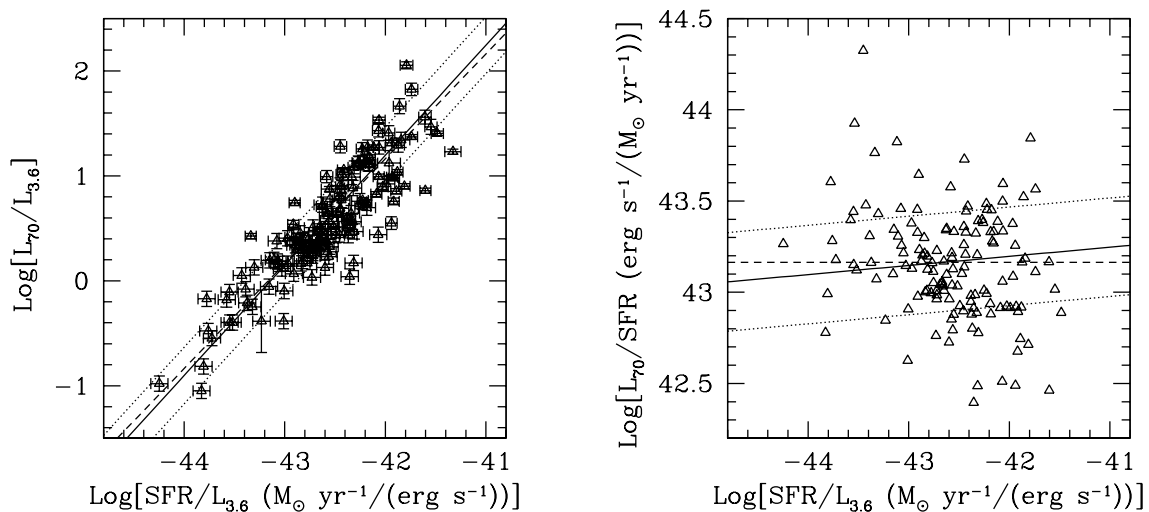


Fig. 10.— **(left)** and **(right)**. The same as Figure 9, with the data normalized by the  $3.6 \mu\text{m}$  luminosity, instead of the emitting area. The same data are shown, together with the best fit (continuous line) and the  $1 \sigma$  envelope to the data (dotted lines, which correspond to the range  $-0.27, +0.27$  dex). The best fit with unity slope is shown as a dashed line in both panels.

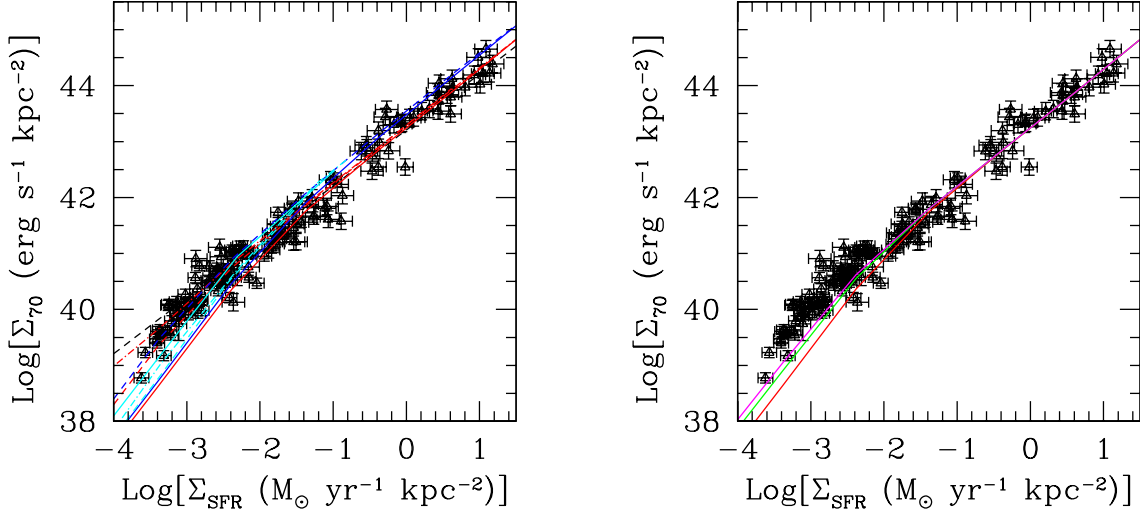


Fig. 11.— The same data on the 70  $\mu\text{m}$  LSD,  $\Sigma_{70}$ , versus the SFSD,  $\Sigma_{SFR}$  as in Figure 9. **Left.** The data are compared with the simple galaxy models described in Figure 4 (see text for more details). The best linear fit with unity slope is shown as a black dashed line (from Figure 9). All stellar population models are for solar metallicity: 10 Gyr old (blue continuous) and 100 Myr old (red continuous) constant star formation populations, and exponentially decreasing star formation (cyan; dash:  $\tau=5$  Gyr; continuous:  $\tau=2$  Gyr), together with our default dust model, in which the dust extinction is described by the starburst curve (Calzetti 2001), the dust emission comes from a medium of increasing opacity, and the peak infrared emission progressively moves into the 70  $\mu\text{m}$  band from longer wavelengths (Draine & Li 2007) for increasing SFSD. Dashed red and blue lines show the same default dust models, but the star-forming regions have a 10% filling factor within the galaxies, for stellar population with 100 Myr and 10 Gyr constant star formation, respectively. The case of a constant and large dust opacity ( $A_V > 10$  mag) with SFSD is shown as a red dot-dash line for a 100 Myr constant star formation population. The cyan dot-dash line is a decreasing star formation model, with  $\tau=2$  Gyr, embedded in a uniform distribution of dust (mixed dust–star case). **Right.** The 100 Myr constant star formation model is here reproduced for three different dust attenuation curves: the starburst curve (red line), used in our default model, according to which stars and ionized gas display differential attenuation (by a factor 0.44, see Calzetti, Kinney & Storchi–Bergmann 1994; Calzetti et al. 2000); the Milky Way (green line) and SMC (magenta line) extinction curves, with foreground dust geometry, and stars and gas attenuated by the same dust column density.

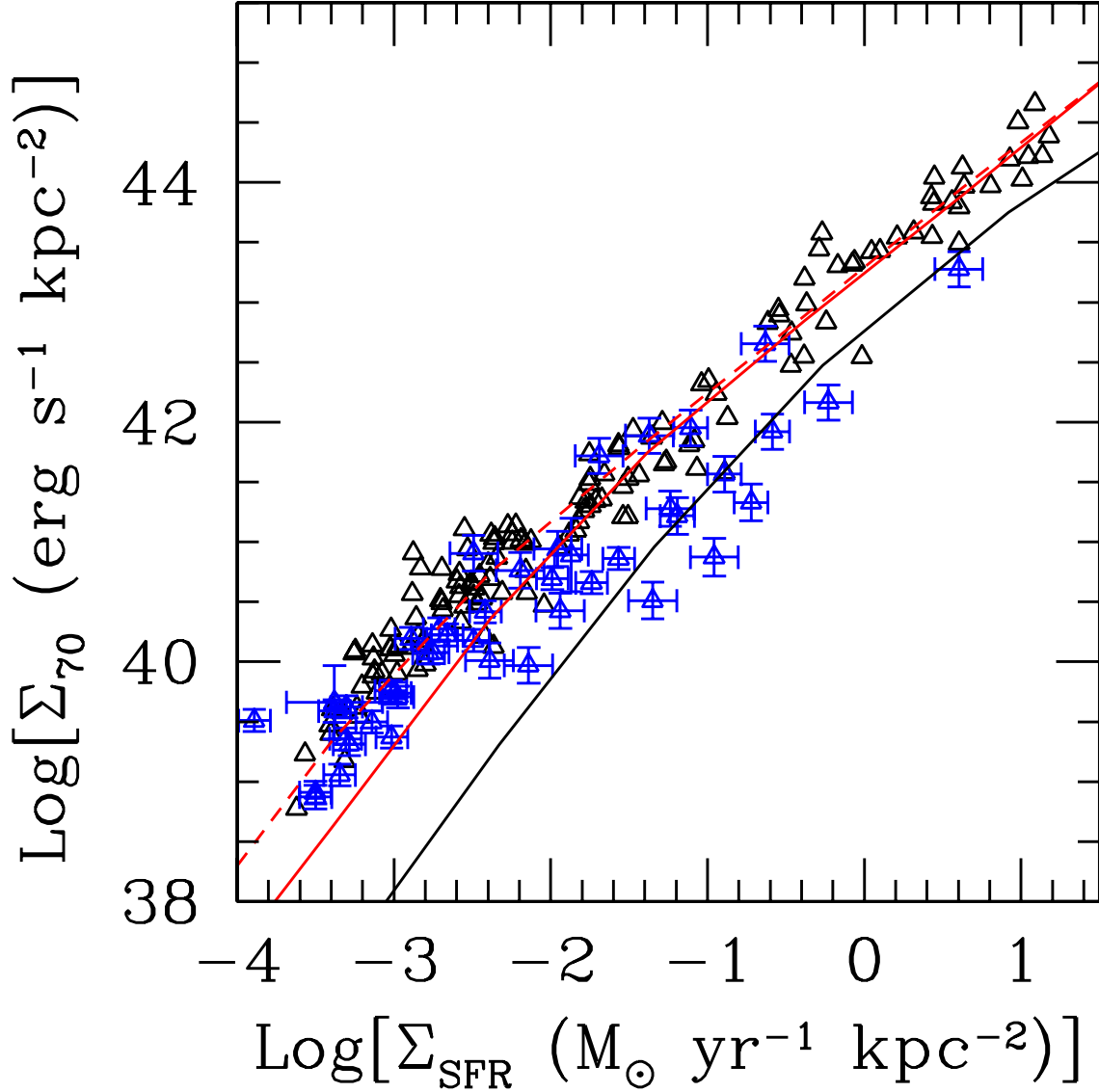


Fig. 12.— The data on the 70  $\mu\text{m}$  LSD,  $\Sigma_{70}$ , versus the SFSD,  $\Sigma_{\text{SFR}}$  now include both metal-rich (black symbols) and metal-poor (blue symbols) galaxies, as defined in the text.  $1\text{-}\sigma$  error bars are reported for the low metallicity datapoints, while the error bars for the metal-rich galaxies are omitted for clarity. The model of a 100 Myr continuous star formation population is used for reference, from Figure 11. The default solar metallicity model (red continuous line) and the 10% filling-factor one (red dashed line) are compared with a 0.1 solar metallicity model (continuous black line, from Calzetti et al. 2007).

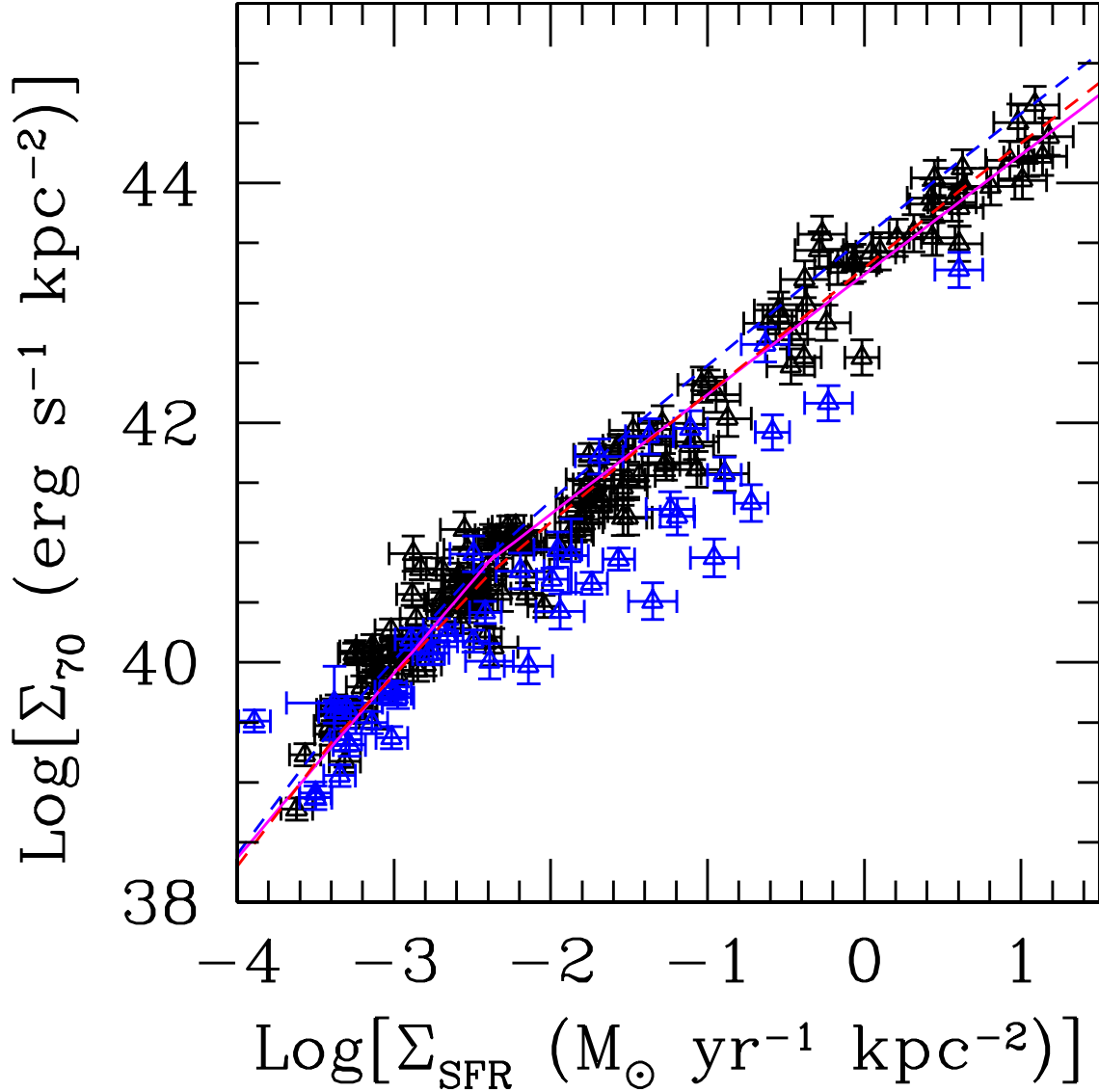


Fig. 13.— The data on the 70  $\mu\text{m}$  LSD,  $\Sigma_{70}$ , versus the SFSD,  $\Sigma_{SFR}$  for both metal-rich (black symbols) and metal-poor (blue symbols) galaxies. The data are the same as in Figure 12, except that now errorbars are included for all datapoints. Models of stellar populations with constant star formation over 100 Myr (red dashed line) and 10 Gyr (blue dashed line) surrounded by a solar metallicity ISM with a 10% area filling factor at the low luminosity end are also reported, and compared with the best power-law fits to the asymptotic trends at the low and high luminosity ends (broken magenta line, equations 20 and 21).

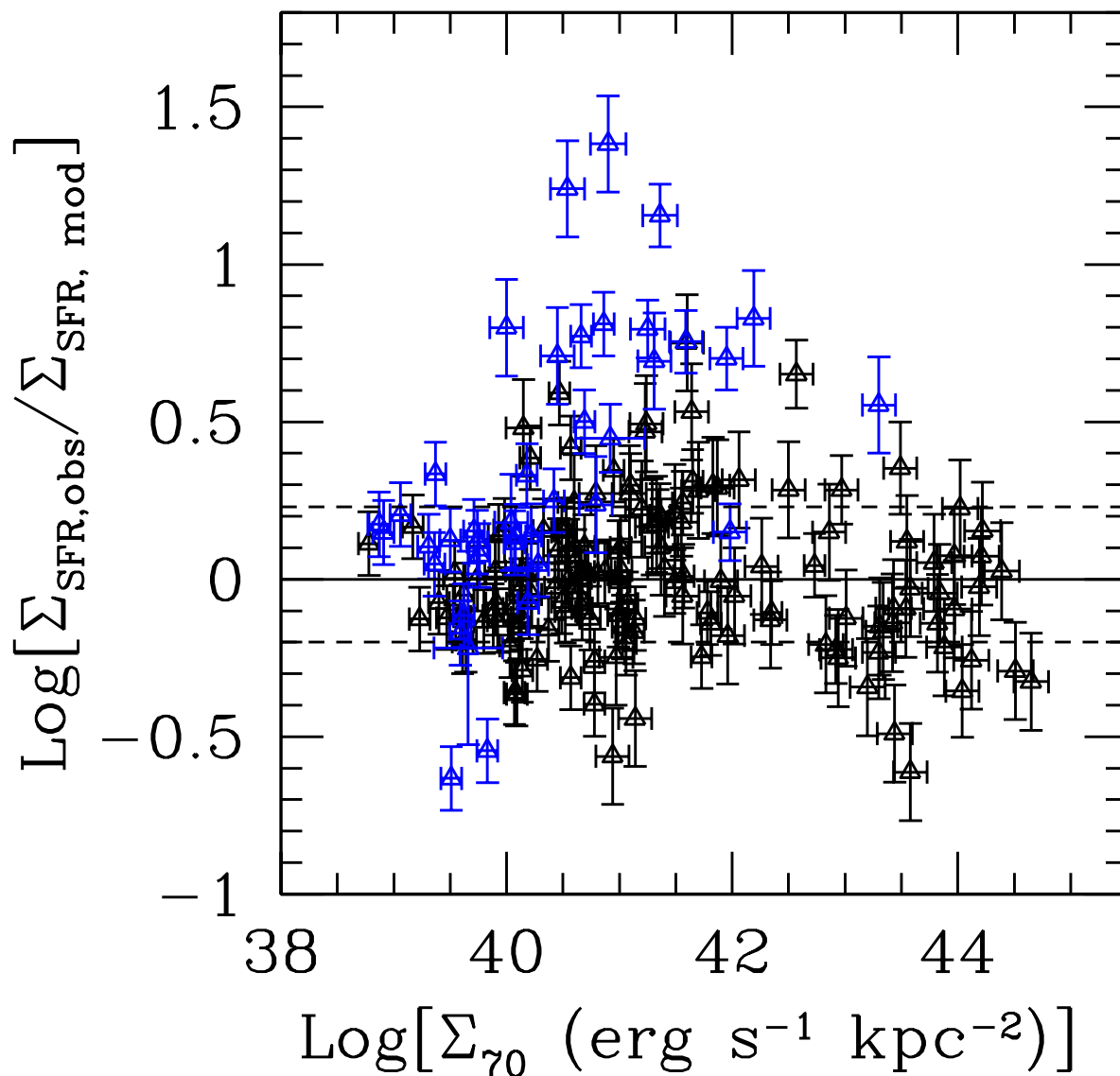


Fig. 14.— The difference between the data and the fitting lines described in equations 20 and 21 is reported as a function of the  $70 \mu\text{m}$  LSD, together with horizontal lines showing identity between models and data (continuous line) and the  $1 \sigma$  dispersion value ( $-0.20$ ,  $+0.23$ , dashed lines).



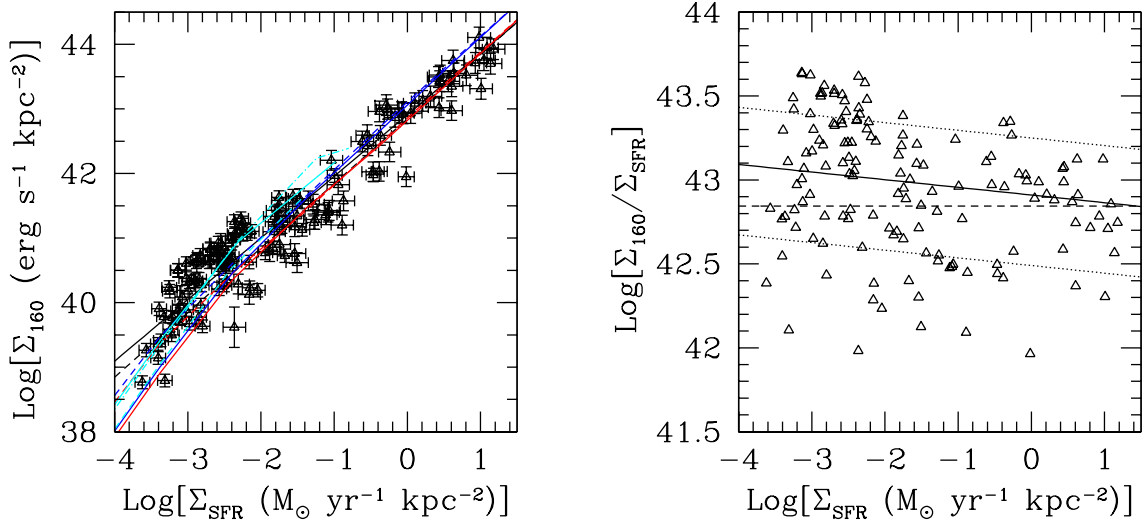


Fig. 15.— **Left.** The 160  $\mu\text{m}$  LSD,  $\Sigma_{160}=\text{L}(160)/\text{area}$ , versus the SFSD,  $\Sigma_{\text{SFR}}$ , for the metal-rich galaxies in our sample. Both best linear fits and model lines are overplotted on the data. The best fit has a slope slightly less than unity (continuous black line), and a line through the data with unity slope is shown as a dashed black line. The models are the same of Figure 11. **Right.** Similarly to the right-hand panel of Figure 9, the same data and best linear fit (continuous line) are shown, together with the  $1\sigma$  envelope to the data (dotted lines), for the ratio  $\text{L}(160)/\text{SFR}$  as a function of the SFSD. The best fit with unity slope is shown as a dashed line. This plot better shows the deviation of the datapoints from a one-to-one relation. The error bars on the data are not shown, for clarity.

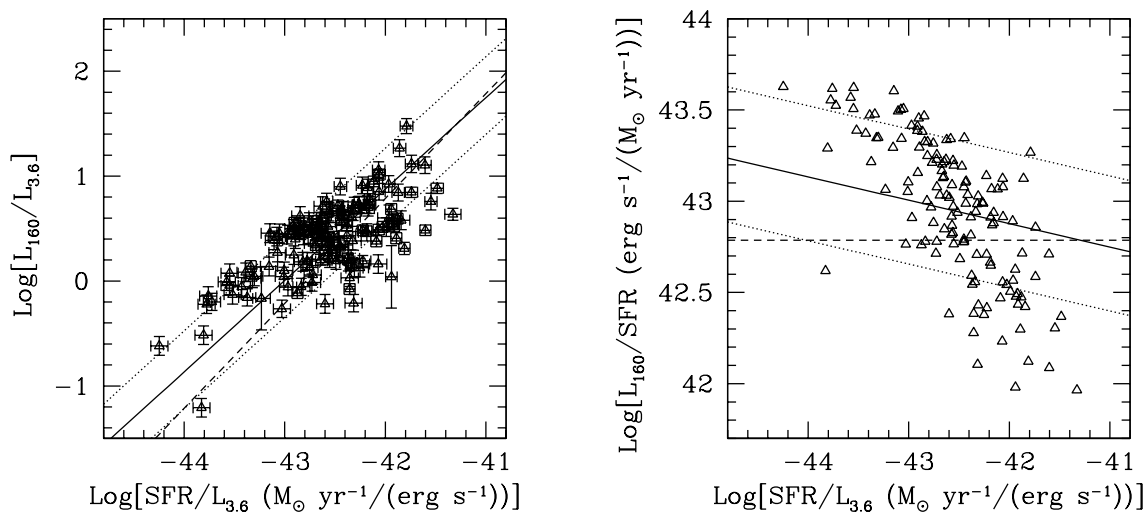


Fig. 16.— **(left)** and **(right)**. The same data as Figure 15, with the data normalized by the 3.6  $\mu\text{m}$  luminosity, instead of the emitting area. The best fit (continuous line) and the  $1\sigma$  envelope to the data (dotted lines, which correspond to the range  $-0.35, +0.39$  dex) are also shown. A line through the data with unity slope is shown in both panels (dashed line).

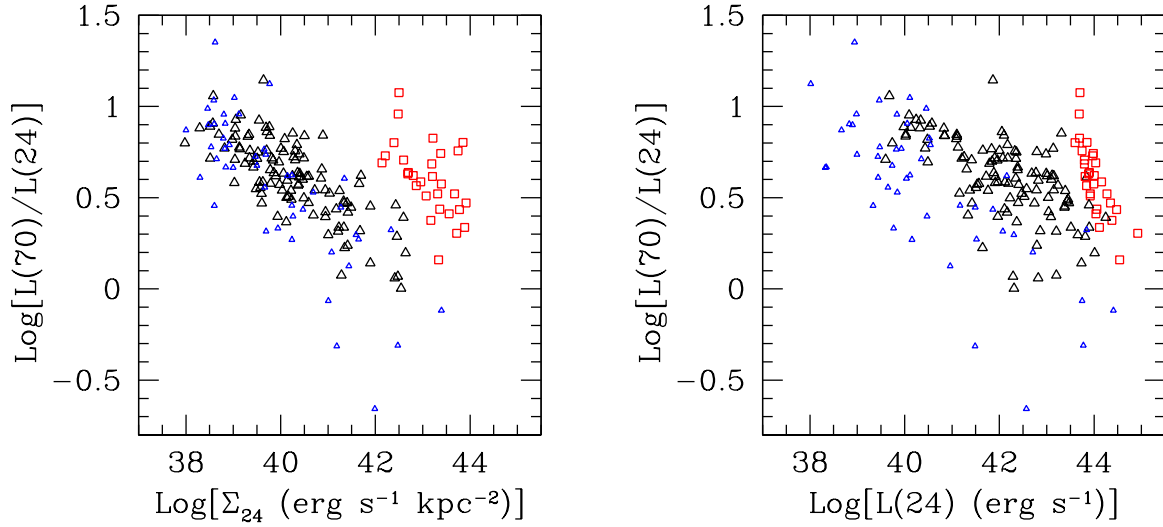


Fig. 17.— The  $L(70)/L(24)$  ratio as a function of the  $24 \mu\text{m}$  LSD,  $\Sigma_{24}$ , (left) and luminosity,  $L(24)$ , (right), respectively. As in Figure 1, colors identify the star-forming and starburst high-metallicity (black symbols) and low-metallicity (blue symbols) galaxies, and the LIRGs (red). Errorbars are omitted for clarity. The decreasing trend of  $L(70)/L(24)$  as a function of increasing LSD for normal star-forming and starburst galaxies (black and blue symbols) observed in the left-hand-side panel shows the variation in the dust effective temperature for these systems, similarly to Figure 3. The trend shows a large scatter, when the ratio  $L(70)/L(24)$  is plotted as a function of luminosity (right-hand-side panel).

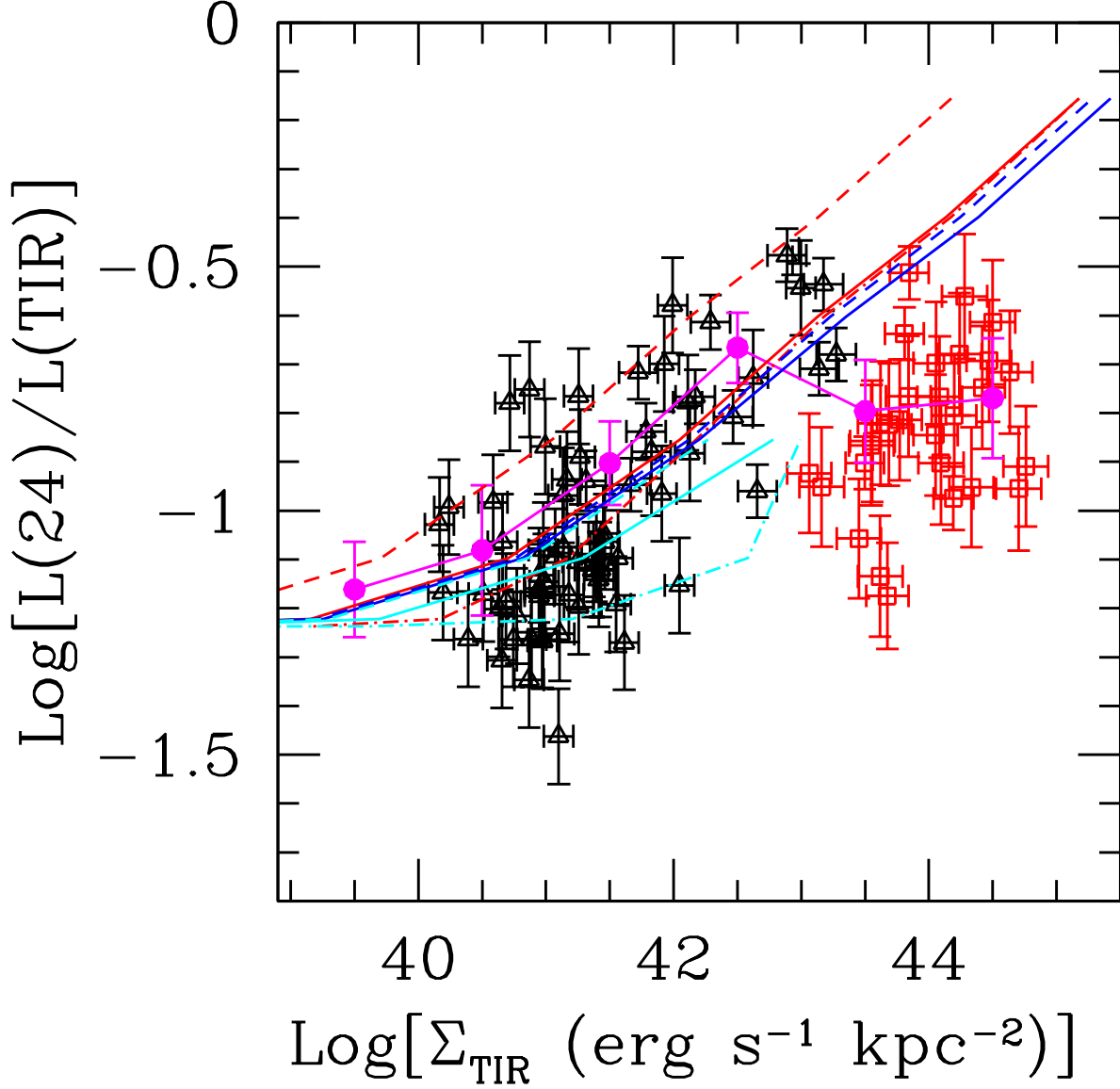


Fig. 18.— The  $L(24)/L(\text{TIR})$  ratio as a function of the IRSD for the galaxies in our sample with oxygen abundance  $12+\text{Log}(\text{O}/\text{H})>8.5$  (black and red points). The magenta points and line, and the models lines (blue, red, and cyan) are as in Figure 4, left. The trend for the galaxies with  $12+\text{Log}(\text{O}/\text{H})>8.5$  is similar to that of the whole sample with  $12+\text{Log}(\text{O}/\text{H})>8.1$ .

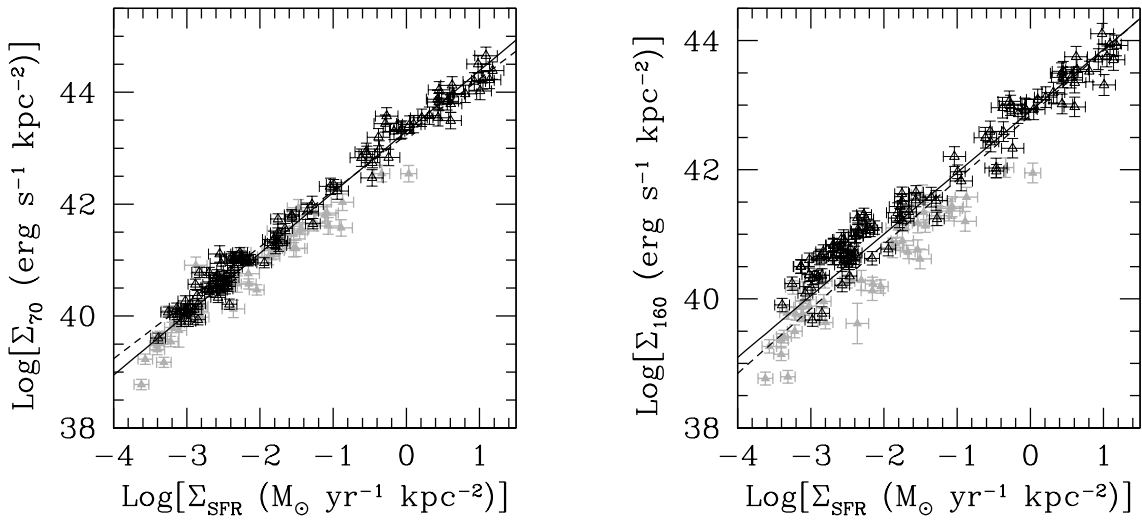


Fig. 19.— The 70  $\mu\text{m}$  LSD (**left**) and 160  $\mu\text{m}$  LSD (**right**) as a function of  $\Sigma_{\text{SFR}}$ . The high-metallicity data have been divided into two subsamples according to the galaxies oxygen abundance:  $8.1 < 12 + \text{Log}(\text{O}/\text{H}) \leq 8.5$  (grey triangles) and  $12 + \text{Log}(\text{O}/\text{H}) > 8.5$  (black triangles). See text for an explanation of the selected metallicity boundary. The scatter in the data decreases when restricting the sample to the highest metallicity values. The continuous and dashed lines in both panels are the best fit (equations 18 and 23) and the unity slope lines, respectively.

Table 1. Properties of Normal Star-Forming and Starburst Galaxies.

Name <sup>a</sup>	D <sup>b</sup> (Mpc)	12+Log(O/H) <sup>c</sup>	R <sub>25</sub> <sup>d</sup> (")	R <sub>H<math>\alpha</math></sub> <sup>e</sup> (")	Ref. <sup>e</sup>	Log( $\Sigma_{SFR}$ ) <sup>f</sup> (M $_{\odot}$ yr <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{70}$ ) <sup>g</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{160}$ ) <sup>g</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{TIR}$ ) <sup>g</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Ref. <sup>h</sup>
High Metallicity Galaxies										
NGC0625	4.07	8.1	172.4	60.2	L	-1.838±0.101	41.10±0.09	40.85±0.10	41.43±0.11	K08,D09
NGC1522	9.99	8.2	36.05	13.5	L	-1.435±0.101	41.56±0.09	41.13±0.10	41.80±0.11	K08,D09
NGC1800	8.83	8.36	59.85	19.6	L	-1.713±0.101	41.34±0.09	41.17±0.10	41.62±0.11	K08,D09
UGC4278	8.13	8.08	140.3	144.8	L	-3.407±0.101	39.47±0.09	39.36±0.10	39.77±0.11	K08,D09
UGC5829	8.44	8.3	140.3	140.8	L	-3.404±0.101	39.40±0.09	39.14±0.10	39.63±0.11	K08,D09
NGC3368	11.27	9.04	227.6	84.8	L	-2.274±0.101	41.13±0.09	41.30±0.10	41.61±0.11	K08,D09
UGC5923	7.67	8.26	28	19.6	L	-2.501±0.101	40.65±0.09	40.27±0.10	40.83±0.12	K08,D09
NGC3510	9.18	8.08	119.5	99.6	L	-3.123±0.101	39.92±0.09	39.75±0.10	40.19±0.11	K08,D09
NGC3623	9.59	9.06	293.1	125.2	L	-2.693±0.101	40.49±0.09	40.84±0.10	41.11±0.11	K08,D09
UGC6900	8	8.1	62.7	27.8	L	-3.107±0.103	39.74±0.30	39.96±0.30	40.26±0.38	K08,D09
NGC4248	7.76	8.15	90.6	27.4	L	-2.508±0.101	40.66±0.09	40.71±0.10	41.08±0.11	K08,D09
NGC4288	8.22	8.52	64.15	48.2	L	-2.486±0.101	40.61±0.09	40.55±0.10	40.95±0.11	K08,D09
UGC7490	9	8.46	99.35	83.5	L	-3.325±0.101	39.59±0.09	39.78±0.10	40.08±0.11	K08,D09
UGC7699	7.34	8.15	114.1	89.6	L	-3.204±0.101	39.80±0.09	39.77±0.10	40.15±0.11	K08,D09
UGC7698	6.1	8.04	193.7	145.2	L	-3.620±0.101	38.78±0.09	38.77±0.10	39.14±0.12	K08,D09
NGC4656	9.2	8.78	454	282.6	L	-2.979±0.101	39.90±0.09	39.67±0.10	40.17±0.11	K08,D09
UGC7950	8.48	8.37	38.65	40.4	L	-3.024±0.101	40.11±0.09	39.89±0.10	40.35±0.12	K08,D09
NGC4707	7.97	8.43	67.15	67.2	L	-3.219±0.101	39.54±0.09	39.50±0.10	39.89±0.12	K08,D09
UGCA320	7.76	8.08	168.7	155.2	L	-3.313±0.101	39.17±0.09	38.79±0.10	39.36±0.12	K08,D09
UGC8320	4.33	8.29	108.9	95.2	L	-3.367±0.101	39.64±0.09	39.42±0.10	39.87±0.12	K08,D09
NGC5068	6.24	8.82	217.4	206.7	L	-2.438±0.101	40.53±0.09	40.59±0.10	40.94±0.11	K08,D09
NGC5477	8.25	8.14	49.8	55.6	L	-2.798±0.101	39.98±0.09	39.63±0.10	40.17±0.12	K08,D09
NGC0024	8.13	8.62	172.7	143.	L,S	-3.075±0.101	39.91±0.09	40.09±0.10	40.39±0.11	K09,D07
NGC0337	24.69	8.56	86.5	109.	S	-2.399±0.101	40.81±0.09	40.71±0.10	41.14±0.11	K09,D07
NGC0628	7.32	8.67	314.1	261.	L,S	-2.583±0.101	40.55±0.09	40.75±0.10	41.07±0.11	K09,D07
NGC0925	9.12	8.51	314.1	296.	S	-2.995±0.101	40.06±0.09	40.18±0.10	40.51±0.11	K09,D07
NGC1097	16.88	8.78	280.	282.	S	-2.453±0.101	40.72±0.09	40.77±0.10	41.16±0.11	K09,D07
NGC1291	9.83	8.72	293.1	329.	L,S	-3.392±0.101	39.61±0.09	39.90±0.10	40.19±0.11	K09,D07
NGC1377	24.17	8.37	53.35	44.3	S	-1.674±0.101	41.36±0.09	40.73±0.10	41.74±0.11	K09,D07
NGC1482	22.03	8.53	73.65	59.1	S	-1.567±0.101	41.81±0.09	41.53±0.10	42.11±0.11	K09,D07

Table 1—Continued

Name <sup>a</sup>	D <sup>b</sup> (Mpc)	12+Log(O/H) <sup>c</sup>	R <sub>25</sub> <sup>d</sup> (")	R <sub>H<math>\alpha</math></sub> <sup>e</sup> (")	Ref. <sup>e</sup>	Log( $\Sigma_{SFR}$ ) <sup>f</sup> (M <sub><math>\odot</math></sub> yr <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{70}$ ) <sup>g</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{160}$ ) <sup>g</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{TIR}$ ) <sup>g</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Ref. <sup>h</sup>
NGC1512	10.14	8.71	267.4	185.	L,S	-2.977±0.101	40.13±0.09	40.32±0.10	40.63±0.11	K09,D07
NGC1566	18.27	8.82	249.6	269.	S	-2.704±0.101	40.52±0.09	40.63±0.10	40.98±0.11	K09,D07
NGC1705	5.1	8.32	57.15	57.2	L,S	-2.043±0.101	40.47±0.09	40.19±0.10	40.69±0.11	K09,D07
NGC2403	3.22	8.56	656.3	413.	L,S	-2.492±0.101	40.55±0.09	40.64±0.10	40.99±0.11	K09,D07
NGC2798	24.68	8.69	77.1	26.0	S	-0.992±0.101	42.35±0.09	41.97±0.10	42.62±0.11	K09,D07
NGC2841	9.81	8.86	243.9	209.	S	-2.418±0.101	40.21±0.09	40.64±0.10	40.89±0.11	K09,D07
NGC2915	3.78	8.16	57.15	51.4	G03	-2.152±0.101	40.57±0.09	40.23±0.10	40.77±0.11	G03,D07
NGC2976	3.56	8.64	176.7	117.	L,S	-2.189±0.101	41.02±0.09	41.07±0.10	41.43±0.11	K09,D07
NGC3049	19.1	8.75	65.65	81.3	S	-2.471±0.101	40.49±0.09	40.35±0.10	40.87±0.11	K09,D07
NGC3031	3.63	8.69	692.7	680.	L,S	-3.020±0.101	40.11±0.09	40.37±0.10	40.65±0.11	P06,D07
NGC3034	3.89	8.72	336.6	143.	L,S	-0.463±0.101	42.74±0.09	41.98±0.10	43.00±0.11	K09,D07
NGC3190	16.81	8.7	130.9	103.	S	-2.882±0.101	40.57±0.09	40.63±0.10	40.97±0.11	K09,D07
NGC3184	8.53	8.81	222.4	217.	S	-2.859±0.101	40.37±0.09	40.66±0.10	40.95±0.11	K09,D07
NGC3198	13.68	8.6	255.4	233.	S	-2.897±0.101	40.12±0.09	40.34±0.10	40.66±0.11	K09,D07
NGC3265	19.49	8.65	38.65	44.3	S	-2.162±0.101	40.99±0.09	40.63±0.10	41.26±0.11	K09,D07
Mrk33	21.92	8.56	30.	26.1	S	-1.276±0.101	41.65±0.09	41.24±0.10	41.99±0.11	K09,D07
NGC3351	9.33	8.91	222.4	189.	L,S	-2.597±0.101	40.67±0.09	40.76±0.10	41.13±0.11	K08,D07
NGC3521	8.99	8.68	328.9	284.	L,S	-2.583±0.101	40.74±0.09	40.92±0.10	41.25±0.11	K08,D07
NGC3621	6.64	8.5	369.1	531.	S	-3.246±0.101	40.09±0.09	40.18±0.10	40.53±0.11	K09,D07
NGC3627	9.38	8.8	273.6	244.	L,S	-2.358±0.101	41.03±0.09	41.07±0.10	41.45±0.11	K09,D07
NGC3773	12.53	8.64	35.25	35.1	S	-1.941±0.101	40.95±0.09	40.77±0.10	41.27±0.11	K09,D07
NGC3938	12.22	8.71	161.1	191.	S	-2.687±0.101	40.43±0.09	40.64±0.10	40.95±0.11	K09,D07
NGC4125	22.91	8.89	172.7	95.	S	-2.848±0.101	39.93±0.09	39.77±0.10	40.24±0.11	K09,D07
NGC4236	4.45	8.31	656.3	582.	L,S	-3.567±0.101	39.23±0.09	39.26±0.10	39.63±0.11	K08,D07
NGC4254	33.29	8.8	161.1	185.	S	-2.126±0.101	41.01±0.09	41.10±0.10	41.46±0.11	K09,D07
NGC4321	14.32	8.84	222.4	195.	S	-2.343±0.101	40.87±0.09	41.05±0.10	41.37±0.11	K09,D07
NGC4450	27.07	8.85	157.4	113.	S	-3.019±0.101	40.27±0.09	40.61±0.10	40.87±0.11	K09,D07
NGC4536	14.45	8.6	227.6	229.	S	-2.574±0.101	40.63±0.09	40.53±0.10	40.99±0.11	K09,D07
NGC4559	11.57	8.51	321.4	267.	S	-2.845±0.101	40.22±0.09	40.36±0.10	40.69±0.11	K09,D07
NGC4569	16	8.88	286.5	270.	K01	-3.254±0.101	40.07±0.09	40.24±0.10	40.58±0.11	B02,D07
NGC4579	20.62	8.93	176.7	116.	S	-2.387±0.101	40.69±0.09	40.97±0.10	41.26±0.11	K09,D07
NGC4594	9.33	8.99	261.3	201.	L,S	-3.138±0.101	40.14±0.09	40.50±0.10	40.77±0.11	K09,D07
NGC4625	9.51	8.65	65.65	38.0	L,S	-2.245±0.101	41.00±0.09	41.06±0.10	41.41±0.11	K09,D07

Table 1—Continued

Name <sup>a</sup>	D <sup>b</sup> (Mpc)	12+Log(O/H) <sup>c</sup>	R <sub>25</sub> <sup>d</sup> (")	R <sub>H<math>\alpha</math></sub> <sup>e</sup> (")	Ref. <sup>e</sup>	Log( $\Sigma_{SFR}$ ) <sup>f</sup> (M <sub>⊙</sub> yr <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{70}$ ) <sup>g</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{160}$ ) <sup>g</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{TIR}$ ) <sup>g</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Ref. <sup>h</sup>
NGC4631	8.92	8.44	464.6	257.	L,S	-2.222±0.101	41.14±0.09	41.12±0.10	41.51±0.11	K09,D07
NGC4725	17.12	8.73	321.4	240.	S	-3.132±0.101	40.03±0.09	40.50±0.10	40.74±0.11	K09,D07
NGC4736	5.01	8.66	336.6	238.	L,S	-2.380±0.101	41.06±0.09	40.98±0.10	41.39±0.11	K08,D07
NGC4826	5.54	8.92	300	84.	L,S	-1.753±0.101	41.73±0.09	41.63±0.10	42.05±0.11	K09,D07
NGC5055	8.32	8.78	377.7	226.	L,S	-2.361±0.101	40.99±0.09	41.25±0.10	41.55±0.11	K09,D07
NGC5194	8.13	8.86	336.6	296.	L,S	-2.244±0.101	41.07±0.09	41.23±0.10	41.56±0.11	K09,D07
Tol89	14.9	8.54	84.55	81.	S	-2.572±0.101	40.33±0.09	40.21±0.10	40.72±0.11	K09,D07
NGC5474	6.8	8.32	143.6	131.	L,S	-2.810±0.101	40.18±0.09	40.27±0.10	40.60±0.11	K09,D07
NGC5713	26.49	8.64	82.65	52.2	S	-1.563±0.101	41.78±0.09	41.65±0.10	42.13±0.11	K09,D07
NGC5866	15.14	8.72	140.3	100.	S	-2.830±0.101	40.78±0.09	40.73±0.10	41.10±0.11	K09,D07
IC4710	8.48	8.37	108.9	109.	S	-2.823±0.101	40.14±0.09	39.96±0.10	40.42±0.11	K09,D07
NGC6822	0.47	8.35	464.6	765.	S	-3.133±0.101	39.88±0.09	39.87±0.10	40.24±0.11	K09,D07
NGC6946	6.8	8.72	344.4	217.	S	-1.755±0.101	41.49±0.09	41.51±0.10	41.91±0.11	K09,D07
NGC7331	14.52	8.7	314.1	295.	S	-2.693±0.101	40.78±0.09	40.82±0.10	41.18±0.11	K09,D07
NGC7552	22.27	8.74	101.7	113.	S	-1.657±0.101	41.57±0.09	41.35±0.10	41.93±0.11	K09,D07
NGC7793	3.82	8.53	255	286.	L,S	-2.549±0.101	40.46±0.09	40.67±0.10	40.97±0.11	K09,D07
Mrk170	20.4	8.09	31.4	18.1	A	> -2.877±0.153	40.91±0.15	40.62±0.15	41.14±0.15	...,E08
Mrk930	77	8.11	25.8	14.9	A	> -1.748±0.153	41.53±0.15	40.90±0.15	41.79±0.15	...,E08
NGC1569	1.9	8.13	108.9	85.9	H04	-0.890±0.153	41.58±0.15	41.20±0.15	41.93±0.15	H04,E08
Mrk1094	41	8.15	20.75	19.8	G03	-1.540±0.153	41.21±0.15	40.76±0.15	41.40±0.15	G03,E08
NGC3310	21.3	8.18	92.7	67.2	J04	-1.263±0.153	41.67±0.15	41.29±0.15	41.98±0.15	J04,E08
Mrk162	98	8.19	18.1	10.5	A	> -1.474±0.153	41.94±0.15	41.61±0.15	42.23±0.15	...,E08
NGC1156	7.8	8.19	99.35	88.0	H04	-2.159±0.153	40.76±0.15	40.13±0.15	40.94±0.15	H04,E08
Tol2	22	8.22	25	19.4	G03	-1.506±0.153	41.21±0.15	40.62±0.15	41.38±0.15	G03,E08
MinkObj	78	8.22	15	13.9	C06	-2.364±0.153	40.13±0.15	39.62±0.31	40.34±0.32	C06,E08
NGC4449	4.2	8.23	185	133.5	L	-1.819±0.153	41.17±0.15	41.10±0.15	41.53±0.15	K08,E08
NGC7714	40	8.26	57.15	24.2	J04	-0.873±0.153	42.04±0.15	41.58±0.15	42.41±0.15	J04,E08
UGC4703	57	8.31	13.7	6.0	M99	-1.068±0.153	41.61±0.15	41.43±0.15	41.96±0.16	M99,E08
NGC1140	21.2	8.32	49.8	44.0	H94	-1.881±0.153	41.06±0.15	40.79±0.15	41.37±0.15	H94,E08
NGC1510	11.8	8.33	39.55	10.9	S	-1.114±0.153	41.81±0.15	41.36±0.15	42.06±0.15	K09,E08
NGC3125	12	8.34	32.9	22.9	G03	-1.083±0.153	41.84±0.15	41.40±0.15	42.11±0.15	G03,E08
NGC4214	2.9	8.36	255.4	188.6	H04	-2.310±0.153	40.57±0.15	40.29±0.15	40.86±0.15	K08,E08
NGC4670	23.2	8.38	42.4	28.1	G03	-1.539±0.153	41.46±0.15	41.18±0.15	41.73±0.15	G03,E08



Table 1—Continued

Name <sup>a</sup>	D <sup>b</sup> (Mpc)	12+Log(O/H) <sup>c</sup>	R <sub>25</sub> <sup>d</sup> (")	R <sub>H<math>\alpha</math></sub> <sup>e</sup> (")	Ref. <sup>e</sup>	Log( $\Sigma_{SFR}$ ) <sup>f</sup> (M <sub><math>\odot</math></sub> yr <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{70}$ ) <sup>g</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{160}$ ) <sup>g</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{TIR}$ ) <sup>g</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Ref. <sup>h</sup>
He2-10	9	8.55	52.15	20.7	J00	-0.468±0.153	42.47±0.15	42.03±0.15	42.89±0.15	J00,E08
NGC3628	13.1	8.57	443.8	240.	F90	-2.530±0.153	40.94±0.15	40.88±0.15	41.31±0.15	K08,E08
NGC3079	21.8	8.57	238.3	180.	R07	-2.551±0.153	41.11±0.15	40.92±0.15	41.38±0.15	L96,E08
NGC2782	42	8.59	104	30.	E96	-1.354±0.153	41.87±0.15	41.58±0.15	42.17±0.15	M06,E08
NGC3077	3.8	8.6	161.1	27.6	J04	-0.944±0.153	42.24±0.15	41.83±0.15	42.47±0.15	K08,E08
NGC3367	49	8.62	75.35	50.3	G01	-1.744±0.153	41.30±0.15	41.21±0.15	41.72±0.15	M06,E08
NGC5236	4.5	8.62	386.4	305.	L99	-1.774±0.153	41.34±0.15	41.43±0.15	41.82±0.15	K08,E08
NGC5953	35	8.67	48.65	17.	H03	-1.037±0.153	42.32±0.15	42.20±0.15	42.66±0.15	K87,E08
NGC4194	42	8.67	54.6	13.5	HT04	-0.242±0.153	42.84±0.15	42.33±0.15	43.17±0.15	HT04,E08
NGC2903	8.9	8.68	377.7	146.	B02	-1.814±0.153	41.37±0.15	41.33±0.15	41.78±0.15	K08,E08
Mrk25	48	8.68	16.5	9.5	A	> -1.287±0.153	41.99±0.15	41.52±0.15	42.29±0.15	...,E08
Mrk331	78	8.76	20.3	10.	HT04	-0.368±0.153	42.99±0.15	42.59±0.15	43.27±0.15	HT04,E08
IC342	3.3	8.85	641.4	380.	H05	-1.787±0.153	41.26±0.15	41.25±0.15	41.67±0.15	W99,E08
IIZw40	9.2	8.11	12.5	10.4	G03	-0.015±0.108	42.54±0.15	41.95±0.15	42.94±0.15	C07,E08
NGC5253	4	8.19	150.	22.2	C99	0.017±0.108	42.95±0.15	42.43±0.15	43.39±0.15	C07,E08
NGC2537	6.9	8.44	52.15	28.4	J04	-1.506±0.108	41.53±0.15	41.34±0.15	41.83±0.15	C07,E08
NGC2146	17.9	8.68	180.8	35.5	J04	-0.540±0.108	42.89±0.15	42.43±0.15	43.15±0.15	C07,E08
Low Metallicity Galaxies										
WLM	0.92	7.77	344.4	107.9	L	-2.751±0.101	40.08±0.09	40.00±0.10	40.40±0.11	K08,D09
UGC5272	7.61	7.83	62.7	62.6	L	-2.991±0.101	39.74±0.09	39.38±0.10	39.93±0.12	K08,D09
NGC3109	1.34	7.77	571.6	286.7	L	-3.020±0.101	39.76±0.09	39.65±0.10	40.06±0.11	K08,D09
SextansA	1.32	7.54	176.7	144.3	L	-3.015±0.101	39.37±0.09	39.15±0.10	39.62±0.12	K08,D09
UGC5764	7.59	7.95	59.85	45.21	L	-3.285±0.102	39.31±0.10	39.11±0.11	39.59±0.12	K08,D09
UGCA281	5.7	7.8	24.95	26.08	L	-1.739±0.101	40.66±0.09	39.89±0.10	40.87±0.12	K08,D09
UGC8508	2.69	7.89	50.95	56.8	L	-3.141±0.101	39.50±0.09	39.26±0.10	39.74±0.12	K08,D09
UGC8651	3.02	7.85	70.35	83.03	L	-3.500±0.102	38.87±0.10	38.76±0.11	39.18±0.13	K08,D09
UGC8837	8.89	7.7	128	93.46	L	-3.308±0.101	39.36±0.09	39.43±0.10	39.78±0.12	K08,D09
UGC9128	2.24	7.75	49.8	23.91	L	-3.379±0.307	39.66±0.31	39.88±0.31	40.18±0.44	K08,D09
UGC9240	2.8	7.95	54.6	39.99	L	-2.891±0.101	40.19±0.09	39.97±0.10	40.47±0.12	K08,D09
Mrk475	9.66	7.97	11.69	10.43	L	-1.566±0.101	40.86±0.09	40.06±0.11	40.99±0.13	K08,D09
UGC9992	9.17	7.88	48.65	45.64	L	-3.378±0.101	39.59±0.09	39.30±0.10	39.83±0.12	K08,D09

Table 1—Continued

Name <sup>a</sup>	D <sup>b</sup> (Mpc)	12+Log(O/H) <sup>c</sup>	R <sub>25</sub> <sup>d</sup> (")	R <sub>H<math>\alpha</math></sub> <sup>e</sup> (")	Ref. <sup>e</sup>	Log( $\Sigma_{SFR}$ ) <sup>f</sup> (M <sub>⊙</sub> yr <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{70}$ ) <sup>g</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{160}$ ) <sup>g</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{TIR}$ ) <sup>g</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Ref. <sup>h</sup>
ESO347–G17	10.04	7.9	41.4	40.39	L	-2.779±0.101	40.07±0.09	39.73±0.10	40.26±0.12	K08,D09
UGC12613	0.76	7.93	150.4	86.51	L	-3.891±0.101	39.51±0.09	39.69±0.10	40.01±0.12	K08,D09
UGCA442	4.27	7.72	106.5	101.4	L	-3.500±0.101	38.91±0.09	38.89±0.10	39.30±0.12	K08,D09
HolmII	3.39	7.88	238.3	224.	L,S	-2.976±0.101	39.71±0.09	39.43±0.10	39.95±0.11	K09,D07
M81DwA	3.55	7.49	0.01	39.3	L,S	-3.595±0.101	39.83±0.09	39.46±0.10	40.10±0.11	K09,D07
DDO053	3.56	7.82	46.45	43.2	L,S	-2.493±0.101	40.18±0.09	39.92±0.10	40.44±0.11	K09,D07
HolmI	3.84	8.00	108.9	83.	L,S	-3.302±0.101	39.62±0.09	39.60±0.10	39.96±0.11	K09,D07
HolmIX	3.28	7.8	75.35	29.4	L,S	-2.712±0.101	40.28±0.09	40.22±0.10	40.70±0.11	K09,D07
M81DwB	7.08	8.02	26.15	20.1	L,S	-2.418±0.101	40.42±0.09	40.47±0.10	40.83±0.11	K09,D07
DDO154	4.3	7.78	90.6	63.5	L,S	-3.346±0.101	39.06±0.09	39.43±0.10	39.71±0.11	K09,D07
DDO165	4.57	7.80	104	25.3	L,S	-2.659±0.101	40.22±0.09	40.20±0.10	40.62±0.11	K09,D07
IC2574	2.8	7.93	395.5	305.	L,S	-3.237±0.101	39.62±0.09	39.58±0.10	39.97±0.11	K08,D07
NGC5408	4.81	8.02	48.65	71.3	S	-1.991±0.101	40.69±0.09	40.19±0.10	40.93±0.11	K09,D07
IZw18	12.6	7.19	8.85	8.39	G03	-1.349±0.153	40.51±0.15	40.51±0.32	40.98±0.32	G03,E08
Tol65	34	7.45	10.45	4.85	G03	-0.957±0.153	40.87±0.16	40.99±0.32	41.63±0.32	G03,E08
UGC4483	3.2	7.55	33.65	16.35	G03	-1.939±0.154	40.43±0.15	39.95±0.19	40.60±0.19	K08,E08
ESO146–G14	23.8	7.66	84.6	48.84	A	> -4.506±0.156	39.45±0.15	39.30±0.17	39.72±0.17	...,E08
Mrk178	4.7	7.82	36.9	36	G03	-2.141±0.154	39.97±0.15	39.46±0.16	40.07±0.16	K08,E08
Mrk153	41	7.83	24.4	14.09	A	> -2.491±0.153	40.91±0.15	40.34±0.16	41.14±0.16	...,E08
UM462	13.4	7.91	19.35	11.17	A	> -1.691±0.153	41.72±0.15	41.15±0.15	41.95±0.15	...,E08
Haro11	87	7.92	14.7	4.3	H	0.604±0.153	43.28±0.15	42.54±0.15	43.76±0.15	S06,E08
UGC4393	35	7.95	67.15	33.72	J04	-2.195±0.153	40.76±0.15	40.98±0.15	41.29±0.15	J04,E08
POX4	52	7.96	16.71	14.6	G03	-1.238±0.153	41.28±0.15	40.64±0.15	41.59±0.16	G03,E08
Tol2138–405	246	8.01	15.5	11.14	A	-1.956±0.153	40.94±0.15	40.55±0.16	41.44±0.16	...,E08
NGC4861	15.2	8.01	119.5	118.5	G03	-2.739±0.153	40.01±0.15	39.53±0.15	40.30±0.15	G03,E08
Mrk206	25.4	8.04	18.1	10.45	A	> -1.371±0.153	41.89±0.15	41.42±0.15	42.19±0.15	...,E08
UM448	87	8.06	13.25	7.65	A	> -0.632±0.153	42.65±0.15	42.16±0.15	42.94±0.15	...,E08
SHOC391	106	8.06	8.49	4.9	A	-0.228±0.153	42.16±0.15	41.43±0.16	42.79±0.16	K04,E08
Mrk1450	20	7.99	13	5.1	C07	-0.586±0.108	41.92±0.15	41.22±0.16	42.19±0.16	C07,E08
SBS0335–052	57	7.25	15	4.1	C07	-0.721±0.108	41.33±0.15	41.05±0.32	42.26±0.32	C07,E08
VIIZw403	4.3	7.71	43.35	6.0	C07	-1.107±0.108	41.95±0.15	41.33±0.40	42.12±0.40	C07,E08
HS0822+3542	11	7.4	8	4.1	C07	-1.192±0.108	41.22±0.15	41.00±0.32	41.53±0.32	C07,E08
UGCA292	3.1	7.27	30	5.0	C07	-1.869±0.108	40.89±0.31	40.65±0.32	41.11±0.39	C07,E08

Table 1—Continued

Name <sup>a</sup>	D <sup>b</sup> (Mpc)	12+Log(O/H) <sup>c</sup>	R <sub>25</sub> <sup>d</sup> (")	R <sub>H<math>\alpha</math></sub> <sup>e</sup> (")	Ref. <sup>e</sup>	Log( $\Sigma_{SFR}$ ) <sup>f</sup> (M <sub>⊙</sub> yr <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{70}$ ) <sup>g</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{160}$ ) <sup>g</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{TIR}$ ) <sup>g</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Ref. <sup>h</sup>
UM461	13.4	7.8	9.08	5.0	C07	-0.892±0.108	41.57±0.15	40.69±0.21	41.89±0.21	C07,E08

<sup>a</sup>Galaxy name, as listed in Dale et al. (2007), Dale et al. (2009), or Engelbracht et al. (2008).

<sup>b</sup>Distance in Mpc, rescaled, where necessary, to H<sub>0</sub>=70 km s<sup>-1</sup> Mpc<sup>-1</sup>, as reported in Dale et al. (2007), Kennicutt et al. (2008), and Engelbracht et al. (2008).

<sup>c</sup>Oxygen abundances. For the SINGS galaxies (marked as ‘S’ in column 6), the reported number is an average of the ‘high’ and ‘low’ oxygen abundance values derived in Moustakas et al. (2009) (see discussion in Calzetti et al. 2007). For the LVL galaxies (marked as ‘L’ in column 6), the reported oxygen abundances are from Marble et al. (2010). For the other galaxies, oxygen abundances are from a variety of literature sources, as reported in Engelbracht et al. (2008).

<sup>d</sup>The optical radius R<sub>25</sub> in arcseconds, defined as D<sub>25</sub>/2. When available, the preferred source is the RC3 catalog (De Vaucouleurs et al. 1991), as reported in NED.

<sup>e</sup>The radius of the ionized gas emission, R<sub>H $\alpha$</sub> , in arcseconds (column 5), and source of the measurement (column 6). Measurements performed on images: S=SINGS (available at: [http://data.spitzer.caltech.edu/popular/sings/20070410\\_enhanced\\_v1/](http://data.spitzer.caltech.edu/popular/sings/20070410_enhanced_v1/); Kennicutt et al. 2009); L=LVL (Kennicutt et al. 2008); H=archival HST/ACS image; C99=Calzetti et al. (1999); C07=Calzetti et al. (2007, using P $\alpha$   $\lambda$ 1.8756  $\mu$ m images). Measurements performed on images retrieved from NED: E96=Evans et al. (1996); G01=Garcia-Barreto & Rosado (2001); G03=Gil de Paz, Madore, & Pevunova (2003); H04=Hunter & Elmegreen (2004); J00=Johnson et al. (2000); J04=James et al. (2004); K01=Koopman, Kenney & Young (2001); L99=Larsen & Richtler (1999). Measurements performed on published hard-copy images only: B02=Boselli & Gavazzi (2002); C06=Croft et al. (2006); F90=Fabbiano, Heckman & Keel (1990); H94=Hunter, van Weerden & Gallagher (1994); H03=Hernandez-Toledo et al. (2003); HT04=Hattori et al. (2004); H05=Hernandez et al. (2005); M99=Mendez, Esteban & Balcells (1999); R07=Robitaille et al. (2007). When no image is available, values adopted from Figure 1 are indicated with ‘A’.

<sup>f</sup>Star formation rate per unit area, in units of M<sub>⊙</sub> yr<sup>-1</sup> kpc<sup>-2</sup>, derived from equation 17, and dividing by the area calculated from the ionized gas emission radius of column 5. The  $\Sigma_{SFR}$  of galaxies for which the H $\alpha$  flux information is missing (see last column of Table) is marked as a lower limit.

<sup>g</sup>Luminosity at 70  $\mu$ m (column 8), 160  $\mu$ m (column 9), and integrated over the full 3–1,100  $\mu$ m wavelength range (TIR, column 10), normalized to the ionized gas emission area, in units of erg s<sup>-1</sup> kpc<sup>-2</sup>.

<sup>h</sup>Literature source of the H $\alpha$  (first reference) and infrared (second reference) flux data. B02=Boselli & Gavazzi (2002); C06=Croft et al. (2006); C07=Calzetti et al. (2007,  $\Sigma_{SFR}$  derived from extinction-corrected P $\alpha$  fluxes); D07=Dale et al. (2007); D09=Dale et al. (2009); E08=Engelbracht et al. (2008); H94=Hunter, van Weerden & Gallagher (1994); K87=Kennicutt et al. (1987); K04=Kniazev et al. (2004); K08=Kennicutt et al. (2008); K09=Kennicutt et al. (2009); L96=Lehnert & Heckman (1996); M99=Mendez, Esteban & Balcells (1999); M06=Moustakas & Kennicutt (2006); P06=Perez-Gonzalez et al. (2006); S06=Schmitt et al. (2006); W99=Wang, Heckman & Lehnert (1999).

Table 2. Properties of Luminous Infrared Galaxies.

Name <sup>a</sup>	D <sup>b</sup> (Mpc)	R <sub>25</sub> <sup>c</sup> ( $''$ )	R <sub>H<math>\alpha</math></sub> <sup>d</sup> ( $''$ )	Log( $\Sigma_{SFR}$ ) <sup>e</sup> (M <sub>⊙</sub> yr <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{24}$ ) <sup>f</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{70}$ ) <sup>f</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{160}$ ) <sup>f</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{TIR}$ ) <sup>f</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )	Log( $\Sigma_{3.6}$ ) <sup>g</sup> (erg s <sup>-1</sup> kpc <sup>-2</sup> )
NGC0023	63.8	62.7	10.	0.602±0.154	43.36± 0.15	43.93± 0.16	43.59± 0.16	44.19± 0.18	43.03± 0.15
MCG+12-02-001	68.9	6.6	4.9	1.179±0.154	43.92± 0.15	44.39± 0.16	43.92± 0.16	44.63± 0.18	...± ...
NGC0633	72.7	38.6	4.9	0.807±0.154	43.56± 0.15	43.97± 0.16	43.52± 0.16	44.24± 0.18	42.68± 0.15
UGC1845	66.4	36.1	6.2	0.428±0.154	43.19± 0.15	43.88± 0.16	43.49± 0.16	44.09± 0.18	42.60± 0.15
NGC1614	62.7	39.5	13.2	0.606±0.146	43.33± 0.15	43.49± 0.15	42.97± 0.15	43.85± 0.15	42.09± 0.15
UGC3351	65.2	47.6	12.7	-0.287±0.154	42.48± 0.15	43.44± 0.15	43.06± 0.16	43.62± 0.17	42.16± 0.15
NGC2369	47.1	106.5	17.5	-0.169±0.148	42.60± 0.15	43.30± 0.15	42.87± 0.15	43.50± 0.15	42.24± 0.15
NGC2388	61.9	30.0	13.3	0.045±0.154	42.80± 0.15	43.42± 0.16	42.94± 0.16	43.62± 0.18	42.01± 0.15
MCG+02-20-003	72.4	19.0	2.8	0.982±0.154	43.75± 0.15	44.51± 0.16	44.11± 0.16	44.70± 0.18	42.84± 0.15
NGC3110	78.7	46.5	20.	0.012±0.154	42.77± 0.15	43.50± 0.15	43.15± 0.16	43.72± 0.17	42.24± 0.15
NGC3256	35.5	114.1	23.3	0.434±0.146	43.17± 0.15	43.55± 0.15	43.02± 0.15	43.81± 0.15	42.17± 0.15
NGC3690/IC694	51.1	85.2	16.2	1.008±0.154	43.72± 0.15	44.02± 0.16	43.31± 0.16	44.28± 0.18	42.56± 0.15
ESO320-G030	40.4	67.2	8.8	0.448±0.147	43.22± 0.15	44.04± 0.15	43.53± 0.15	44.19± 0.15	42.51± 0.15
MCG-02-33-098	77.7	49.8	3.5	1.367±0.154	44.12± 0.15	44.45± 0.16	43.93± 0.16	44.73± 0.18	43.33± 0.17
IC860	63.3	16.5	3.3	1.088±0.154	43.85± 0.15	44.65± 0.15	43.94± 0.16	44.76± 0.17	42.83± 0.15
NGC5135	55.9	77.1	7.4	0.638±0.154	43.39± 0.15	43.97± 0.16	43.55± 0.16	44.20± 0.18	42.82± 0.15
NGC5653	65.1	52.2	12.7	-0.076±0.154	42.68± 0.15	43.31± 0.15	42.95± 0.16	43.55± 0.17	42.23± 0.15
NGC5734	63.5	45.4	13.0	-0.379±0.154	42.39± 0.15	43.20± 0.15	42.96± 0.16	43.45± 0.17	42.20± 0.15
IC4518	74.9	8.4	6.2	0.606±0.154	43.36± 0.15	43.79± 0.16	43.35± 0.16	44.06± 0.18	...± ...
Zw049.057	63.3	26.8	13.0	-0.269±0.154	42.50± 0.15	43.58± 0.15	43.00± 0.16	43.67± 0.17	41.52± 0.15
NGC5936	65.1	43.4	8.4	0.315±0.154	43.07± 0.15	43.58± 0.15	43.20± 0.16	43.84± 0.17	42.47± 0.15
IRAS17138-1017	81.2	22.5	5.1	0.929±0.154	43.67± 0.15	44.19± 0.16	43.71± 0.16	44.42± 0.18	...± ...
IC4687/6	79.4	38.7	5.8	1.046±0.154	43.78± 0.15	44.21± 0.16	43.76± 0.16	44.47± 0.18	42.65± 0.15
IC4734	73.5	38.7	5.6	0.626±0.154	43.38± 0.15	44.12± 0.15	43.75± 0.16	44.33± 0.17	42.69± 0.15
NGC6701	60.6	46.5	6.8	0.443±0.154	43.21± 0.15	43.82± 0.16	43.43± 0.16	44.05± 0.18	42.70± 0.15
NGC7130	70.7	45.4	11.7	0.210±0.154	42.96± 0.15	43.54± 0.16	43.12± 0.16	43.77± 0.18	42.27± 0.15
IC5179	50.0	70.4	16.5	-0.062±0.154	42.70± 0.15	43.34± 0.16	42.93± 0.16	43.56± 0.18	42.21± 0.15
NGC7591	70.2	58.5	5.9	0.559±0.154	43.32± 0.15	43.84± 0.16	43.43± 0.16	44.08± 0.18	42.73± 0.15
NGC7771	61.2	75.4	30.	0.080±0.154	42.83± 0.15	43.52± 0.15	43.19± 0.16	43.76± 0.17	42.51± 0.15

<sup>a</sup>Galaxy name, as listed in Alonso-Herrero et al. (2006).

<sup>b</sup>Distance in Mpc from Sanders et al. (2003) and Surace, Sanders & Mazzarella (2004), rescaled to H<sub>0</sub>=70 km/s/Mpc.

<sup>c</sup>The optical radius R<sub>25</sub> in arcseconds. When available, the preferred source is the RC3 catalog (De Vaucouleurs et al. 1991), as reported in NED.

<sup>d</sup>The radius of the ionized gas emission,  $R_{H\alpha}$ , in arcseconds. When available, the  $H\alpha$  sizes are from the images in Hattori et al. (2004); otherwise they are from Alonso-Herrero et al. (2006), using HST NICMOS  $P\alpha$   $\lambda 1.8756 \mu\text{m}$  images. For NGC1614, NGC3256, NGC3690, NGC5653, and Zw049.057 the data are from Alonso-Herrero et al. (2000, 2001, 2002).

<sup>e</sup>Star formation rate per unit area, in units of  $M_{\odot} \text{yr}^{-1} \text{kpc}^{-2}$ , derived from equation 17, and dividing by the area calculated from the ionized gas emission radius of column 4.

<sup>f</sup>Luminosity per unit area at  $24 \mu\text{m}$ ,  $70 \mu\text{m}$ ,  $160 \mu\text{m}$ , and integrated over the full  $8\text{--}1,000 \mu\text{m}$  wavelength range (TIR), in units of  $\text{erg s}^{-1} \text{kpc}^{-2}$ , using the ionized gas emission area. The luminosities in the MIPS bands are extrapolated from the IRAS measurements as described in section 3. Exceptions are NGC1614 and NGC3256, with MIPS data reported in Engelbracht et al. (2008), and NGC2369, ESO320-G030, and Zw049.057, with IR SED published in Rieke et al. (2009).

<sup>g</sup>Luminosity per unit area at  $3.6 \mu\text{m}$ , in units of  $\text{erg s}^{-1} \text{kpc}^{-2}$ , extrapolated from the  $K_S$ -band luminosity, as described in section 3.  $K_S$ -band luminosities are from 2MASS, as reported in NED, and corrected for the effect of both Galactic foreground (Schlegel, Finkbeiner & Davis 1998) and internal extinction.