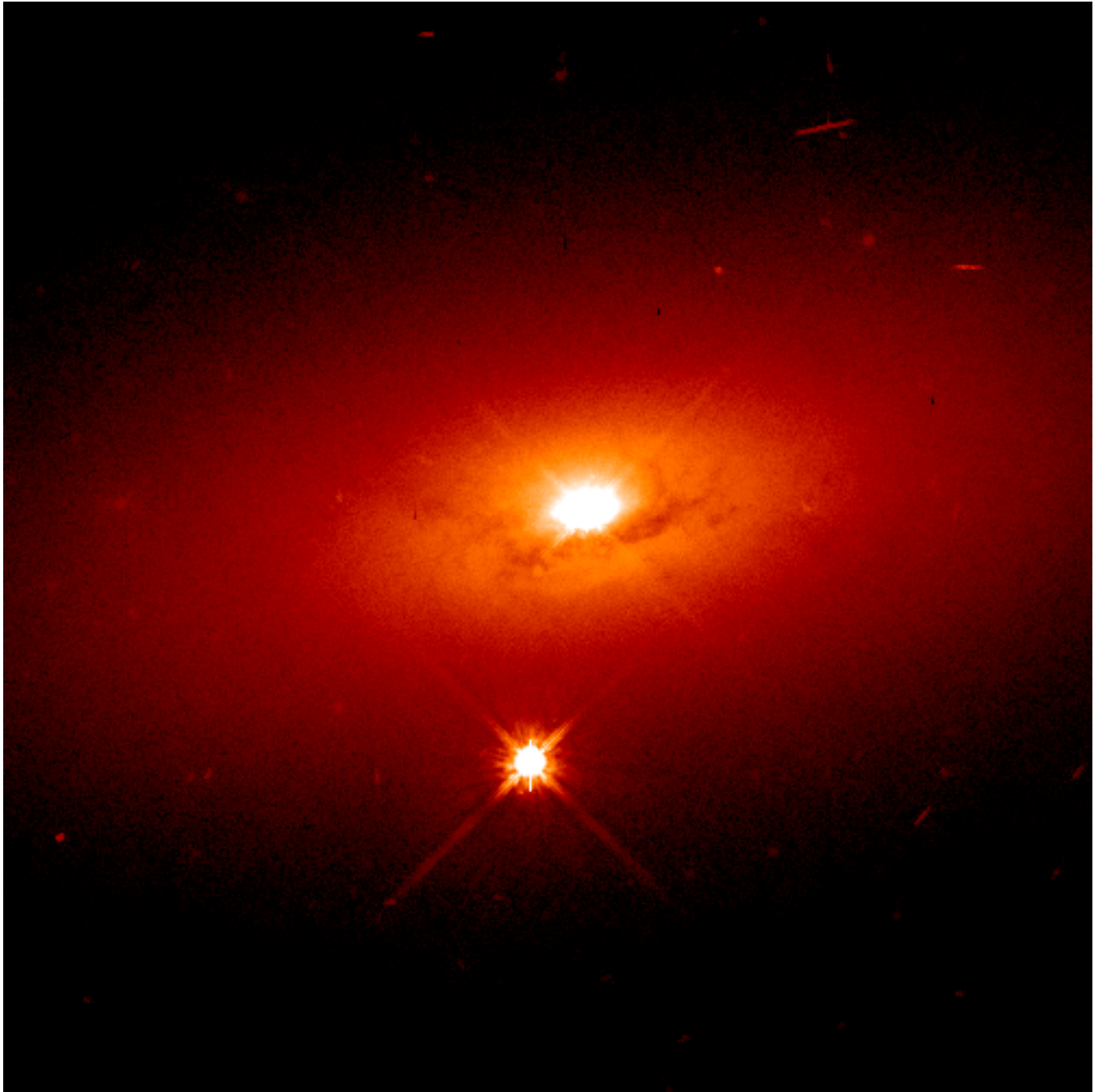


## Strong gravity and accreting black holes

- Finish how to get the masses of black holes
- The AGN Zoo
- Black Hole systems
  - The spectrum of accreting black holes
  - X-ray “reflection” from accretion disks
  - Strong gravity effects in the X-ray reflection spectrum

# Spectra of accreting black holes

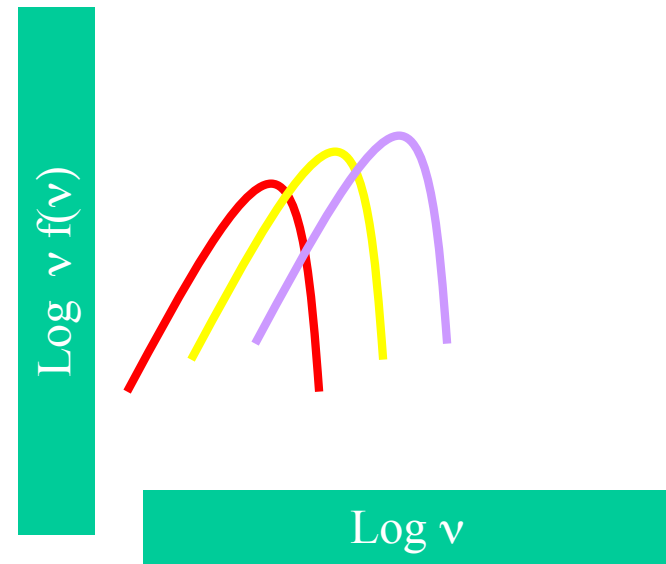
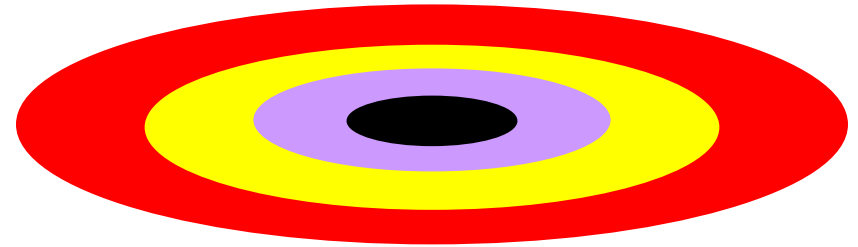




# Spectra of accretion flow: disc-

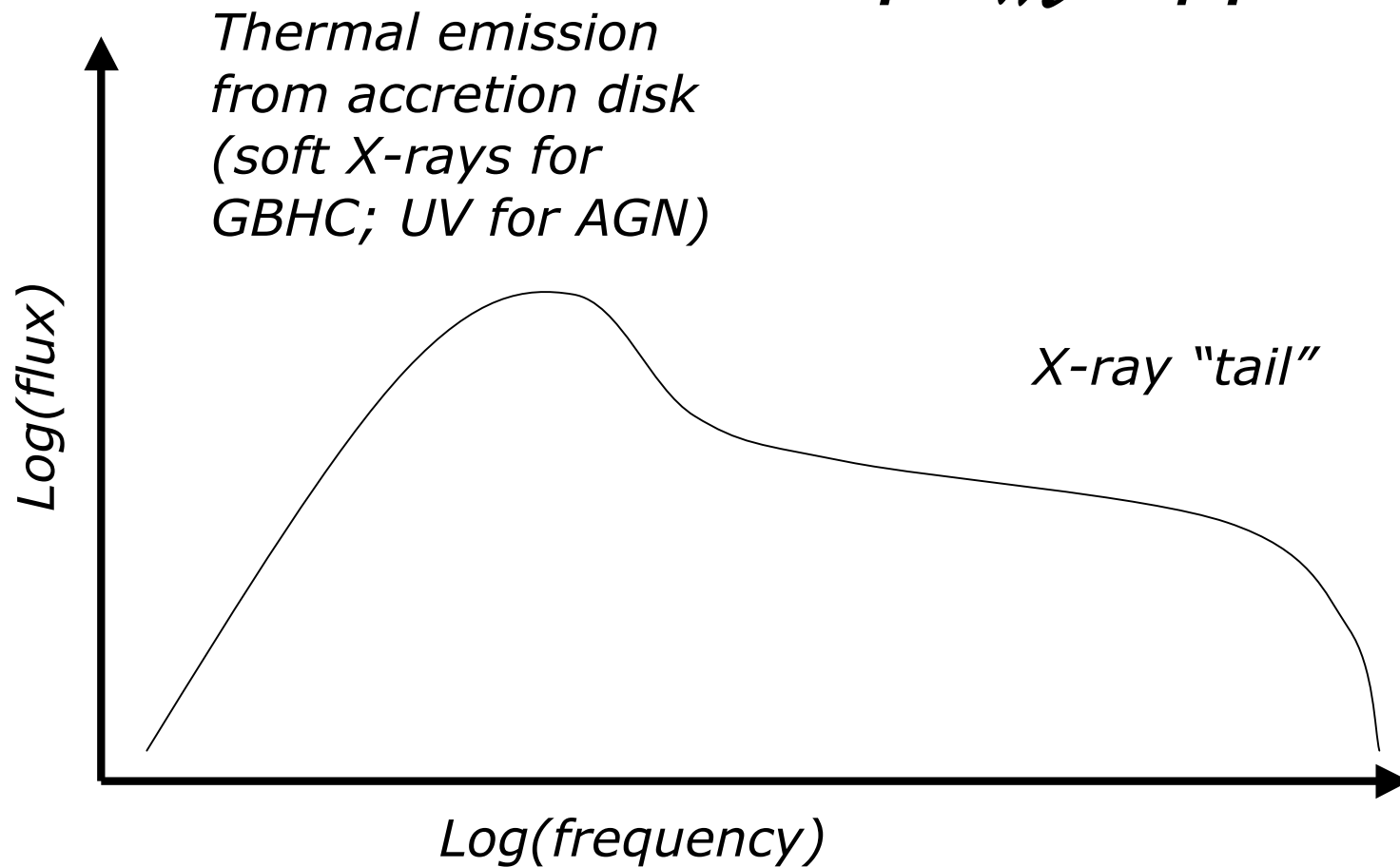
## C. Done

- Differential Keplerian rotation
- Viscosity : gravity  $\rightarrow$  heat
- Thermal emission:  $L = A\sigma T^4$
- Temperature increases inwards
- GR last stable orbit gives minimum radius  $R_{\text{ms}}$
- $a=0$ :  $T_{\text{max}} = (M/M_{\odot})^{-1/4} (L/L_{\text{Edd}})^{1/4}$ 
  - 1 keV ( $10^7$  K) for  $10 M_{\odot}$
  - 10 eV ( $10^5$  K) for  $10^8 M_{\odot}$
- $a=0.998$   $T_{\text{max}} \sim 2.2 T_{\text{max}} (a=0)$
- AGN: UV disc, ISM absorption, mass more uncertain. XRB...



# What Do Broad Band Spectra of Black Holes Look Like

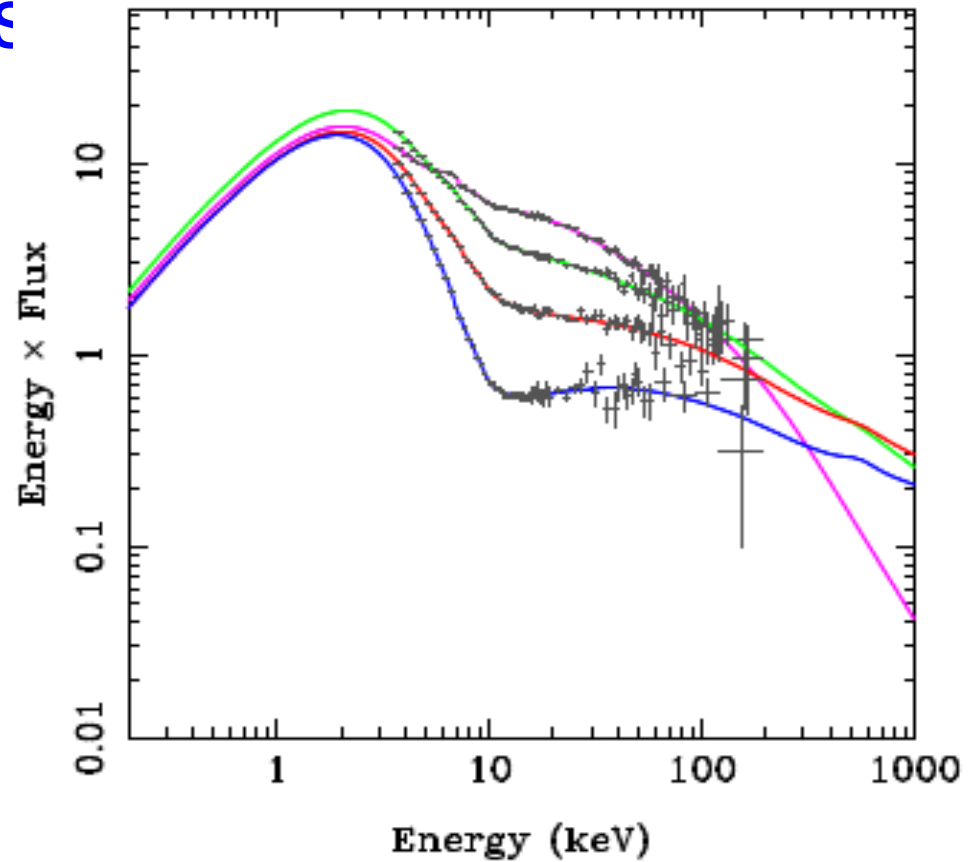
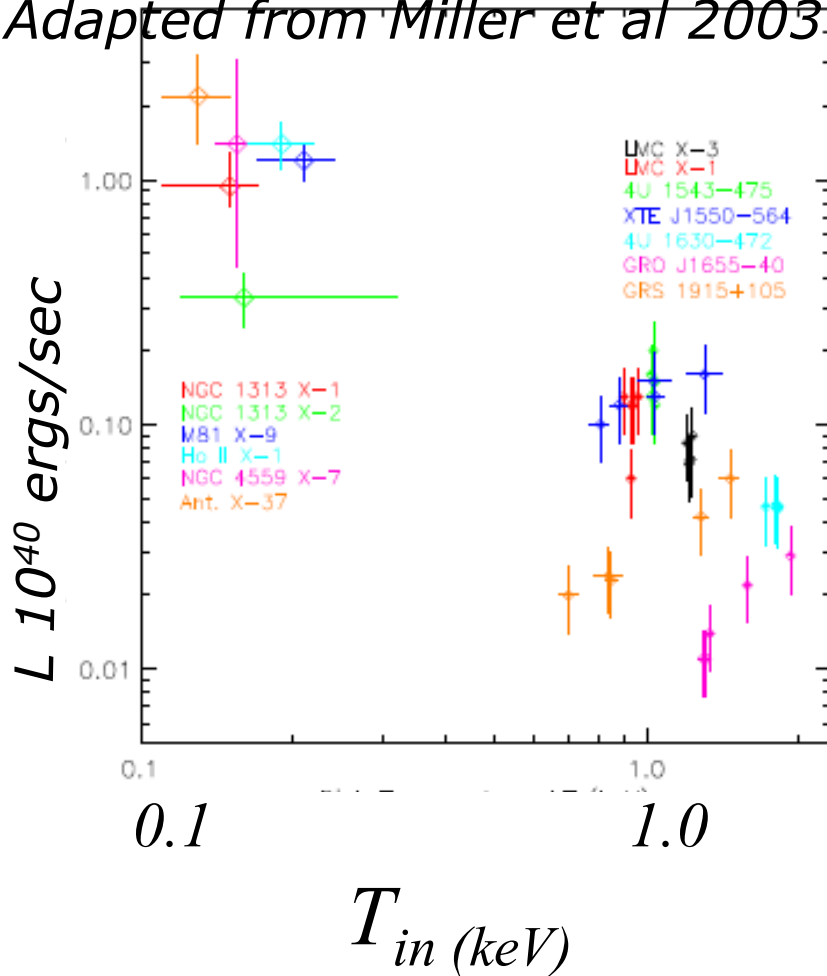
$$T \sim M^{1/4} M^{-1/4}$$



# Galactic Black Holes

- Relatively low mass and so the disk are 'hot'

Adapted from Miller et al 2003



$$L \sim 7 \times 10^{38} f (M_{10})^2 (T_{in})^4 \text{ erg/s}$$

Where  $M_{10}$  is the mass in units of  $10 M_{\odot}$  and  $T_{in}$  is in keV

***f* is a factor taking some physics into account**

## Derivation (See Rosswog and Bruggen sec 8.4)

- Derivation of previous eq
- $L=2\pi R_{in}^2 f(\cos i)^{-1}$ ;  $f$  is the flux from the surface of the disk,  $R$  is the radius
- Using the black body law

$$L=4\pi\sigma R^2 T_{in}^4 \quad \sigma \text{ is the Stefan- Boltmann constant}$$

In fitting the spectrum  $T_{in}$  is directly observable

We can thus take the 2 equations to get the innermost radius

$$R_{in} = \text{sqrt}(L/4\pi\sigma T_{in}^4) \text{ and}$$

$$T_{in} \sim 3M_{10}^{-1/4} \text{keV}$$

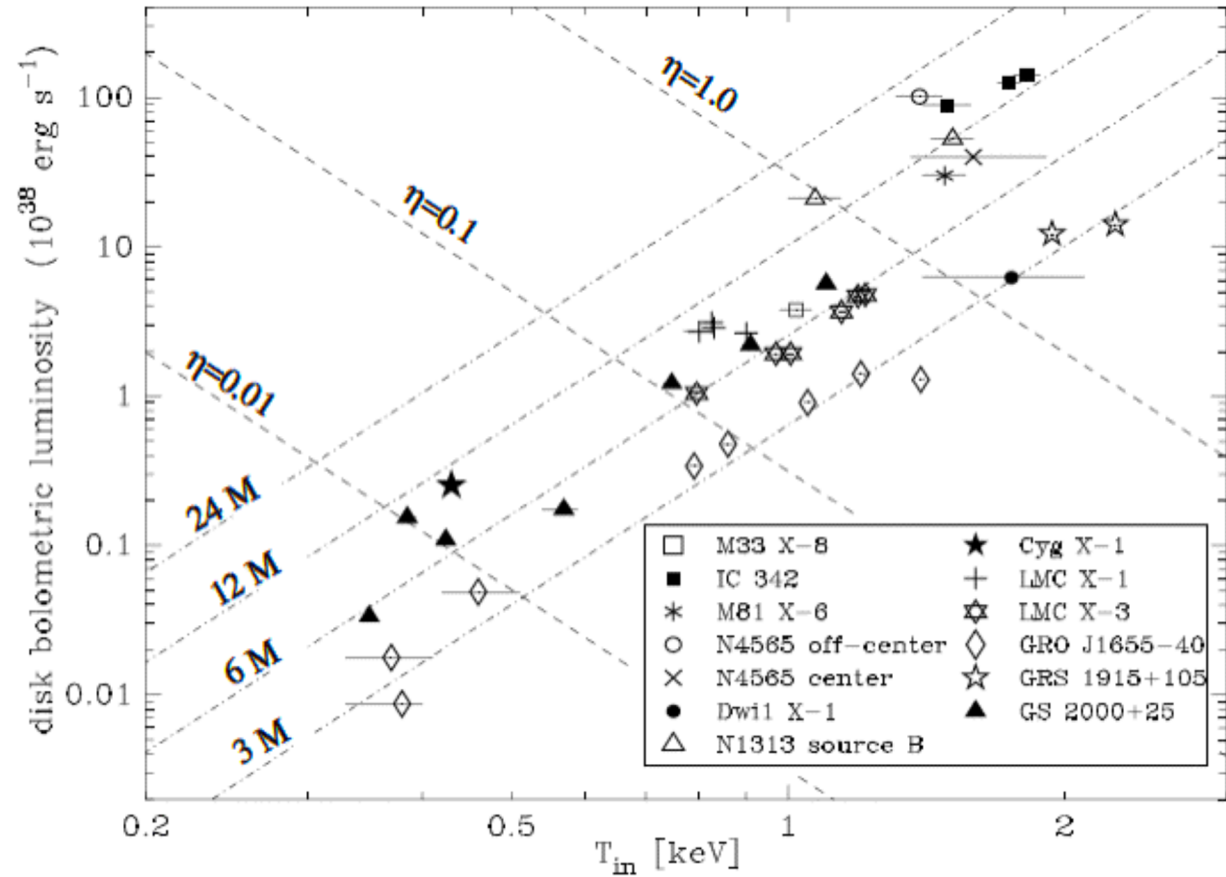
$$T(r) = 6.3 \times 10^5 \mathcal{M}_{Edd}^{1/4} M_8^{-1/4} (r/r_s)^{-3/4}$$

( $\mathcal{M}_{Edd}$  is the accretion rate in Eddington units,  $T=T_{in}$  for  $r=r_s$ )

# Real Objects

- Amazingly data for galactic black holes agrees with the simple theory

Makishima et al 2000

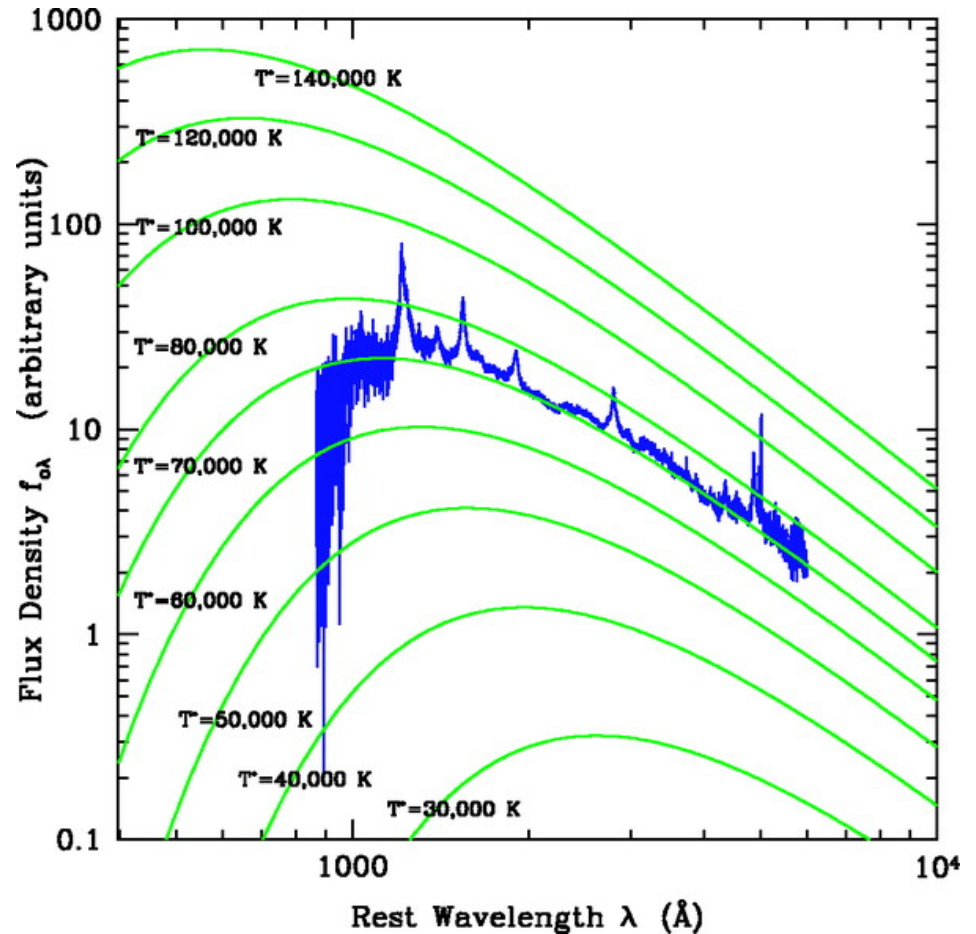


$$L_{bol} = \eta L_E \propto \eta M ,$$



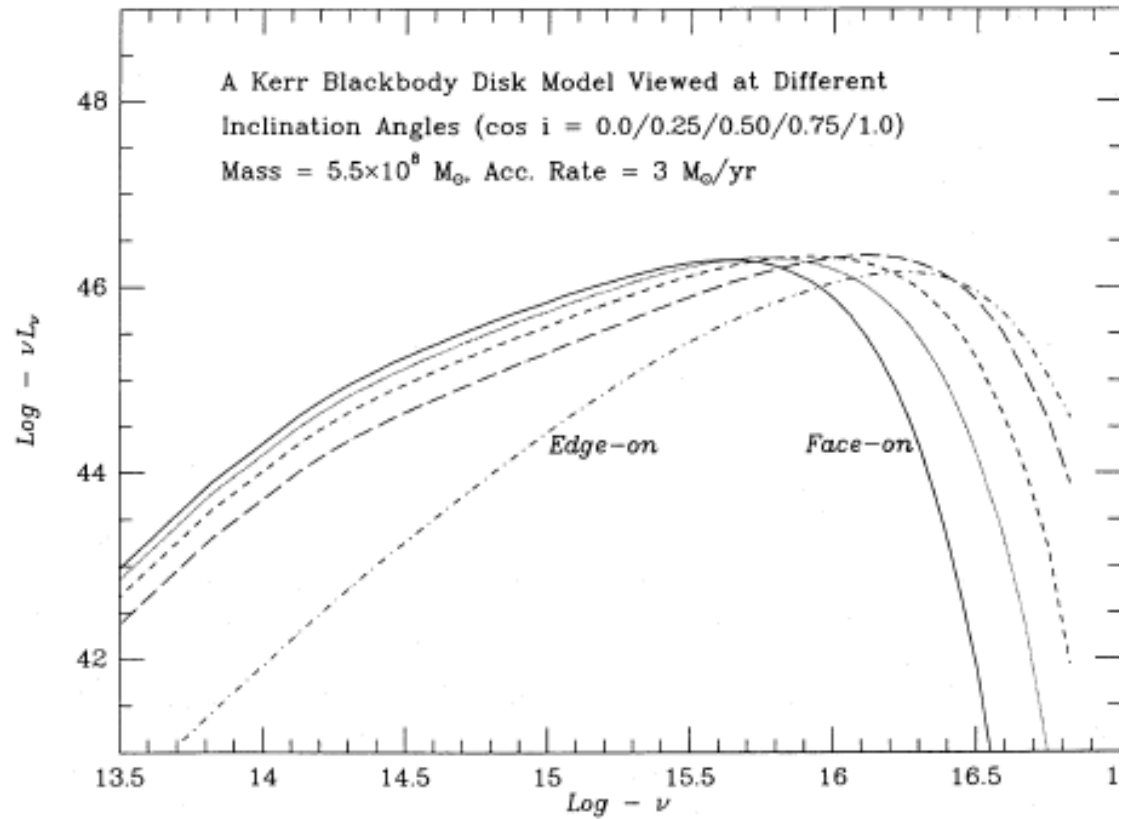
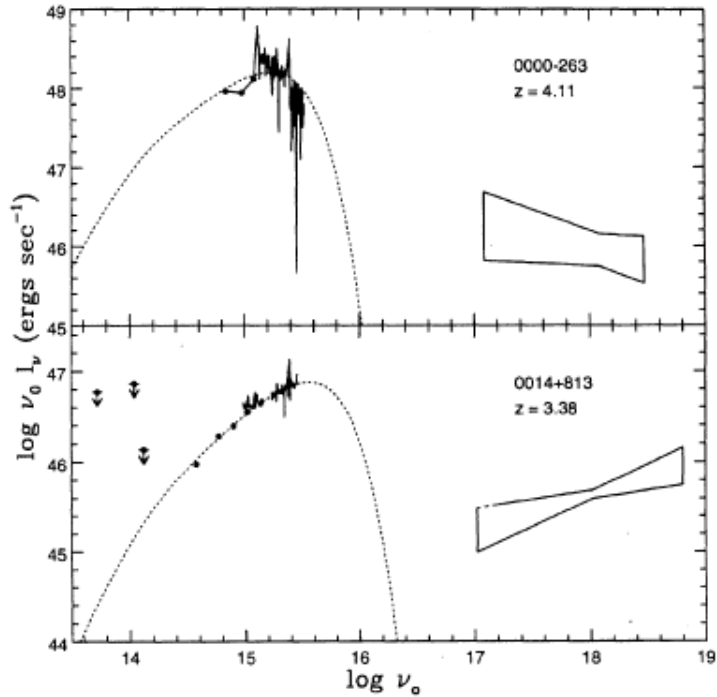
# AGN

- AGN are very massive and so the predicted spectrum of the accretion disk is 'cool'
- $T \sim 8 \times 10^4$  K for an Eddington limited  $M \sim 10^8 M_{\odot}$  black hole



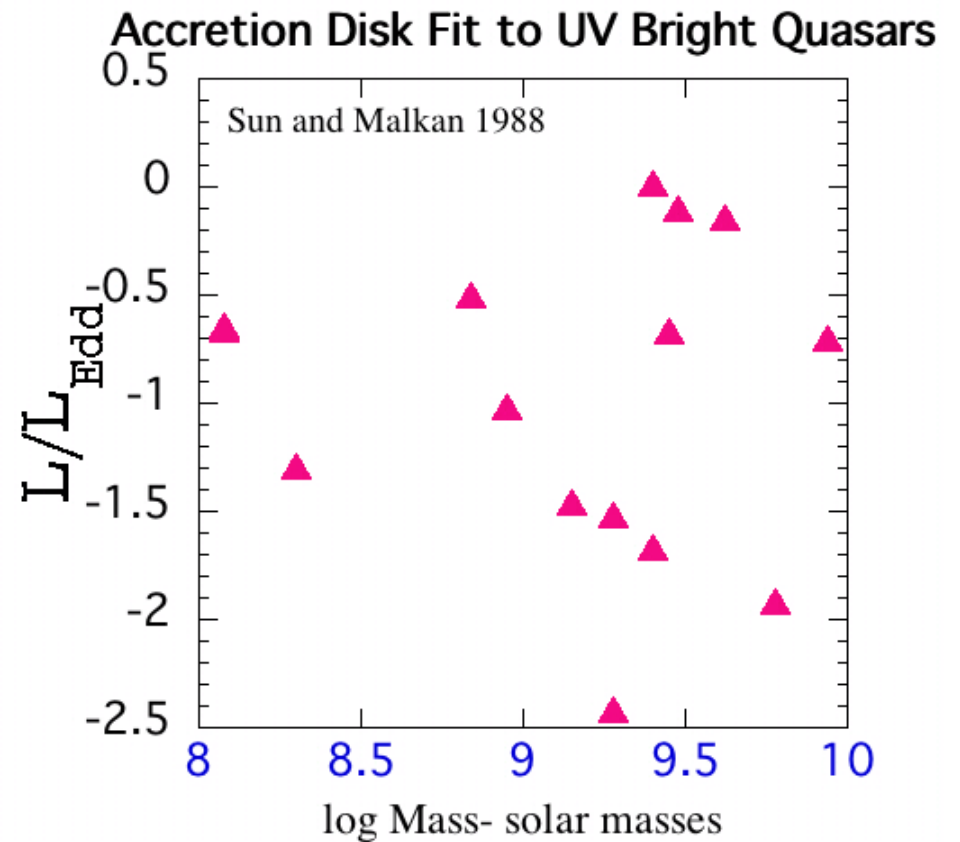
# Malkan and Sun 1989

- Can Fit AGN UV-optical data with accretion disk models

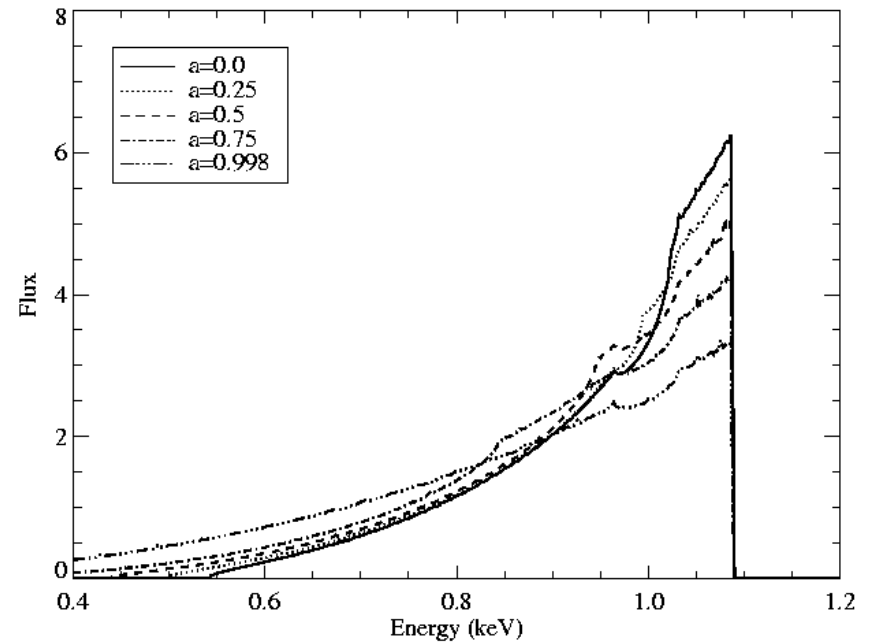
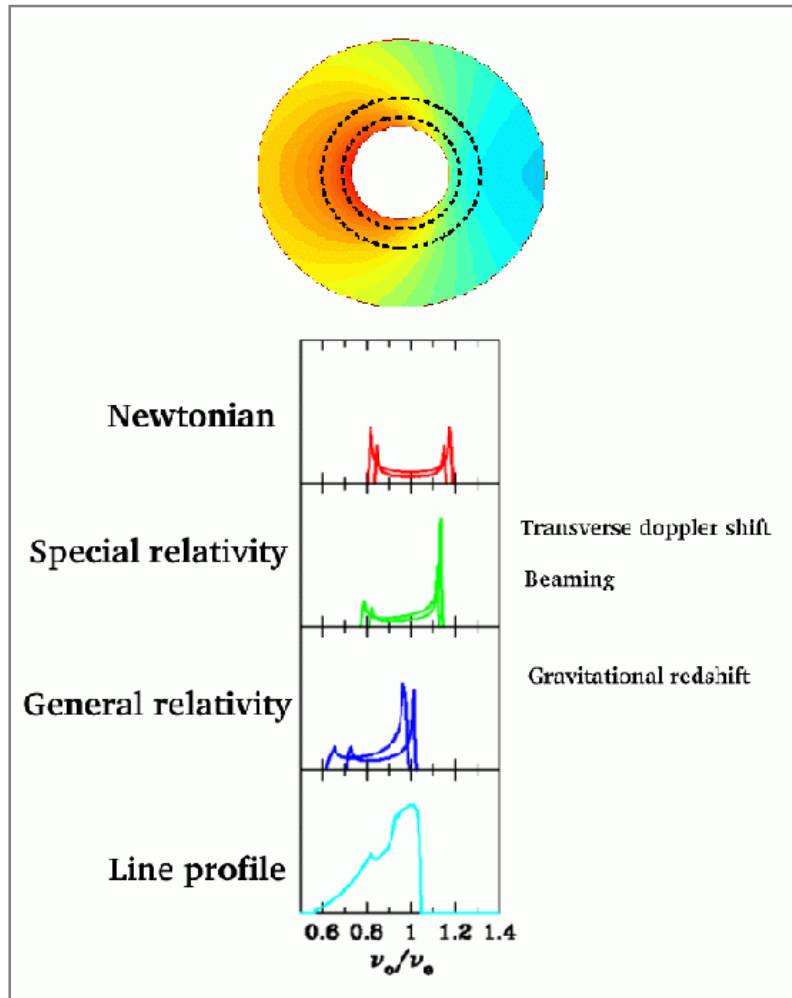


# Fitted Parameters for UV Disk Fits

- Results are 'reasonable' but not unique
- Now have independent mass estimates- results can be checked
- Find that values are not quite right- need more complex accretion disk models (surface is not BB relativistic effects)



# Relativistic effects imprint characteristic profile on the emission line...

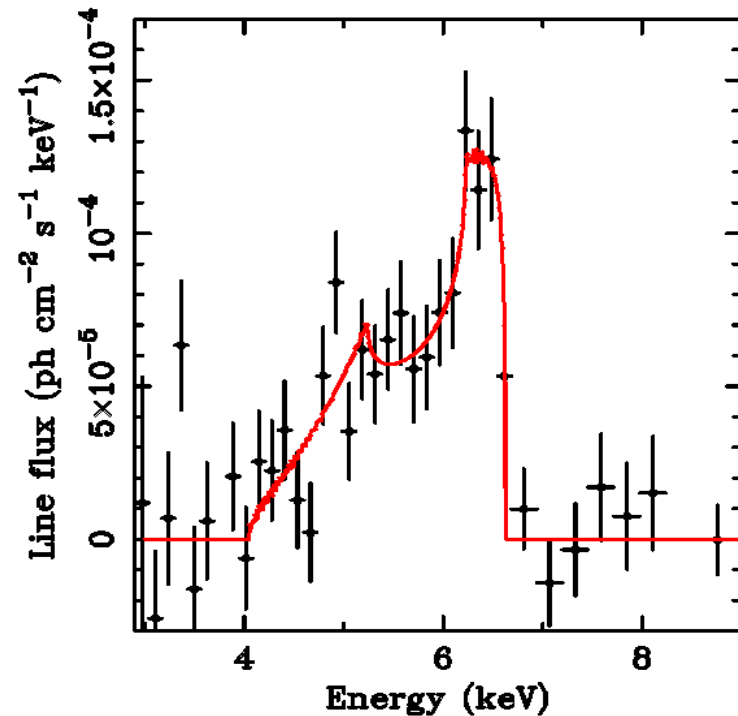


*Theoretical line profiles  
[Laura Brenneman]*

**Andy Young**

# Observations of relativistic emission lines

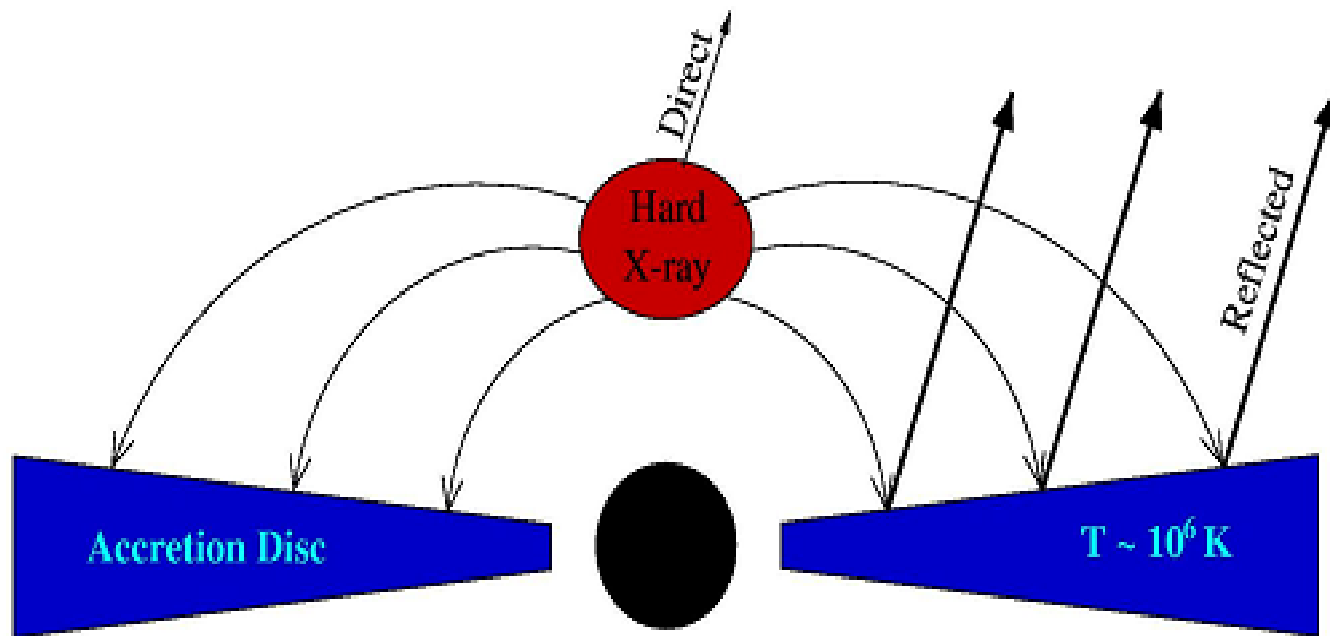
- First seen in 1994 with ASCA observatory
- 5 day observation of Seyfert-1 galaxy MCG-6-30-15
- Needed long observation to collect enough photons to form detailed spectrum



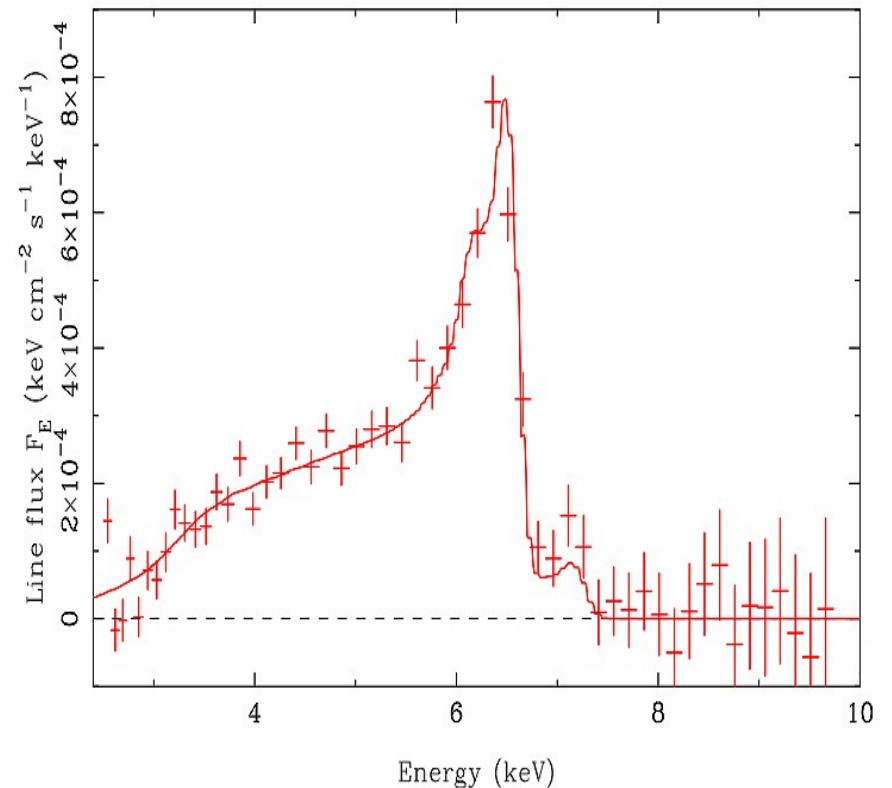
Power-law continuum subtracted  
ASCA: Tanaka et al. (1995)

# Relativistic Effects

- Light rays are bent by strong gravity- making the geometry rather complicated
- Do not know 'where' x-ray source is - try to use data to figure it out

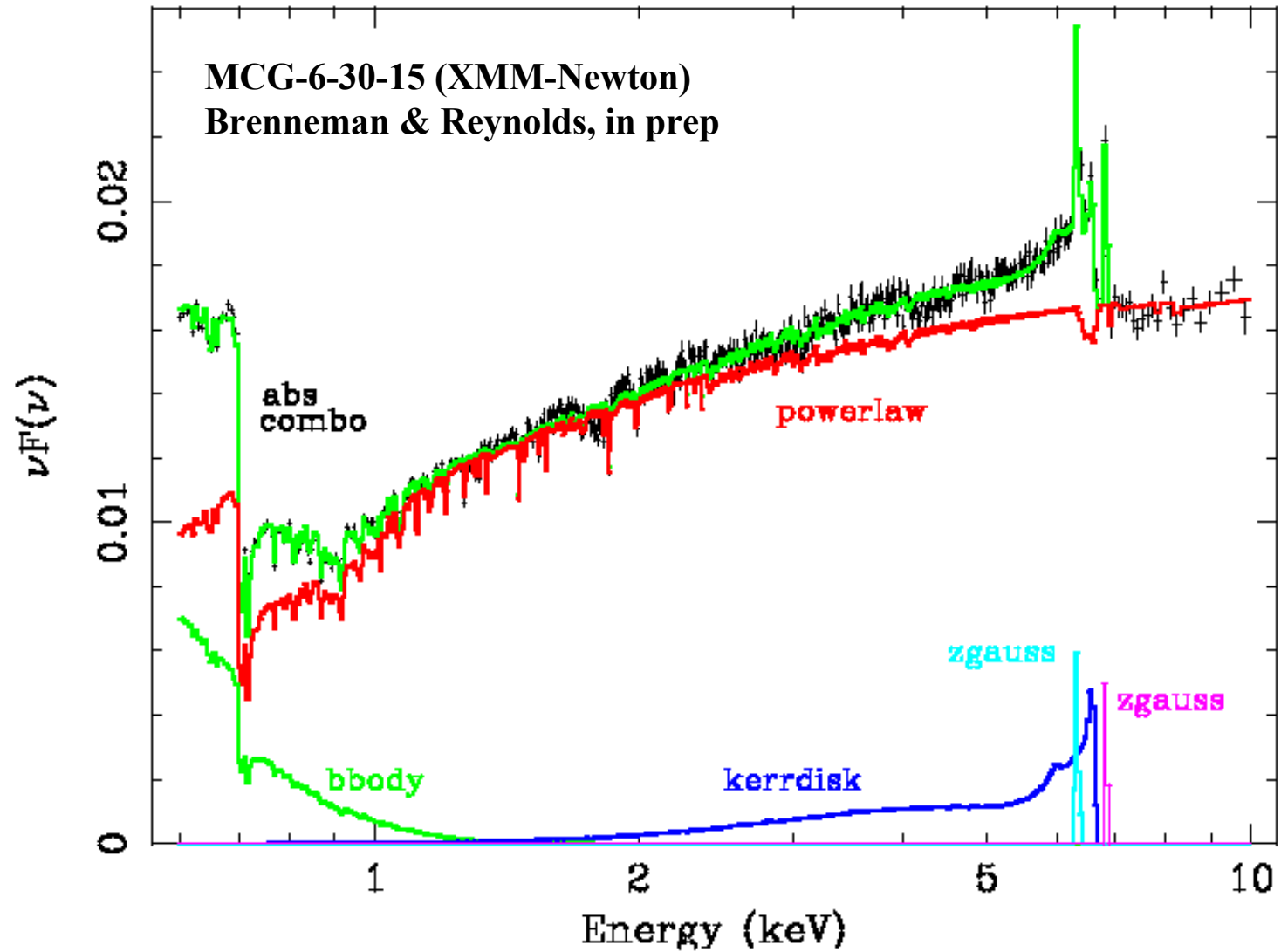


- Modern XMM-Newton observations
- Confirm relativistic line with extreme redshifts
- If no line emission from within ISCO, need to invoke spinning black hole to get strong enough redshifting



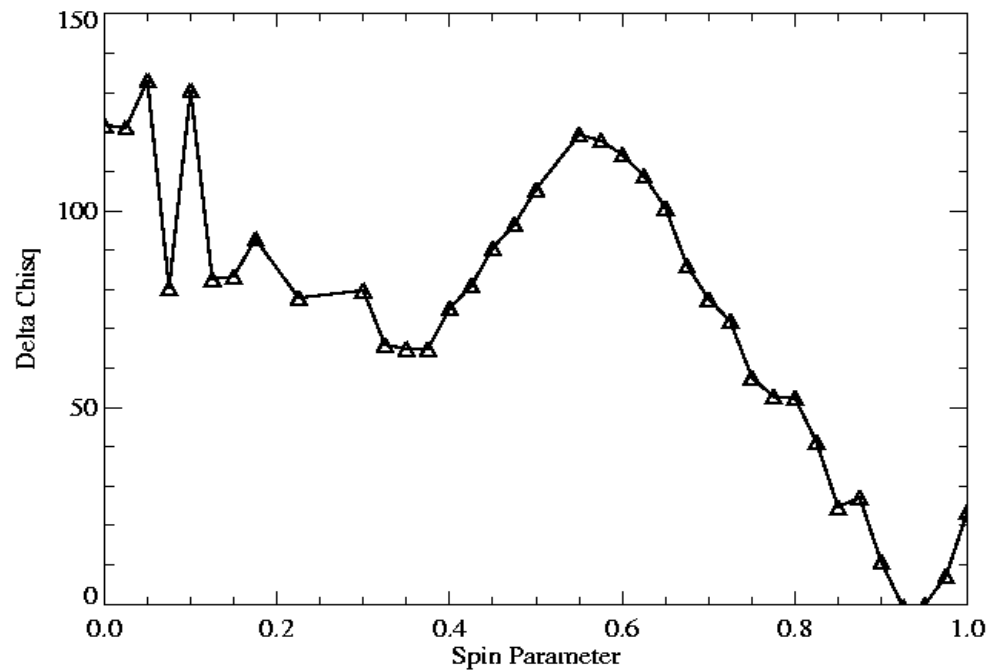
Power-law continuum subtracted  
XMM: Fabian et al. (2002)

Spectra are quite complex...





- Applied models to long (350ks) XMM dataset for MCG-6-30-15
  - Data strongly prefers rapidly spinning BH solution
  - $a \sim 0.93$



**L.Brenneman**

