

## NEW TOPIC- Star Formation MBW ch 9

- One of the most important processes for galaxy formation and evolution
- Big questions
  - When and how does star formation occur ?
  - How is it related to the evolution of galaxy properties?
  - What are the physical processes that drive star formation ?
    - star formation occurs (at least in spirals at low  $z$ ) almost exclusively associated with molecular clouds
      - what is the rate at which stars form in this cloud
      - what mass fraction of the cloud forms stars
      - what controls the IMF?
- for a review see

### **THE CURRENT STATUS OF GALAXY FORMATION**

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[https://ned.ipac.caltech.edu/level5/March12/Silk/Silk\\_contents.html](https://ned.ipac.caltech.edu/level5/March12/Silk/Silk_contents.html)

## Star Formation in Spirals

- This is an enormous subject- lots of recent work (see Kennicutt 1989 and Kennicutt and Evans 2012 for reviews)
- Broadly.. Observations of nearby galaxies have shown, over a broad range of galactic environments and metallicities, that star formation occurs only in the molecular phase of the interstellar medium (ISM).
  - Star formation is inextricably linked to the molecular clouds
  - Theoretical models show that this association results from the correlation between chemical phase, shielding, and temperature.
  - See MWB sec 9.1.-9.3 for a discussion
- Interstellar gas converts from atomic to molecular only in regions that are well shielded from interstellar ultraviolet (UV) photons, and since UV photons are also the dominant source of interstellar heating, only in these shielded regions does the gas become cold enough to be subject to Jeans instability (Krumholz 2012)

# Star Formation Occurs Primarily in Spirals

- HST imaging of low-moderate redshift universe shows that star-forming galaxies at all redshifts are *dominated* by disks, while passive (non-starforming or non-active) galaxies have spheroidal structures
  - (Eales et al 2015)- estimate that  $\approx 83\%$  of the stellar mass-density formed over the history of the Universe occurred in LTGs (jargon, late type galaxies, aka spirals)
- However since  $\sim 50\%$  of all stellar mass lies in passive galaxies- need either to transform spirals into E's or merge them

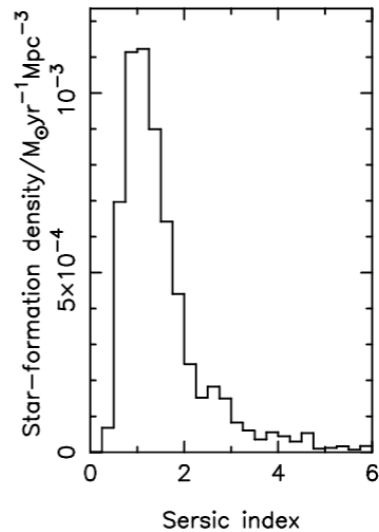


Figure 1. Star-formation rate per unit comoving volume in the Universe today as a function of Sérsic index.

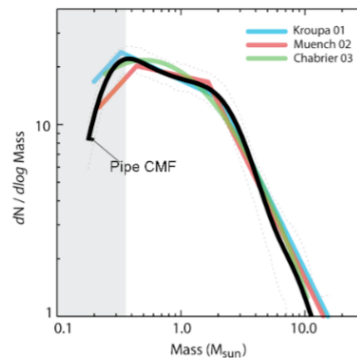
## Star Formation in Spirals

In the MW and other well studied nearby galaxies S occurs mostly in

(Giant molecular clouds (GMCs, which are predominantly molecular, gravitationally bound clouds with typical masses  $\sim 10^5 - 10^6 M_\odot$ )- but GMC formation is a local, not a global process

There is a strong correlation between the mass spectr of molecular clouds and the stellar mass spectrum (Lada 2015 [arXiv:1508.02711](#))

- Observationally one uses CO as a tracer for  $H_2$  (not perfect but the best we have right now Bolatto, Wolfire, and Leroy ARA&A 2013)). This is time consuming but lots of work has been done (Leroy et al 2008)



3 colored lines  
is mass spectrum of  
stars from 3 models  
and black line is mass  
spectrum of molecular  
clouds from a particular  
molecular cloud  
complex

# Star Formation- How to Measure It

The physics of star formation (what processes produce stars) and the astrophysics (where and when were the stars produced) are two of the dominant issues in astrophysics at present-

- Stars form from dense, cold gas either in disks or in gas that is violently shock compressed (in mergers)

Current SF can be estimated from a variety of techniques

- $H\alpha$  observations, which gives the number of ionizing photons if one assumes that all ionizing photons are used and eventually re-emitted - ionizing photons are almost exclusively emitted by massive (hot) stars which have short lifetimes; so the effects of dust can be large
- far-IR flux - this assumes that a constant fraction of the emitted stellar energy is absorbed by dust
- radio continuum emission - this statistically correlated very well with the IR radiation- physics is complex since radio emission comes from synchrotron radiation from relativistic electrons+ thermal bremsstrahlung from hot gas
- far-UV flux (- which is primarily emitted by young (hot) stars- but older /less massive than those responsible for  $H\alpha$ )
- X-ray emission- produced by 'high mass' x-ray binaries (a Neutron star or black hole with a massive companion )

## How to Normalize SFR

- Since essentially all techniques measure the total (or ionizing) luminosity of massive stars we need to transform to ALL the stars
- Use the IMF (Initial mass function; **please read sec 9.6 of MBW** but not 9.6.2)
- For Kroupa IMF
  - $\Psi(M) \sim M^{-1.4}$   $0.1 < M_{\odot} < 1$
  - $\Psi(M) \sim M^{2.5}$   $1 < M_{\odot} < 100$
- Integrate  $\Psi$  from  $10-100M_{\odot}$  get 0.16 of all the mass (correction factor)- these are the stars which have short lifetimes and are hot and thus produce the signatures of star formation. Formation of low mass stars can only be detected in MW and Magellanic clouds

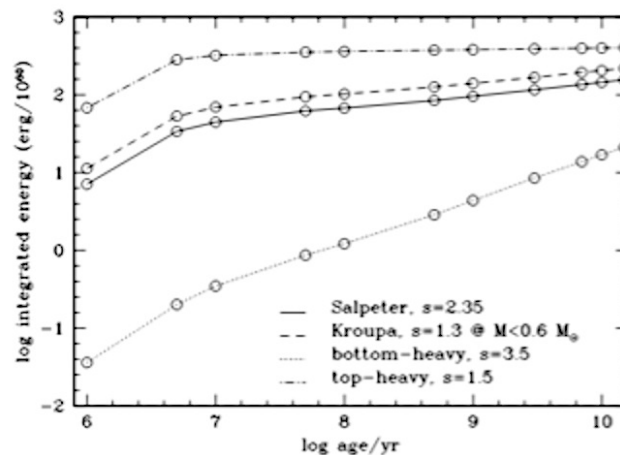
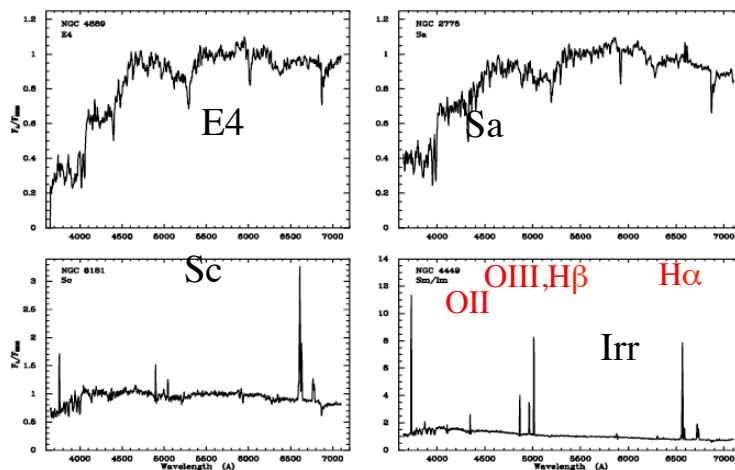


Figure 3. The total energy that has been emitted by a given age for a single stellar population with a mass of  $10^{11} M_{\odot}$ . We

- How to correct various indicators
- $H\alpha$ : emitted by gas ionized by stars with  $T_{\text{eff}} > \sim 20,000\text{K}$  ( $M > 10M_{\odot}$ ) which emit photons that can ionize Hydrogen ( $E_{\text{ioniz}} = 13.6\text{eV}$ ) -  $t < 20\text{Myrs}$
- IR Continuum- UV light absorbed by dust
- UV continuum- direct signature of massive, young stars

## Importance of Emission Lines

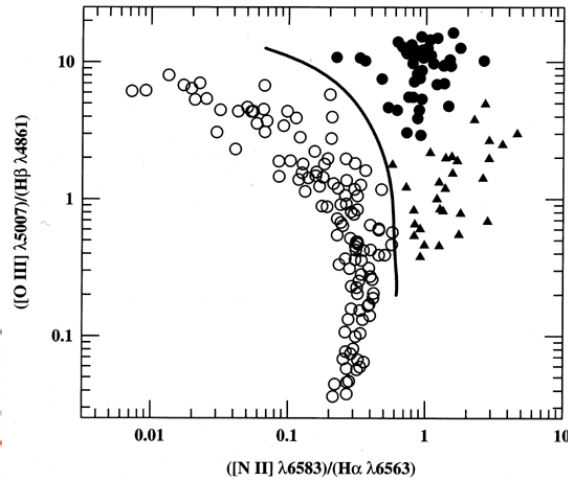
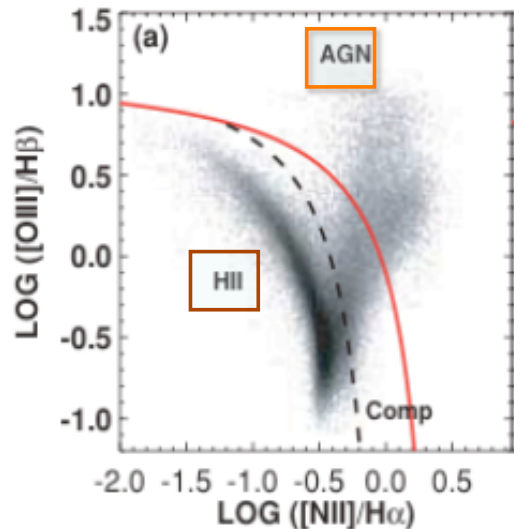
- As one moves on the Hubble sequence the galaxy spectra get more and more emission line dominated and relative prominence of lines changes
- Thus many authors use  $H\alpha$  or  $OII$  as SFR indicators
  - these are strong optical lines produced by gas ionized by hot stars ( $OIII$  is also produced by active galaxies and so it is often difficult to separate AGN from star formation)



Kennicutt 1998

Digression-From spectroscopy how does one classify a galaxy as star forming or an AGN??

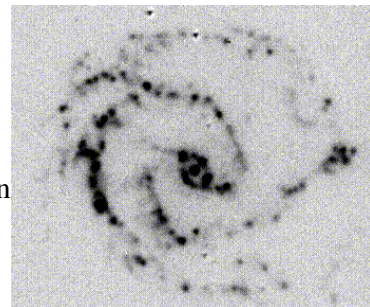
Observe strong lines to make life easier- but these are not necessarily to most diagnostic.



Different lines have different dependences on temperature excitation mechanism (collisions, photoionization)  
Ratios of certain lines (chosen to be close in wavelength do dust is not an issue)  
AGN have 'harder' radiation field (higher UV/optical) and collisional excitation less important than in star forming regions.

## How to Determine SFR from Observables-Hα or Hβ see 10.3.7,10.3.8

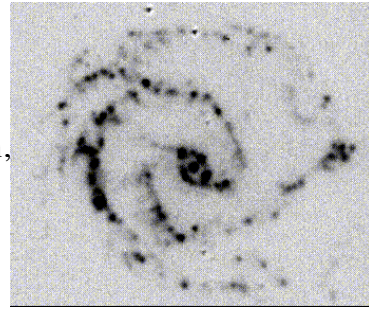
- The strength of the emission lines is the convolution of the number of ionizing photons, the fraction of them that are absorbed and the physical conditions of the gas.
- Simplifying assumptions: gas of constant temperature, given IMF, gas is internally dust free, Case B (optically thick to ionizing continuum)(Hα/Hβ=2.9) and an analytic approximation to the SFR
  - Hα only comes from ionized gas (HII regions)- very non-uniform images (pearls on a string)
- For one type of star (O7) one can calculate the number of Hα photons=  $10^{38}$ ph/sec
- Using stellar models and the IMF one ends up with  $SFR(M_{\odot}/yr)=L(H\alpha)/7 \times 10^{41}$  ergs/sec for  $M > 10M_{\odot}$  stars or
- $SFR(M_{\odot}/yr)=L(H\alpha)/1.1 \times 10^{41}$  ergs/sec for **all** stars
- while this seems great, have to worry about dust, the age of the population- the equation assumes a **zero age** IMF. The older the population is, the less Hα there is- harder to see how much star formation occurred if it has turned off and the system is more than 20Myrs old.



Hα image of a star forming galaxy

## How to Determine SFR from Observables- $H\alpha$ or $H\beta$

- Young, massive stars produce copious amounts of ionizing photons that ionize the surrounding gas. Hydrogen recombination cascades produce line emission, including the well-known Balmer series lines of  $H\alpha$  (6563Å) and  $H\beta$  (4861Å), which are strong.
- Only stars more massive than  $20M_{\odot}$  produce an ionizing photon flux.
- In a stellar population formed through an instantaneous burst with a Kroupa IMF the ionizing photon flux decreases by two orders of magnitude between 5Myr and 10Myr after the burst.
- So  $H\alpha$  measures the '**instantaneous**' star formation rate

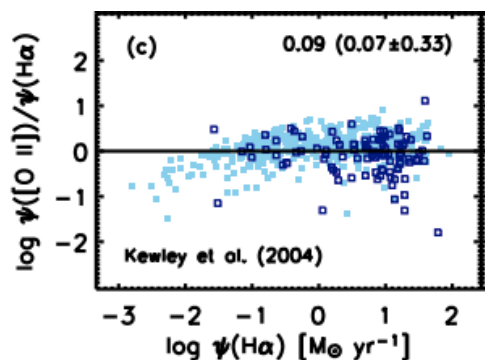


$H\alpha$  image of a star forming galaxy

• <http://www.astr.ua.edu/keel/galaxies/sfr.html>

## How to Determine SFR from Observables-[OII]

- [OII] (a forbidden line, collisionally de-excited in dense gas) is the next most prominent line and is visible until  $z \sim 1.4$  from the ground ( $H\alpha$  is only visible to  $z \sim 0.4$ )
- Calibrate it empirically using  $H\alpha$  since its luminosity is not directly coupled to the ionizing continuum (it is collisionally excited, not a cascade from photoionization) - but fairly wide variation in  $H\alpha$  / [OII] makes it noisier.



Ratio of SFR from [OII] to  $H\alpha$  rate vs  $H\alpha$  rate (Moustakas 2006)

## Summary and Look Forward

- The star formation rates is determined using many different indicators.
- The most important of are
  - far infrared emission tracing deeply embedded star formation
  - H $\alpha$  emission tracing H II regions;
  - and far ultraviolet emission tracing young, massive stars that have dispersed their natal gas and dust.
  - Radio emission tracing relativistic particles created by SF processes (e.g. supernova)
- Molecular hydrogen surface density correlates linearly with star formation rate -HI seems not to matter-

## Star Formation

- One of the most important processes for galaxy formation and evolution
- What are the general conditions for star formation?
  - in the low  $z$  universe star formation in spirals occurs mostly in molecular clouds
  - in ellipticals it is not understood; but is it clear that in some ellipticals stars are forming now.
  - special class of star forming galaxies- star bursts
- General scenario gas cloud collapses, fragments, stars form (somehow).

## SFR indicators

- SFR indicators are derived across the full electromagnetic spectrum, from the X-ray, through the ultraviolet (UV), via the optical and infrared (IR), radio, and using both continuum and line emission (review Kennicutt 1998, Kennicutt & Evans (2012)).
- The importance of these indicators change over cosmic time: it seems that most of the star formation at redshift  $z \sim 1-3$  was enshrouded in dust but at  $z > 3$  dust was much less important.

## UV

- The youngest stellar populations emit the bulk of their energy in the rest frame UV ( $< 0.3\mu$ ); **in the absence of dust attenuation, this is the wavelength range ‘par excellence’ to investigate star formation in galaxies over timescales of  $\approx 10-300$  Myr,**
- both O and B stars are brighter in the UV than at longer wavelengths.
  - the lifetime of an O6 star is  $\sim 6$  Myr, and that of a B8 star is  $\sim 350$  Myr.

The luminosity ratio of a O6 to B8 star at  $0.16 \mu$  is  $\sim 90$ , but, weighting by a Saltpeter IMF SSP for every O6 star formed, 150 B8 stars are formed.

Thus, at age zero, the UV emission from the collective contribution of B8 stars is comparable to that of O6 stars. And since B8 stars live a lot longer they dominate the UV flux on longer timescales.

(Calzetti 2012)

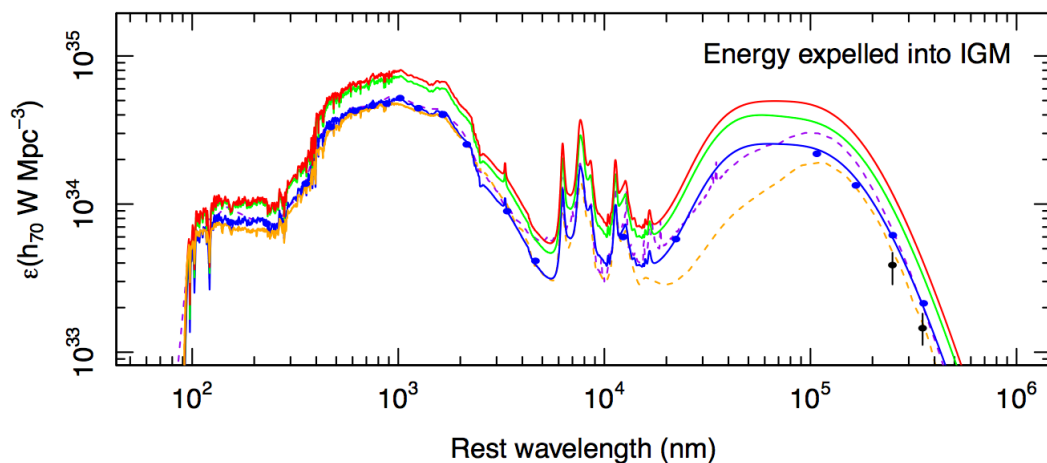


# UV Continuum

- in principle great- direct measure of total luminosity of young massive stars.
- Three big problems
  - DUST- UV extinction is much larger than in optical - light that is absorbed is re-emitted in the IR -the most active and luminous systems are also richer in dust, implying that they require more substantial corrections for the effects of dust attenuation; (MBW-10.3.8(b))
    - effects of dust are **BIG**-  $A_V = 0.9$  produces a factor ten reduction in the UV continuum at  $1300\text{\AA}$  (see MBW pg 479, S+G pg 33-34 for discussion of reddening- more later in lectures on dust)
  - Observations show that at 'low' SFR dust is not a big effect, at high values critical
  - at low redshift must observe from space – e.g. UV does not get thru the atmosphere
  - VERY sensitive to IMF- at best can only constrain 15% of all the stars forming
  - For a Kroupa IMF with constant star formation
 
$$\text{SFR}(\text{UV})M_{\odot}/\text{yr} = 3.0 \times 10^{-47} L_{\text{UV}}(\text{ergs/sec})(912-3000\text{\AA})$$

# IR Continuum

- Direct observations show that  **$\sim 1/2$  of total galaxy light in spirals appears in IR**
- This is thermal emission emitted by dust as a grey body



**Figure 25.** The energy originating (i.e., unattenuated, top), and emanating (i.e., attenuated following dust reprocessing, lower) at intervals equivalent to 0.75, 1.5 and 2.25 Gyr lookback time. The data are normalised to the energy output per  $\text{Mpc}^3$  for  $H_o = 70\text{km/s/Mpc}$ . The data show clear trends in the evolution of the total energy output over this timeline.

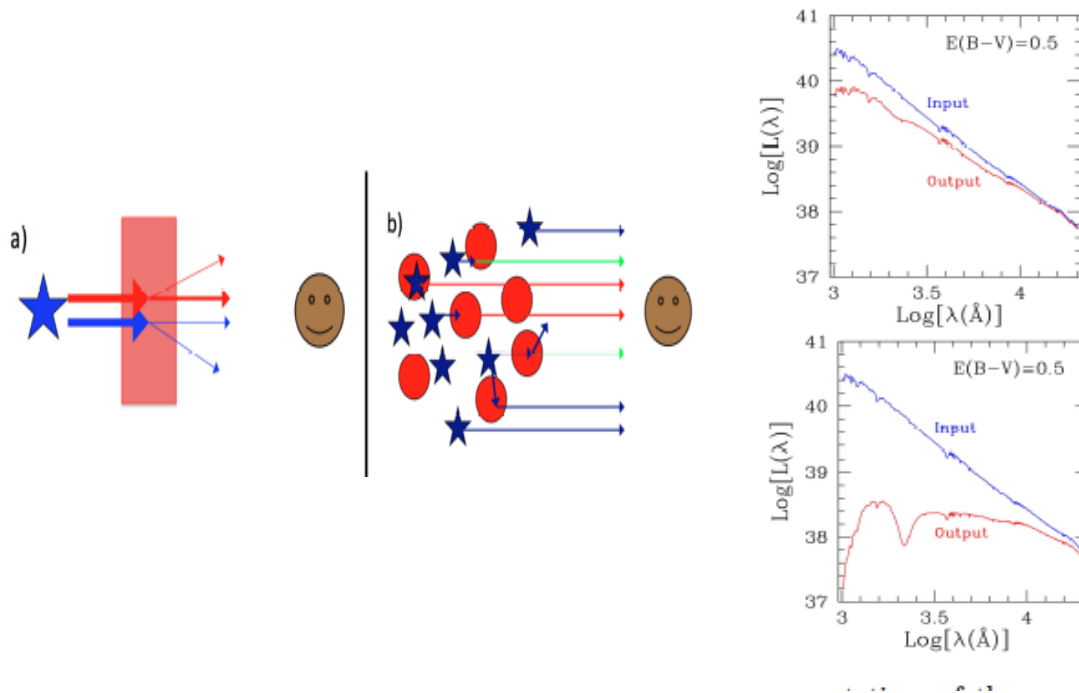
## IR Continuum

- Wavelength at which emission peaks is related to temperature of dust  
 $8\mu \sim 360\text{k}$ ,  $24\mu \sim 121\text{k}$ ,  $70\mu \sim 40\text{k}$ ,  $160\mu \sim 20\text{k}$  based on Black Body Formula  
 $\lambda_{\text{peak}} \sim 29\mu / T_{100}$   $\lambda_{\text{peak}}$  in units of microns and  $T$  in units of 100k  
 (these are the common wavelengths for IR space borne instruments IRAS, Spitzer, WISE, Herschel)  
 $T \sim \lambda^{-1}$  but  $L \sim AT^4$  so to get a lot of luminosity at long wavelengths needs a large emitting area,  $A$   
 Temperature is set primarily by equilibrium; energy absorbed=energy emitted and physics of dust grains.  
 Most galaxies are dominated by  $T \sim 20\text{-}40\text{K}$  dust, rapid star forming galaxies up to  $T \sim 100\text{k}$ .  
 Need wide range of temperatures to produce observed IR spectra.  
 $\text{Roughly SFR (M}_{\odot}/\text{yr}) = L_{\text{total IR}} \times 4.5 \times 10^{-44} \text{ ergs/sec (integrating IR from } 8\text{-}1000\mu)$   
 Advantages- relatively free from extinction, can do at high  $z$  with Herschel  
 Problems- requires lots of assumptions and scaling. Need to assume SF rate law (do problem 2.4 in S&G)

## IR Continuum

- Ideal for starburst galaxies because:
- Young stars dominate UV-optical radiation,  $\tau > 1$ ,  $L_{\text{IR}} \sim L_{\text{SB}}$   
 and cross-section of the dust grains for stellar light is higher in the UV than in the optical
- Not ideal for SF in disks of normal galaxies because: a fair fraction of the IR luminosity is produced by dust re-radiation of emission from 'old' stars e.g. cirrus in the MW. - that is the calibration between SFR and  $L_{\text{IR}}$  depends on the age of the system

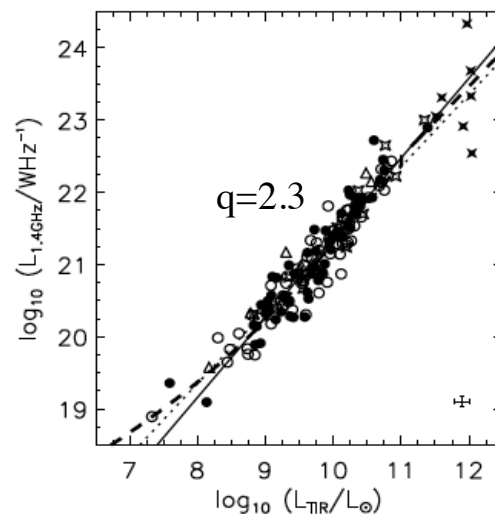
Geometry is a serious issue- the same amount of dust has different effects depending on the relative position of the stars and the dust



## Star Formation- Radio View

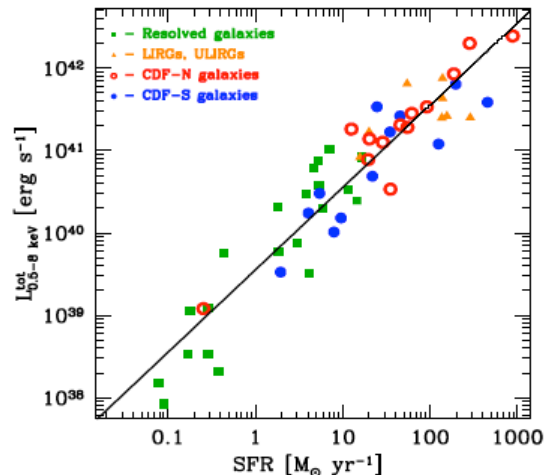
- Radio continuum emission from star-forming galaxies has two components: thermal bremsstrahlung from ionized Hydrogen and non-thermal synchrotron emission from cosmic ray electrons spiraling in the magnetic field of the galaxy
- The relative ratio is frequency dependent because of the different spectral slopes of the 2 processes ( $F_\nu \sim \nu^\alpha$ ,  $\alpha = -0.7$  for synch,  $-0.1$  for TB)
- This method does not depend on how one handles dust or ionizing continuum,
- But physics is not fully understood- why cosmic rays/magnetic field are so finely tuned so that radio synchrotron traces star formation

$$q = \log \left( \frac{FIR}{3.75 \times 10^{12} \text{ W m}^{-2}} \right) - \log \left( \frac{S_{1.4\text{GHz}}}{\text{W m}^{-2} \text{ Hz}^{-1}} \right)$$



# Star Formation X-rays Mineo et al 2012

- In a rapidly star forming galaxies x-rays are produced by
  - 1) high mass x-ray binaries with a lifetime  $\tau \sim 2 \times 10^7$  yrs  
surprisingly the luminosity function of these sources is very similar from galaxy to galaxy with only the normalization  $\sim$  SFR changing
  - 2) hot gas from Supernova- results imply that only 5% of SN energy is needed to produce "diffuse" x-rays



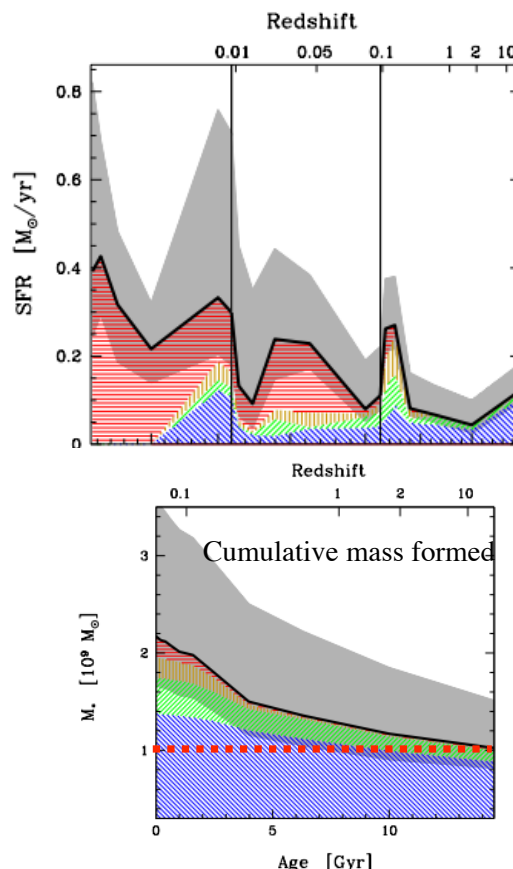
major advantage of x-rays: do not need to be concerned about dust, can do this at high redshift

## How to Infer SFR from Optical Data

- Construct stellar evolutionary tracks containing parameters such as  $T_{\text{eff}}$ ,  $L_{\text{bol}}$ ,
  - These are typically obtained via atmospheric models & spectral libraries
- Construct IMFs containing parameters such as Luminosity, Color, Spectra of Single Age Population
- Add together IMFs from step 2 to get spectra & colors of a galaxy with an arbitrary star formation history
  - Lots of parameters to determine (see <http://arxiv.org/pdf/1208.5229.pdf>) for a detailed discussion of the steps and uncertainties
- 1) Star Formation History
  - 2) Galaxy Age
  - 3) Metal Abundance
  - 4) IMF
- One iterates by comparing the actual galactic emission to the output of a set of galactic stellar population models. The models that best fit the observed data are then used to estimate the galactic properties of interest (e.g. stellar mass, present star formation rate, internal extinction etc.);
- How to handle dust??

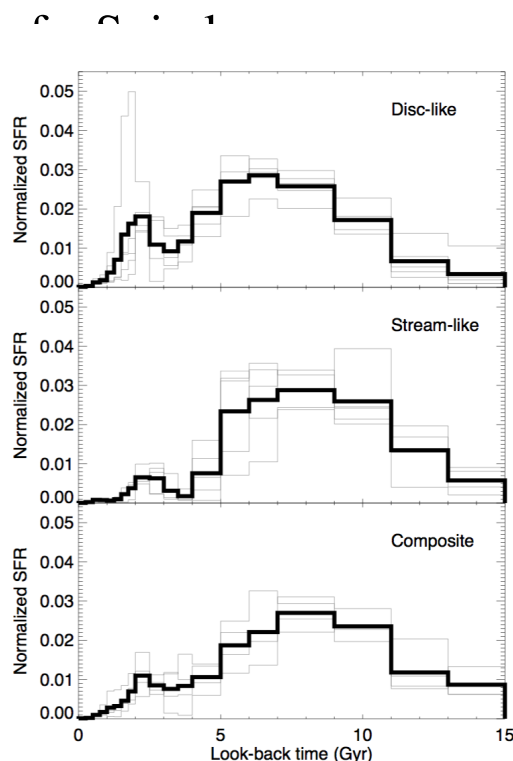
# Uncertainties in Estimating Stellar Masses

- Star formation history- only in a few nearby galaxies can the star formation history be determined
  - e.g. LMC an initial burst of star formation (1/2 mass formed), then a quiescent epoch from ~12 to 5 Gyr ago. Star formation then resumed and continues at an average rate of roughly  $0.2 M_{\odot}/\text{yr}$ , with variations at the factor-of-two level (Harris and Zaritsky 2010)
- IMF uncertainty: fundamental, factor of 2 in transformation of light to mass (also how many binaries!)
- Metallicity: less important (30% effect)
- Different stellar evolution codes- can be very important at different ages (factor of 2)
- Spatial variation in SF history/rate



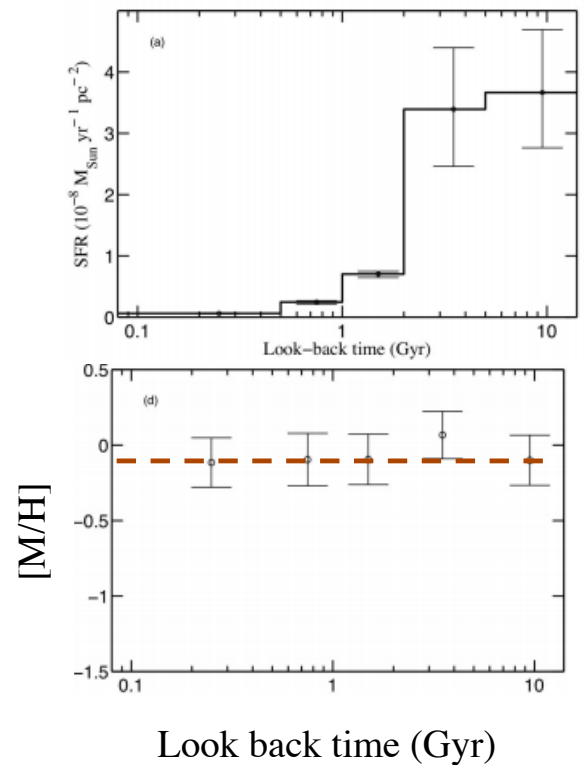
## SF History

- M31 has some of the best data
- In general
- The disc formed most of its mass (~65 percent) since  $z \sim 1$  giving a median age of 7 Gyr,
- with one quarter of the stellar mass formed since 5 Gyr.



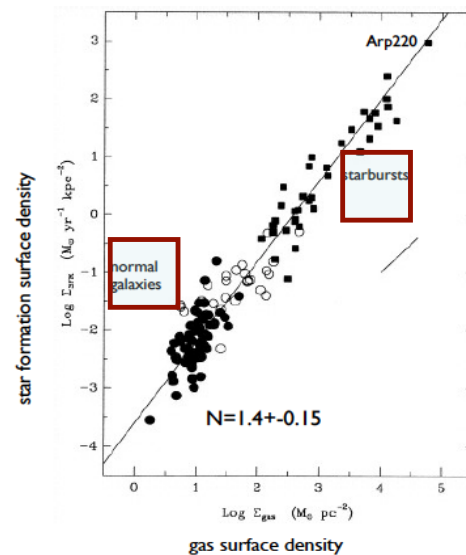
# Star Formation History of an Elliptical

- M32- a dwarf elliptical companion of M31 is close enough to have a CMD for resolved stars-
- very different history than the LMC
- ~95% of its mass formed 5-14 Gyr ago. 2 dominant populations; ~30%  $\pm$  7.5% of its mass 5-8 Gyr old population, ~65%  $\pm$  9% of the mass in a 8-14 Gyr old population (Monachisi et al 2012)
- Metallicity does not change with time (!)- where do the created metals go (another lecture)
- M31 has yet another history the stellar populations of the inner regions of the disk and spheroidal components of M31 are older and more metal-poor than M32



## Kennicutt Schmidt Law (MBW sec 9.5)

- Assume that SFR rate is proportional to total amount of gas
- $SFR \sim \rho_{gas} \sim d\rho_{gas}/dt$ ; sol't  $\rho_{gas} \sim \rho(0)_{ga} e^{-t/\tau}$
- More generally assume  $SFR \sim \rho_{gas}^n$
- e.g. as gas compresses stars form more easily or there maybe another timescale in the process such as the free-fall time of the gas  $SFR \sim \rho_{ga}/t_{freefall}$
- **Frequently this expressed in terms of surface density (an observable)**



Kennicutt 1998

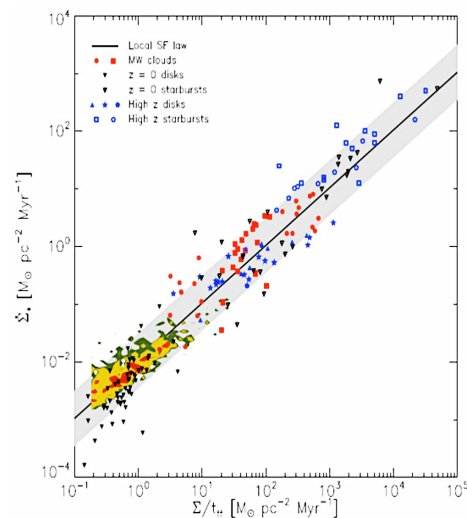
- The motivation for this scaling comes from the gravitational instability of cold gas-rich disks, although the normalization depends on feedback physics.

- $\Sigma_{\text{SFR}} = A \Sigma_{\text{gas}}^n$   $n \sim 1.4$  can be explained by

stars form with a characteristic timescale equal to the free-fall time in the gas disk, which in turn depends inversely on the square root of the gas volume density,  $\tau_{\text{ff}} \sim \rho_{\text{gas}}^{-1/2}$  **for a fixed scale height**  $\rho_{\text{gas}} \sim \Sigma_{\text{gas}}$

gas consumption efficiency is low; takes  $\sim 1.5 \times 10^9$  yrs to convert the gas into stars

## Kennicutt Schmidt Law (MBW sec 9.5)



Not only for whole galaxies but also for parts of them ([Krumholz et al. 2012](#)).

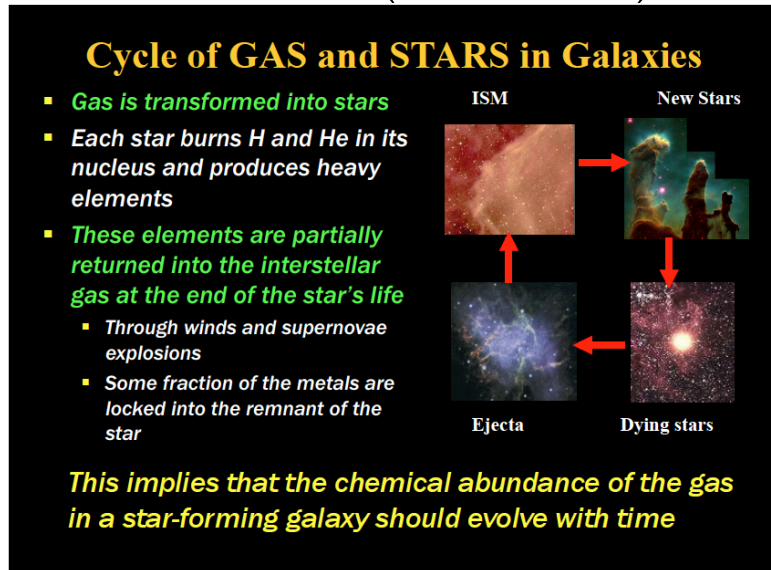
## Basic Equations of Star Formation- see S+G 4.3.2

(1) $M = M_s + M_g$	{	$M$ = total mass in baryons $M_s$ = mass in stars $M_g$ = mass in gas
(2) $\frac{dM}{dt} = f - e$	{	$f$ = rate of infalling gas $e$ = rate of ejected gas
(3) $\frac{dM_s}{dt} = \Psi - E$	{	$\Psi$ = star formation rate $E$ = gas ejection rate of all stars
(4) $\frac{dM_g}{dt} = -\Psi + E + f - e$		

- D. Elbaz; based on Tinsley 1980, Tinsley 1980, Fund. Of Cosmic Physics, 5, 287-388; Maeder 1982

# Basics of Chemical Evolution (MBW 10.4)

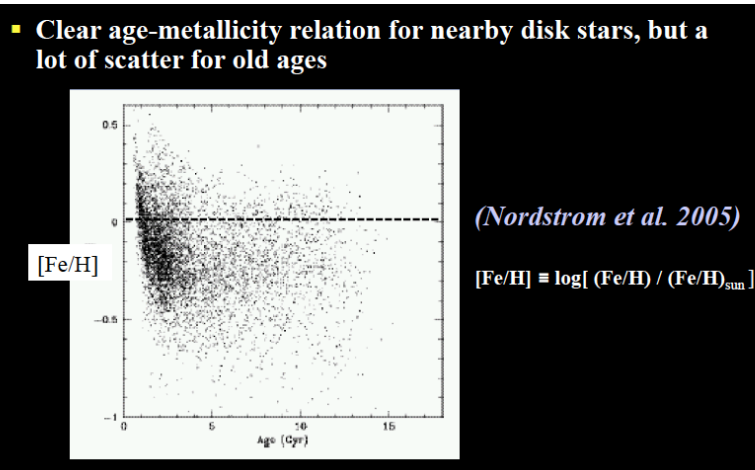
- H and He were present very early on in the Universe, while all metals (except for a very small fraction of Li) were produced through nucleosynthesis in stars
- The fraction by mass of heavy elements is denoted by  $Z$



The Sun's metal abundance  $Z_{\text{sun}} \sim 0.02$

– The most metal-poor stars in the Milky Way have  $Z \sim 10^{-5} \text{ -- } 10^{-4} Z_{\text{sun}}$

## Generic Predictions



- If a galaxy is a **closed box** predict increase of metallicity with time
- Since alpha elements are mostly produced by SnII (from massive short lived stars) while significant amounts of Fe from type Is (longer lived white dwarf binaries) change in chemical composition with age

Pagels Nucleosynthesis and chemical evolution of galaxies 1997 - intro at <http://ned.ipac.caltech.edu/level5/Pagel/frames.html>

Review article

1997ARA&A..35..503

Andrew McWilliam,

Abundance Ratios and Galactic Chemical Evolution



## Closed Box Model (MBW 10.4.2, S+G 4.3.2)

- One-Zone, Closed Box
- – Galaxy's gas is well-mixed
- – No infall, no outflow
  - $M_{\text{tot}} = M_{\text{gas}} + M_{\text{star}} = M_{\text{g}} + M_{\text{s}} = M_{\text{baryons}} = \text{constant}$
  - $M_{\text{H}}$  mass of heavy elements in gas  $= Z_{\text{g}} M_{\text{g}} = Z M_{\text{g}}$
- Instantaneous recycling approximation:
  - The (high-mass) stars return their nucleosynthetic products rapidly (much faster than the time to form a significant fraction of the stars)
  - $dM_{\text{s}}' = \text{total mass made into stars}$
  - $dM_{\text{s}}'' = \text{amount of mass instantaneously returned to ISM .....from SNe, etc; enriched with metals}$
  - $dM_{\text{s}} = dM_{\text{s}}' - dM_{\text{s}}'' = \text{net matter turned into stars}$
  - $y = \text{yield of heavy elements (made instantaneously)}$
  - So  $y dM_{\text{s}} = \text{mass of heavy elements returned to ISM}$

## Closed Box Model- Reminder

- Stellar evolution theory says  
Only stars more massive than  $\sim 8 M_{\text{sun}}$  make heavies (SNe)
- $dM_{\text{s}}'' / dM_{\text{s}} \sim 0.20 = \text{fraction of mass returned to ISM}$
- $y \sim 0.01$  (depends on stellar evolution and Initial Mass Function -IMF)
- $Z(\text{shed gas}) = (\text{heavies shed}) / (\text{mass shed}) = y dM_{\text{s}} / dM_{\text{s}}'' = 0.01/0.2 = 0.05$   
(compared with  $Z_{\text{sun}} \sim 0.02$ )
- Mass conservation implies:  $dM_{\text{g}} + dM_{\text{s}} = 0$  (1)  
Net change in metal content of the gas:
  - $dM_{\text{h}} = y dM_{\text{s}} - Z dM_{\text{s}}$
  - $dM_{\text{h}} = (y - Z) dM_{\text{s}}$  (2)
- Change in  $Z$ 
  - Since  $dM_{\text{g}} = -dM_{\text{s}}$  and  $Z = M_{\text{h}} / M_{\text{g}}$
  - $dZ = dM_{\text{h}} / M_{\text{g}} - M_{\text{h}} dM_{\text{g}} / M_{\text{g}}^2$
  - $= (y - Z) dM_{\text{s}} / M_{\text{g}} + (M_{\text{h}} / M_{\text{g}}) (dM_{\text{s}} / M_{\text{g}}) = y dM_{\text{s}} / M_{\text{g}}$
  - $dZ/dt = -y (dM_{\text{g}}/dt) / M_{\text{g}}$

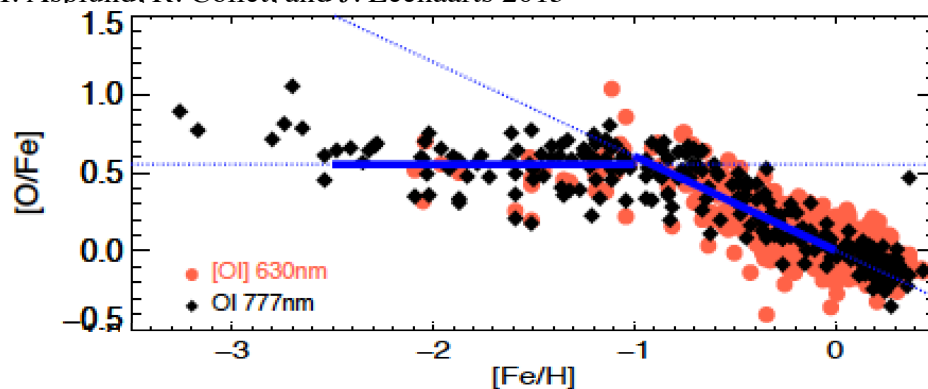
## Closed Box- continued

- Assuming  $y = \text{constant}$  (i.e. independent of time and  $Z$ ):
- $Z(t) = Z(0) - y \ln [M_g(t)/M_g(0)] = Z(0) - y \ln \mu(t)$   
where  $\mu = \text{gas (mass) fraction} = M_g(t) / M_g(0) = M_g(t) / M_t$
- The metallicity of the gas grows with time, as new stars are formed and the gas is consumed
- Metallicity Distribution of the Stars
- The mass of the stars that have a metallicity less than  $Z(t)$  is
- $M_s [< Z(t)] = M_s(t) = M_g(0) - M_g(t)$
- or  $M_s [< Z(t)] = M_g(0) * [1 - \exp(-(Z(t)-Z(0))/y)]$
- When all the gas has been consumed, the mass of stars with metallicity  $Z, Z + dZ$  is  
 $dM_s(Z) \propto \exp(-(Z-Z(0))/y) dZ$

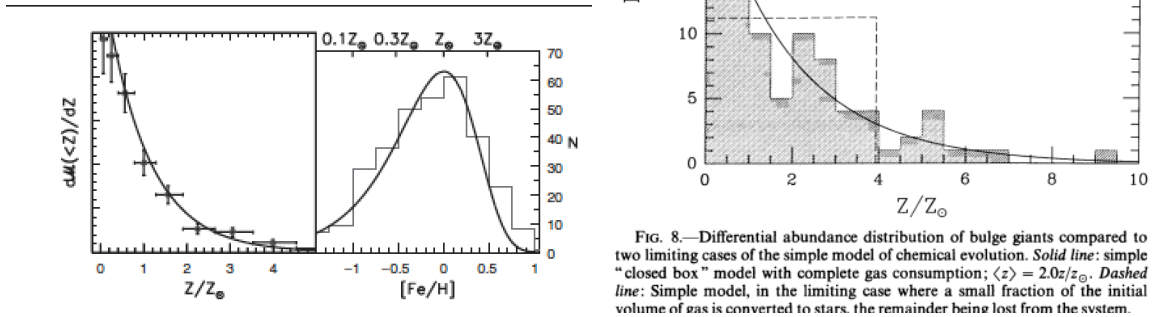
## Yield Derived From Observations

- $Z(\text{today}) \sim Z(0) - y \ln [M_g(\text{today}) / M_g(0)]$
- The average metal content of the gas in the disk near the Sun is  $Z \sim 0.7 Z_{\text{sun}}$
- The initial mass of gas  $M_g(0) = M_s(\text{today}) + M_g(\text{today})$  where  
 $M_s(\text{today}) \sim 40 M_{\text{sun}}/\text{pc}^2$  and  $M_g(\text{today}) \sim 10 M_{\text{sun}}/\text{pc}^2$
- Assuming that  $Z(0) = 0$ , we derive  $y \sim 0.43 Z_{\text{sun}}$

Amarsi, M. Asplund, R. Collet, and J. Leenaarts 2015



## Solutions with Infall and Outflow

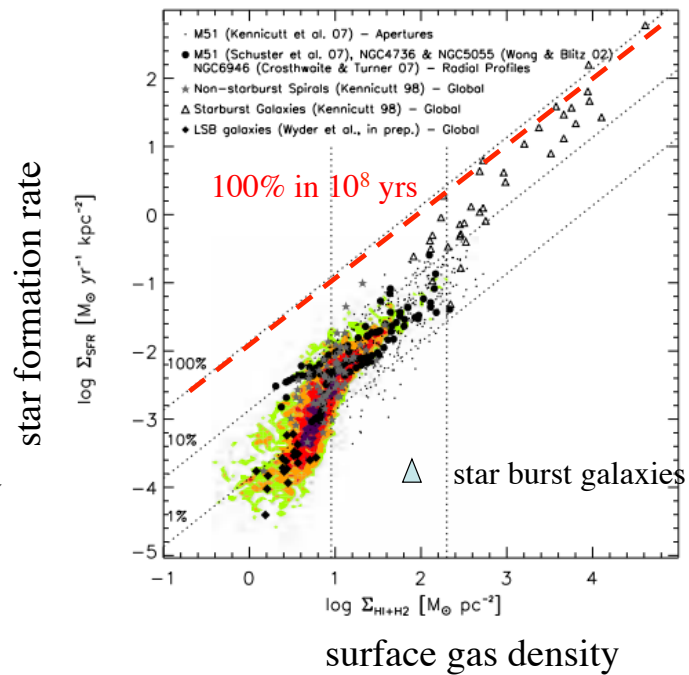


## Solutions with Infall and Outflow

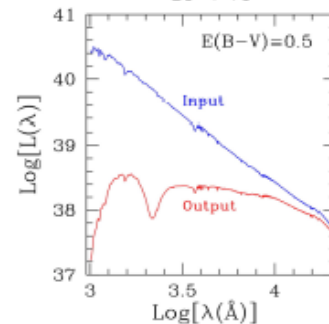
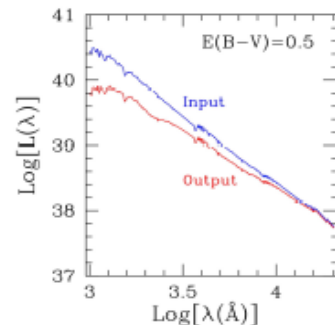
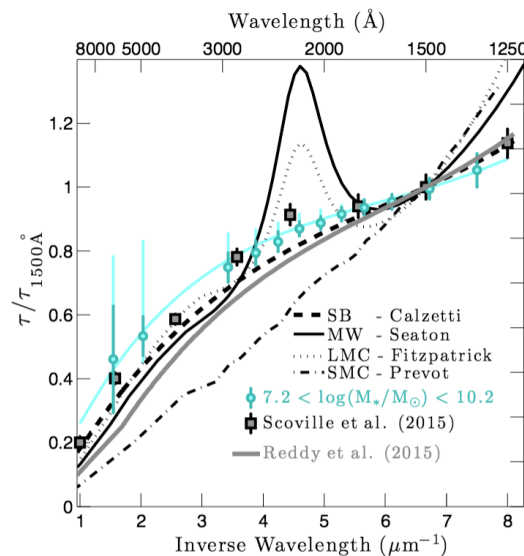
- $dM_{\text{met}} = (y-Z) \cdot \text{SFR}(1-R) - [\text{metals in outflow}] + [\text{metals in infall}]$
- $dM_{\text{gas}} = -\text{SFR}(1-R) - \text{outflow} + \text{infall}$
- $dM^* = \text{SFR}(1-R)$
- $Z = M_{\text{met}}/M_{\text{gas}}$
- Here  $y$ =yield,  $R$ =recycled fraction (fraction of mass that was in stars that ends up back in the ISM)
- credit Eric Bell

# Kennicutt Law with Starbursts

- Newer data show starburst galaxies- the highest star forming rate galaxies obey a different law
- They seem to convert the gas into stars 'as fast as possible'-e.g on freefall timescale.
- This produces a wind as a large amount of energy is injected by star formation in a short time.

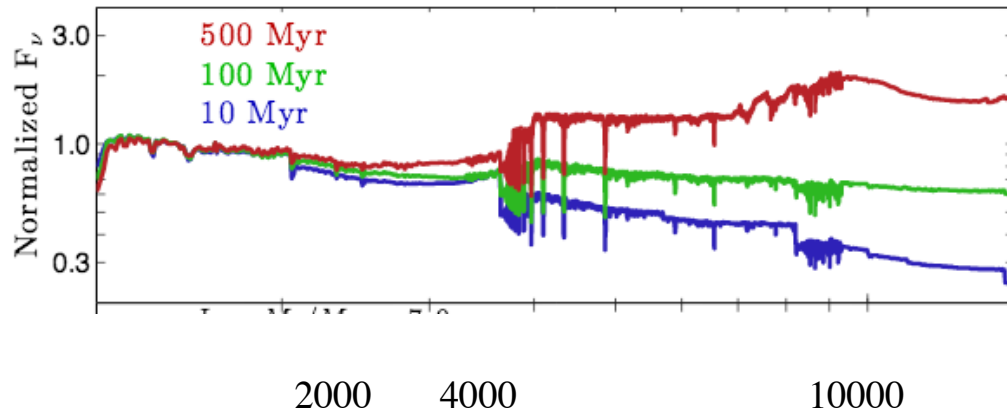


Geometry is a serious issue- the same amount of dust has different effects depending on the relative position of the stars and the dust – today on Astro-ph 1511.00651 Zeiman et al 'The Dust Attenuation Curve Versus Stellar Mass For Emission Line Galaxies At  $Z \sim 2$ '



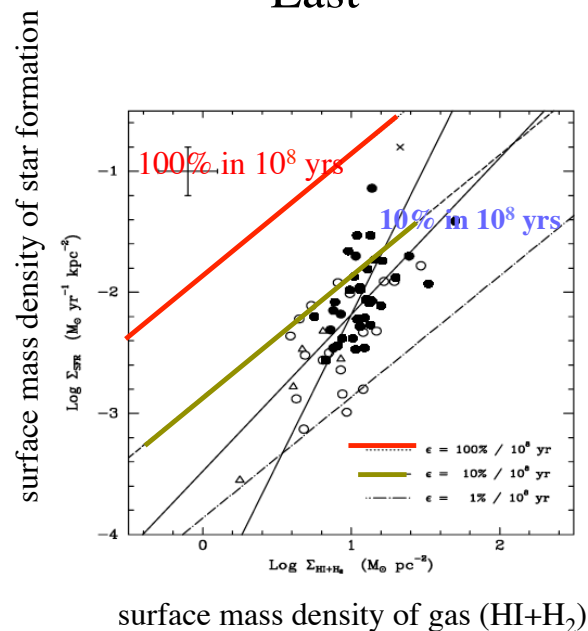
## Broad Band SED of a SSP

- Forming stars at a constant rate for 10,100,500 Myrs (blue, green, red)



- SF in normal galaxies use about 5% of available gas every  $10^8$  yrs !
  - But this does not include 'recycling'- e.g. when stars die they recycle gas back into the ISM
- Since the typical gas mass fraction in disks  $\sim 20\%$  of baryonic mass (but changes a lot as a function of mass), implies that stellar mass of the disk grows by about 1% per  $10^8$  years, i.e. the time scale for building the disk (at the present rate) is  $\sim$  Hubble time.
- The average gas depletion timescale,  $\sim 2.1$  Gyr.
- Recycling of interstellar gas from stars extends the actual time scale for gas depletion by factors of 2–3

## How Long Does the Gas Last



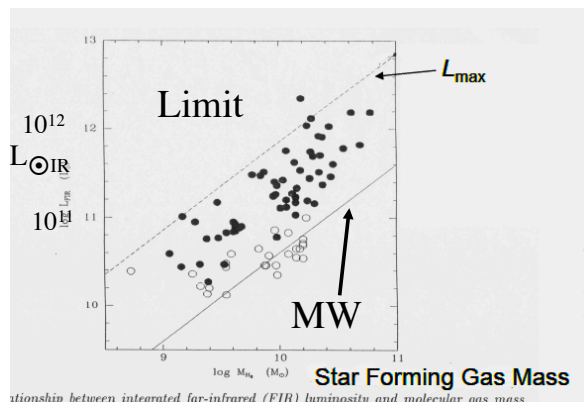
Relationship for 'normal' star formation  
Kennicutt 1998

- Starburst use up their gas much faster
- $\langle 30\% \rangle$  of gas used every  $10^8$  yr
- Depletion timescale  $\approx 0.3$  Gyr
- How luminous are these objects?

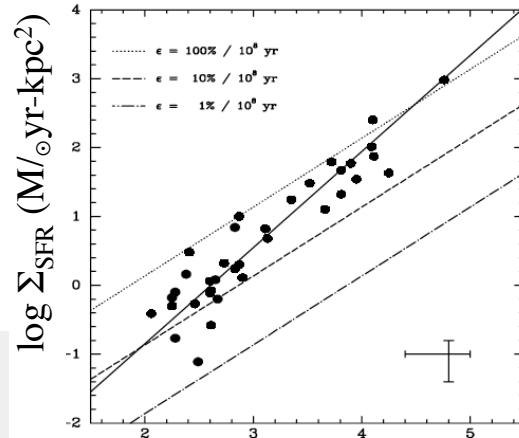
$$\text{SFR}_{\text{max}} \sim 100 \text{ M/yr} (M_{\text{gas}}/10^{10} M_{\odot}) (10^8 \text{ yrs} / \Delta t_{\text{dyn}})$$

nuclear fusion is  $\sim 0.7\%$  efficient

the fraction of rest mass converted to energy for a Saltpeter IMF is  $\epsilon \sim 0.05$  during  $10^8$  yrs



## How Long Does the Gas Last- Star Bursts



$$\log \Sigma_{\text{H1+H2}} (M_{\odot}/\text{pc}^2)$$

$$\text{This gives } L_{\text{max}} \sim 0.07 \epsilon (dM/dt) c^2$$

$$L_{\text{max}} \sim 10^{11} L_{\odot} (M_{\text{gas}}/10^{10} M_{\odot}) (\epsilon/0.05)$$

## Possible Star Formation 'Laws' (MBW sec 9.5.1)

- Define star formation efficiency  $\text{SFE} = \Sigma_{\text{SFR}} / \Sigma_{\text{gas}}$
- to form stars in spirals need
  - cold phase ( $n \sim 4\text{--}80 \text{ cm}^{-3}$ ,  $T \sim 50\text{--}200 \text{ K}$ )
  - and gravitationally bound clouds
- A star formation law *should* predict the SFE from local conditions (physics)
- 1) Kennicutt-Schmidt law for star formation,  $\Sigma_{\text{SFR}} \sim \Sigma_{\text{gas}}^{1.5}$  is often interpreted as indicating that the star formation rate is controlled by the self-gravity of the gas. In that case, the rate of star formation
  - will be proportional to the gas mass divided by the time scale for gravitational collapse
- stars form on a characteristic timescale equal to the free-fall time in the gas disk,  $\sim \rho^{-1/2}$
- since  $\rho_{\text{gas}} \sim \Sigma_{\text{gas}}$  and  $\Sigma_{\text{SFR}} \sim \Sigma_{\text{gas}}^{1.5}$
- expect  $\text{SFE} \sim \Sigma_{\text{gas}}^{0.5}$

## Possible Star Formation 'Laws' (MBW sec 9.5.3)

Disk free-fall time : if scale

height of disk set by hydrostatic

equilibrium then

$t_{\text{ff}} \sim \rho^{-1/2}$  related to the velocity field and density of stars and gas

or some other timescale such as orbital timescale  $t_{\text{orb}} = \Omega/2\pi = 2v(r)/2\pi r$

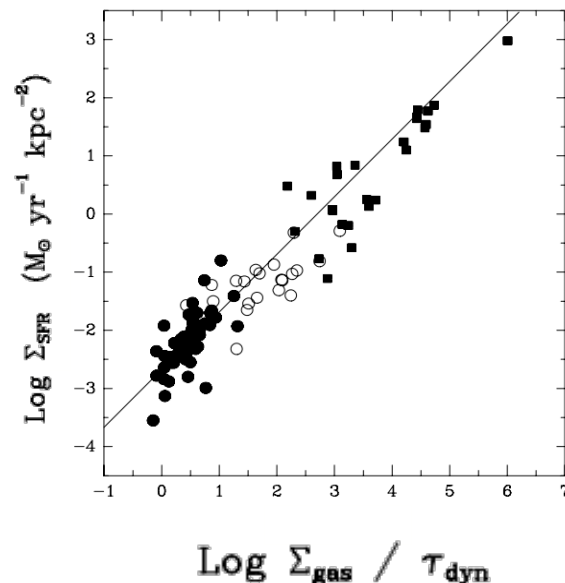
or perhaps gravitational instability - gas unstable against collapse

when Toomre  $Q = \sigma_g \kappa / \pi G \Sigma_{\text{gas}} < 1$  ;

$\kappa$  is the epicyclic frequency; velocity dispersion of the gas  $\sigma_g$

- What other scaling relations seem to hold?
- Kennicutt shows that  $\Sigma_{\text{SFR}}$  is also correlated with  $\Sigma_{\text{gas}}/\tau_{\text{dynamical}}$  in a galaxy sample (but not inside a galaxy!)
- where  $\tau_{\text{dynamical}}$  is the orbital time at the radius of the star forming region  $2\pi R/V_{\text{rot}}(R)$
- In a disk galaxy there is another restoring force (other than pressure) which is important for the Jeans criterion- conservation of angular momentum (Coriolis force)
- Perturbations are unstable to gravitational collapse if
- $Q = c_s \kappa / p G \Sigma_{\text{mass}} < 1$  where  $\kappa$  is the epicyclic freq (MBW 9.10)
- The K-S  $\tau_{\text{dynamical}}$  law follows if  $Q < 1$  (Silk)

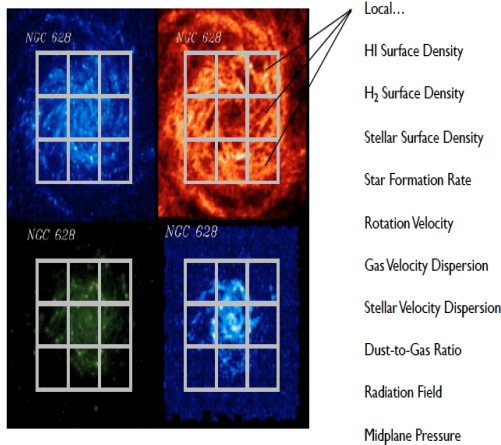
### Kennicutt Schmidt Continued



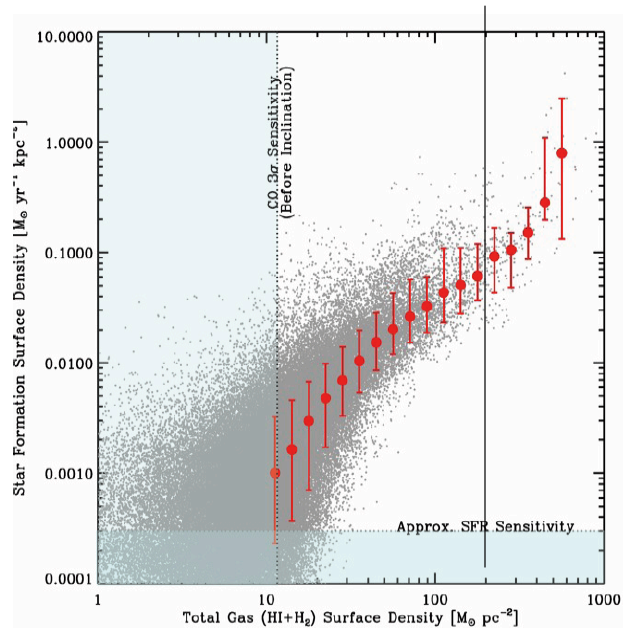
# Kennicutt-Schmidt Updated

GMC density

- SFR depends on surface density of **molecular gas**
- (red points are averages, gray points are individual samples in galaxies)



F. Walter



## Jeans Criterion for collapse of spherical cloud

MBW pg 167, sec 8.2.3, 9.1.2

- Gravitational instability sets in if the free-fall time is **less than** the sound crossing time
- $t_{\text{ff}}^2 = 1/G\rho < (R/c_s)^2 = 10^8 n_H^{-1/2} \text{ yrs}$ ;
- free fall time from  $d^2r/dt^2 = -GM/r^2$ ;  $n_H$  is the number density of gas
- hydrodynamical timescale from  $d^2r/dt^2 = (-1/\rho(r))dP/dr = R/c_s$

**Characteristic mass for system to collapse is Jeans Mass**



## Jeans Criterion for collapse of spherical cloud

Jeans mass  $M_J = 4/3\pi\lambda_J^3\rho = 4/3\pi c_s^3\rho^{-1/2}$

Jeans length  $\lambda_J = \sqrt{\pi c_s^2/G\rho}$  – distance a sound wave travels in a grav free-fall time

For typical values

- $M_J = [\pi^{5/2}/6][c_s^3/(G^3\rho)] \sim 40M_\odot [c_s/0.2\text{kms}^{-1}]^3 [n_{\text{H}_2}/100\text{cm}^{-3}]^{-1/2}$

$M_{J \text{ SOLAR UNITS}} = (T/10\text{k})^{3/2} (n_{\text{H}}/10^5\text{cm}^{-3})^{-1/2}$

units of surface mass density  $\lambda_J = c_s^2/G\Sigma$

$c_s = \text{sound speed} = \sqrt{dP/d\rho} = \sqrt{k_B T/\mu m_{\text{H}}}$  for hydrogen ( $k_B =$  Boltzmann's constant,  $m_{\text{H}} =$  mass of hydrogen atom,  $\mu =$  mean molecular weight)

- For typical values  $c_s = 0.3\text{km/sec}(T/10\text{k})^{1/2}$

However the gas cannot collapse unless it can radiate away the heat from conversion of potential energy so need  $t_{\text{cool}} < t_{\text{ff}}$  the rate at which gas cools depends on a strong function of temperature and the density squared.

if one has an external pressure MBW eq 9.3

- For an isothermal sphere in pressure equilibrium with its surrounding,
- $M_{\text{BE}} = 1.182 c_s^3 / (G^3 \rho)^{1/2} = 1.182 c_s^4 / (G^3 P_{\text{th}})^{1/2}$
- where  $P_{\text{th}} = \rho c_s^2$  is the surface pressure

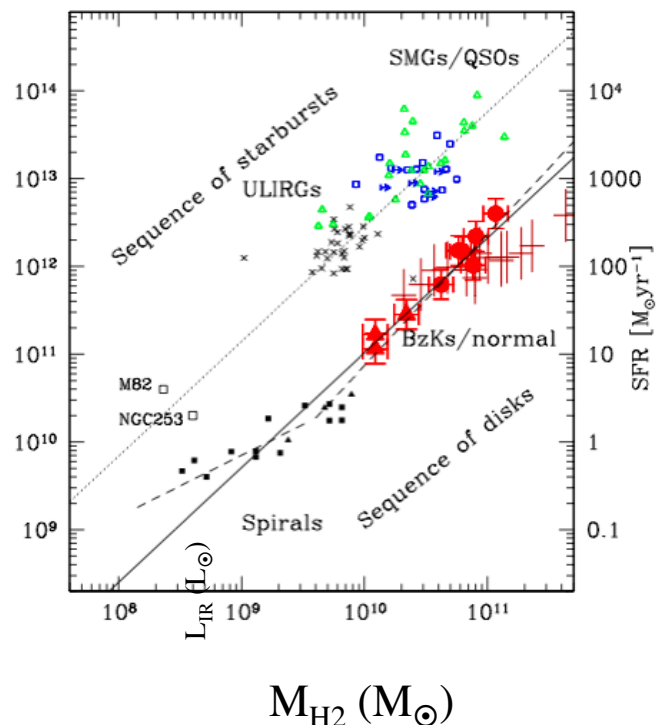
# Gravitational Instability

- Another derivation of Jeans length/mass
- Balance pressure and gravity (pg 355 of S+G)
- Potential energy =  $-1/2 \int \rho(x)\phi(x)d^3x \sim G\rho^2 r^5$
- if gas moves as sound speed  $KE = c_s^2 M$
- $M = 4/3\pi\rho r^3$
- In equilibrium virial theorem says  $KE = PE/2$  so define a length  $\lambda_j$  where that is true and get  $\lambda_j = c_s \sqrt{\pi/G\rho}$

The cloud's radius is the Jeans' Length and its mass ( $4/3\pi\rho\lambda_j^3$ ) is the Jeans mass -when thermal energy per particle equals gravitational work per particle. At this critical length the cloud neither expands nor contracts. Dimensionally this is  $kT = GM/r$

## Starbursts- Higher Redshifts

- In the low redshift universe there are very few, very high SFR objects- these are much more important in the high  $z$  universe
- It appears that the relations for very rapid SF galaxies are different



## Star formation Occurs in Giant Molecular Clouds

- Cooling to  $10^4$  K is not sufficient to form stars.
- The gas has to cool well below 100K and must be shielded from UV radiation by dust.
- Star formation occurs in giant molecular clouds with masses of  $10^3$ - $10^7$  M and radii of 1-100pc.
- These clouds can become gravitationally unstable and collapse and form stars .
- The effects of feedback (e.g. stellar winds and SNR) are not at all clear

## Molecular Clouds MWB sec 9.2

- As the gas density increases the fraction that is molecular increases rapidly (a sharp transition)-  $H_2$  forms on dust grains when it is cold
- These clouds are in rough virial equilibrium  $2GM/\sigma^2=R$ ,  $M\sim R^2$ ,  $\delta V\sim R^{1/2}\sim\rho^{-1/2}$
- $M\sim 10^5$ - $5\times 10^6 M_\odot$ ,  $r\sim 10$ 's pc  $n_{H_2}\sim 100$ - $500\text{cm}^{-3}$  but there is a lot of structure, in protostellar cores density much higher
- Cold  $T\sim 10$ k in MW) - UV light cannot penetrate- heating by Cosmic rays (?) - quite turbulent
- Strongly associated with young star clusters- short lived (?)  $t\sim 10^7$  yrs

## Molecular Clouds MWB sec 9.2

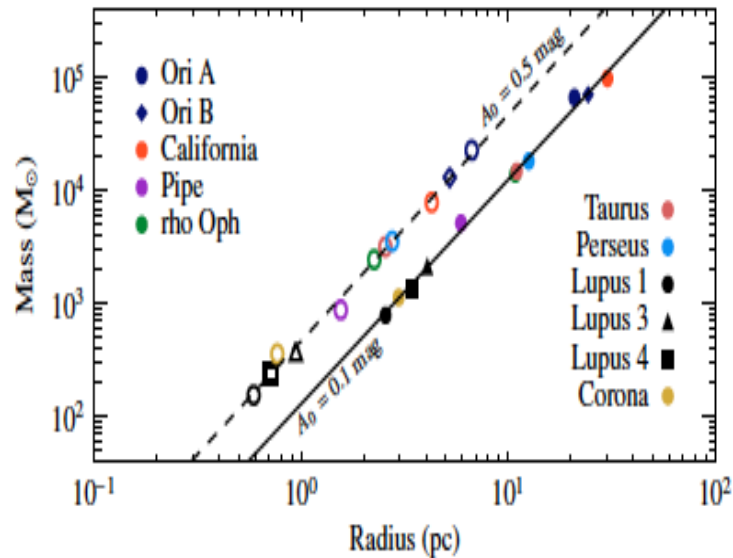
If self gravitating, isothermal spheres collapse if mass exceeds the Jeans mass  $M_J \sim 40 M_\odot (c_s/0.2 \text{ km/sec})^3 (n_{\text{H}_2}/100)^{-1/2}$

which they do by a lot~!

collapse on free fall time

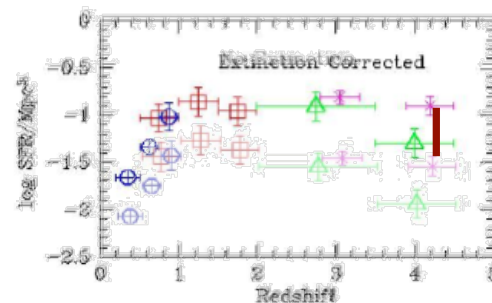
$$t_{\text{ff}} = (3\pi/32G\rho)^{1/2}$$

$$\sim 3.6 \times 10^6 (n_{\text{H}_2}/100)^{-1/2} \text{ yrs}$$



## Dust

- As we discussed before the effects of dust and how one treats it can be a very large effect.
- As an example take the star formation history of the universe as revealed by deep 'optical' studies- it shows that 'correcting for dust' introduces a factor of 3 change!
- Correcting for dust is not easy to do



For those interested in more details on starbursts see Peter Barthels course notes

<http://www.astro.rug.nl/~pdb/starbursts.htm>

## Low Z SFR-Summary

Property	Kennicutt 1998	Spiral Disks	Star Bursts
Radius		1 – 30 kpc	0.2 – 2 kpc
SFR		0 – 20 $M_{\odot} \text{ yr}^{-1}$	0 – 1000 $M_{\odot} \text{ yr}^{-1}$
Bolometric Luminosity		$10^6 - 10^{11} L_{\odot}$	<u><math>10^6 - 10^{13} L_{\odot}</math></u>
Gas Mass		$10^8 - 10^{11} M_{\odot}$	$10^6 - 10^{11} M_{\odot}$
Star Formation Timescale		1 – 50 Gyr	0.1 – 1 Gyr
Gas Density		1 – 100 $M_{\odot} \text{ pc}^{-2}$	$10^2 - 10^5 M_{\odot} \text{ pc}^{-2}$
Optical Depth ( $0.5 \mu\text{m}$ )		0 – 2	<u>1 – 1000</u>
SFR Density		0 – 0.1 $M_{\odot} \text{ yr}^{-1} \text{ kpc}^{-2}$	1 – 1000 $M_{\odot} \text{ yr}^{-1} \text{ kpc}^{-2}$
Dominant Mode		steady state	steady state + burst

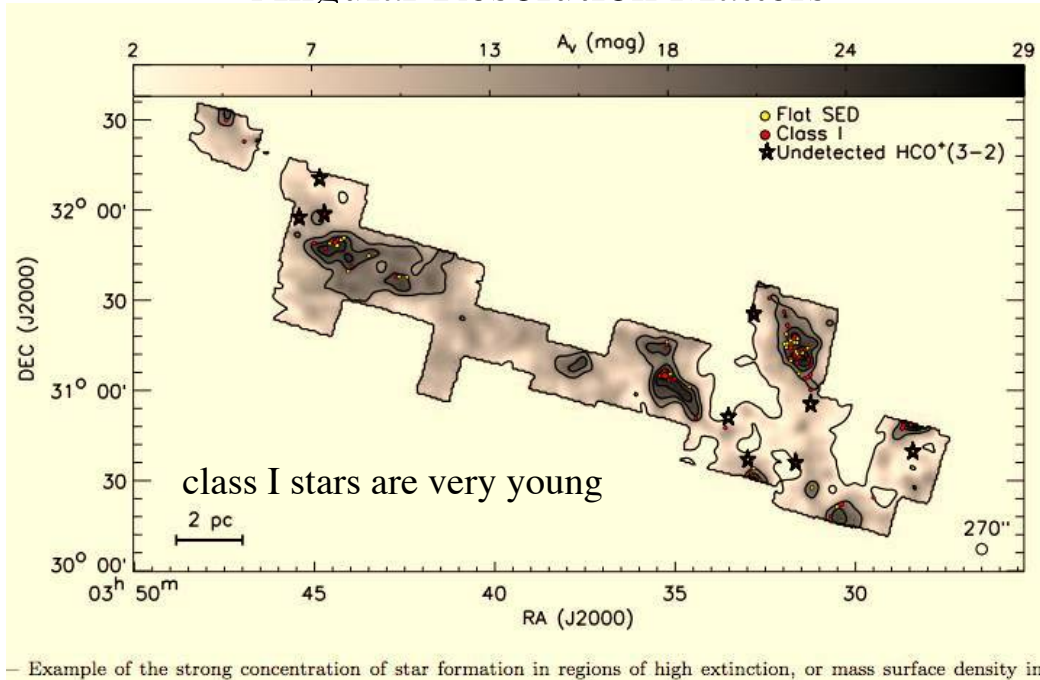
- $t_{\text{freefall}} = (R/G\Sigma)^{1/2}$
- $t_{\text{cross}} = (R/\sigma)$
- the fastest things can happen is when this are equal and make R the Jeans length  
 $R_{\text{Jeans}} \sim \sigma^2/G\Sigma$

## Summary of Situation

- Large scale SFR is determined by a hierarchy of physical processes spanning a vast range of physical scales:
  - the accretion of gas onto disks from satellite objects and the intergalactic medium(Mpc)
  - the cooling of this gas to form a cool neutral phase (kpc)
  - the formation of molecular clouds (10-100 pc);
  - the fragmentation and accretion of this molecular gas to form progressively denser structures such as clumps ( $\sim 1 \text{ pc}$ ) and cores ( $\sim 0.1 \text{ pc}$ )
- The first and last of these processes operate on galactic (or extragalactic) and local cloud scales, respectively, but the others occur at the boundaries between these scales and the coupling between processes is not yet well understood.
- the challenge of explaining the low efficiency of star formation remains.
- Similarly, an understanding of the full IMF, remains elusive.

**Kennicutt and Evans 2012**

# Angular Resolution Matters



In Perseus molecular cloud all the young stars lie in very dusty regions

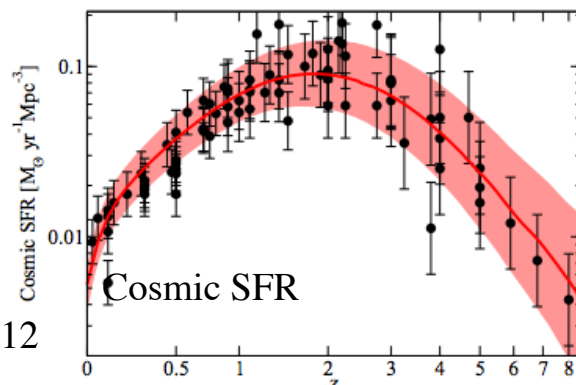
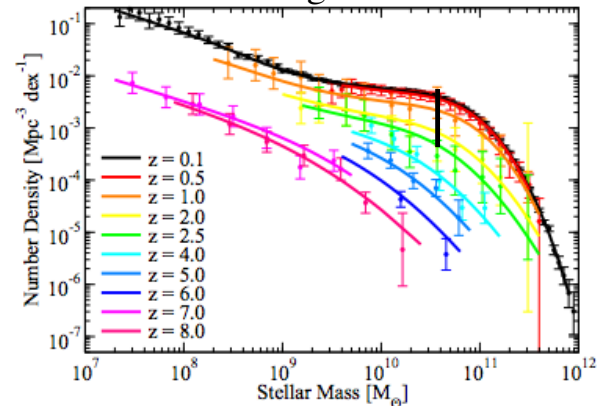
## Cosmic History of Star Formation

### General Results

- 90% of all stars formed since  $z \sim 3$
- SFR has dropped by  $\sim 10\times$  since  $z \sim 1$ .
- The most massive galaxies grow 50:50 by merging  $z < 1$
- Form stars vigorously at  $z > 2$
- Also at earlier epochs: most stars lived in the most massive galaxies at the time
- Effects of dust are dominant at  $z > 3$  and results are uncertain...

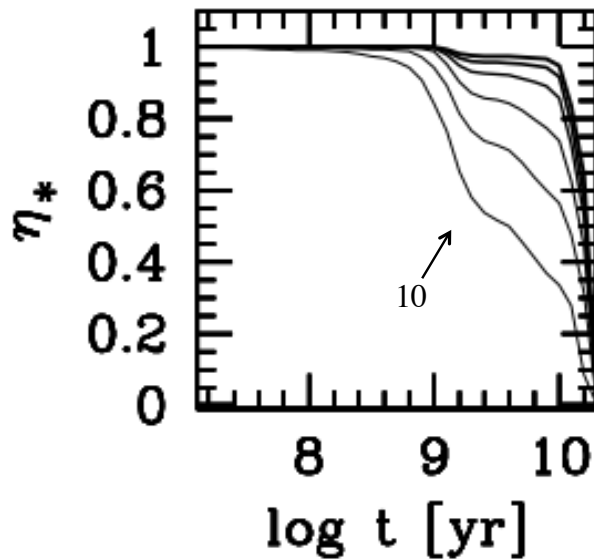
Behroozi et al 2012

## Growth of galaxies



## Results from Stellar Paleontology

- History of **stellar growth** in **6 mass bins 10, 10.3, 10.6, 10.9, 11.2, 11.5 vs time** - *big objects form first, evolve rapidly and then remain the same for long times* (Vale Asari et al 2009) and metallicity  $Z_*$

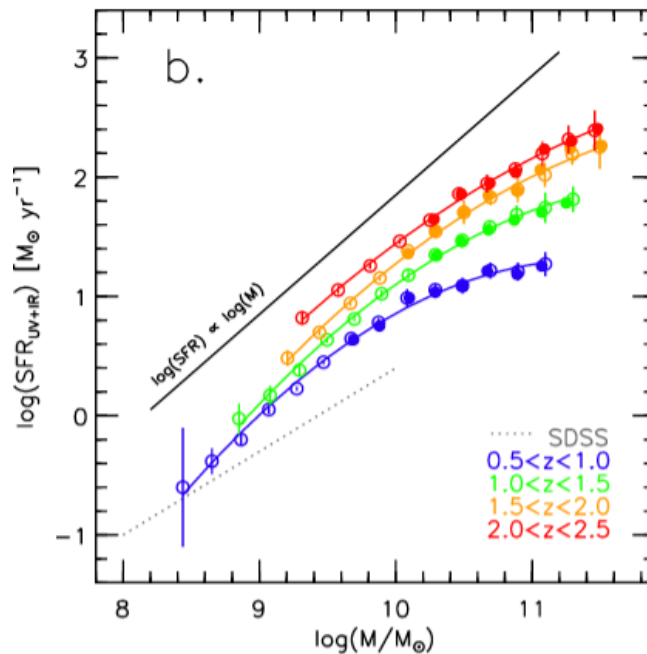


## 'Main Sequence' of Star Formation

- Galaxy surveys out to  $z \sim 4$

show that the majority of star-forming galaxies follow a relatively tight relation between star formation rate ( $\text{SFR} \equiv \Psi$ ) and stellar mass ( $M_*$ ) (Brinchmann et al. 2004; Noeske et al. 2007; Elbaz et al. 2007)

- this is called the main sequence of star formation



Whitaker et al 2015

# Main Sequence of Star Formation

- The MS of Star formation has 3 parameters
- the normalization, intrinsic scatter, and slope
- These are related to the fundamental physical quantities that regulate star formation.
  - The changing normalization of is due predominantly to the changing cosmological gas accretion rates with redshift.
  - The intrinsic scatter of this relation reveals the level of stochasticity in the gas accretion history. Lastly,
  - the measured slope is related to the star formation efficiency.

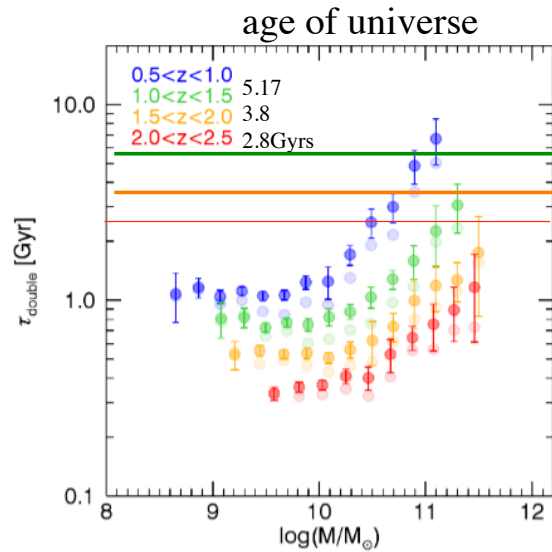
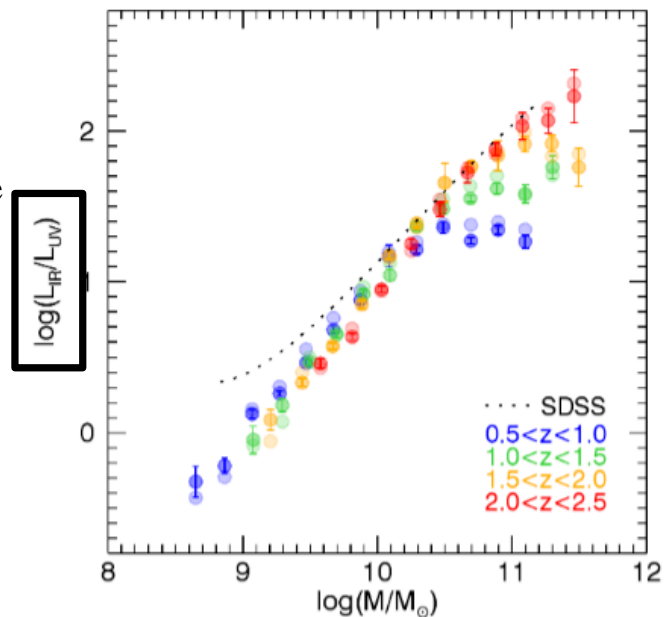


FIG. 12.— The mass-doubling timescale ( $\text{sSFR}^{-1}$ ) as a function of stellar mass for star-forming galaxies. The filled circles

## The Effects of Dust are Mass Dependent

- The luminosity in the IR ( $L_{\text{IR}}$ ) divided by the luminosity in the UV ( $L_{\text{UV}}$ ) for star forming galaxies is a measure of the effects of dust.
- The more massive the galaxy (and thus the higher the SFR) the more important dust is!

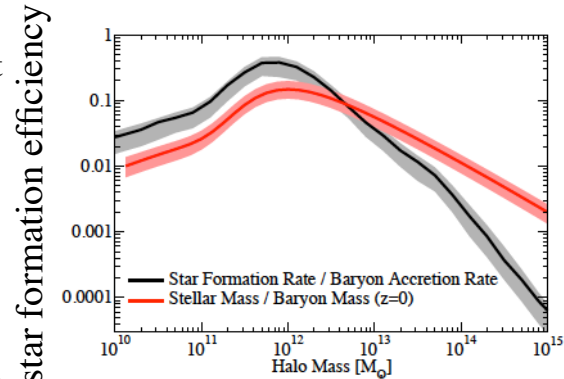


(Whitaker et al 2014)



# Theoretical ideas About Galaxy Wide Star Formation

- Theoretical predictions are that galaxy formation is most efficient near a mass of  $10^{12} M_{\odot}$  based on analyses of supernova feedback and gas cooling times (Silk 1977; Rees & Ostriker 1977; Dekel & Silk 1986; White & Rees 1978; Blumenthal et al. 1984).
- Hydrodynamical simulations indicate that the host dark matter halo mass strongly influences gas accretion onto galaxies
- For low halo masses, simulations predict that gas accretes in cold filaments (“cold mode accretion”) directly to the galaxy disk, efficiently forming stars.
- Above a transition halo mass of  $\sim 10^{11} M_{\odot}$  a shock develops at the virial radius which heats accreting gas (“hot mode accretion”) and rapidly quenches instantaneous star formation

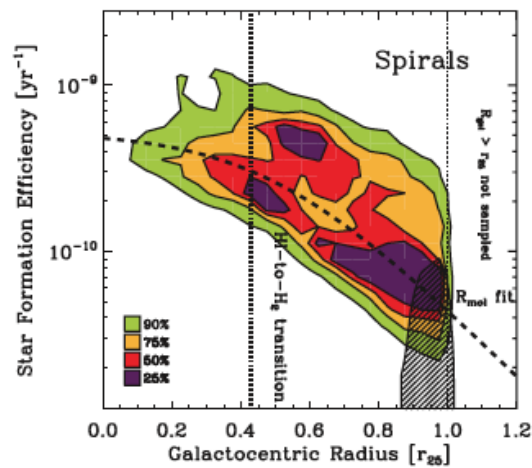


## Star Formation Efficiency Vs Galactocentric Radius

- Where  $H_2$  dominates HI SFE is roughly constant and declines exponentially at larger radii.
- where the ISM is mostly  $H_2$  in spiral galaxies, the SFE does not vary strongly with any of the obvious quantities : including radius,  $\Sigma_{ga}$ ,  $\Sigma_*$ ,  $P_{th}$ ,  $\Omega_{orb}$ ,

### Summary of results

- Molecular gas, star formation, and stellar surface density all decline with nearly equal exponential scale lengths,  $\sim 0.2 r_{25}$ , giving the appearance of a long-lived star-forming disk embedded in a sea of HI.
- The ISM is mostly  $H_2$  within  $\sim 0.5 r_{25}$  and where  $\Sigma_* > 80 M$



## What Does Not Work-Leroy et al 2008

- the disk freefall time for a fixed scale height disk
  - orbital timescale
  - Rotation curve dependences (e.g. orbital timescales)
  - $Q_{\text{gas}}$
- BUT
  - SFE ( $\text{H}_2$ ) is constant as a function of a range of environmental parameters for disks of spiral galaxies, and **not** starbursts or low metallicity dwarf galaxies.
- Despite enormous amount of work at a resolution of 800pc NO unique driver for the SFE, but perhaps
- ISM physics — balance between warm and cold  $\text{H I}$  phases,  $\text{H}_2$  formation, and perhaps shocks and turbulent fluctuations driven by stellar feedback— govern the ability of the ISM to form GMCs out of marginally stable galaxy disk

## Star Formation Efficiency Vs Radius and $\text{H}_2/\text{H I}$

- Leroy et al (2008) show that star formation efficiency (how long it would take to convert all the gas into stars) changes with galactic radius and the relative fraction of gas in  $\text{H}_2$  and  $\text{H I}$ .
- Where  $\text{H}_2$  is dominant the SFE is constant at about  $5 \times 10^{-10}$  or it takes  $\sim 2 \times 10^9$  yrs to convert ALL the local gas into stars.

