

General Introduction to Galaxy Clusters: Scaling Relations

Brian McNamara
University of Waterloo

With thanks to Helen Russell

Astronomical Units & Terminology

$$r_{2500} = 500 \text{ kpc} \left(\frac{T}{5 \text{ keV}} \right)^{\frac{1}{2}}$$

$$M_{2500} = 2 \times 10^{14} M_{\oplus} \left(\frac{T}{5 \text{ keV}} \right)^{\frac{1}{2}}$$

scaled to critical density of Universe

Kiloparsec Distance $1000 \text{ parsecs} = 3 \times 10^{21} \text{ cm}$

Light year Distance 3.26 parsecs

Milky Way Diameter = 100,000 ly = 30 kpc

10^{53} erg Energy supernova explosion/GRB

keV Temperature $1 \text{ keV} = 1.16 \times 10^7 \text{ K}$

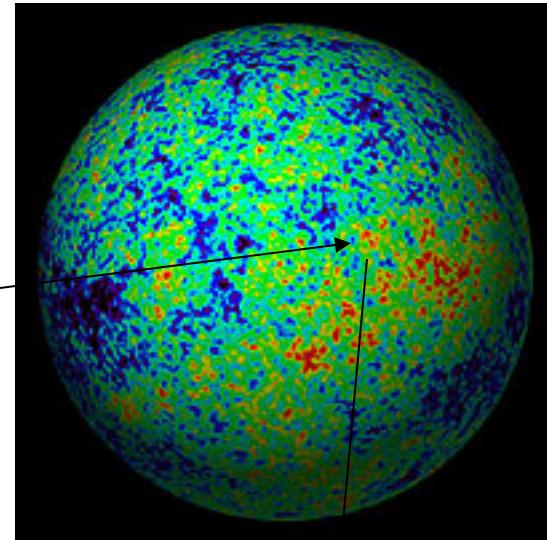
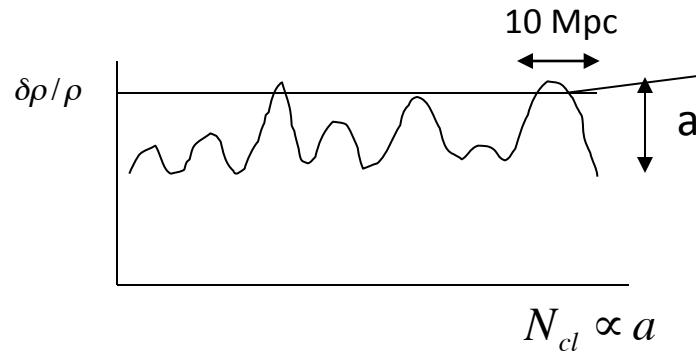
dynes/cm² Pressure cluster center P= $10^{-10} \text{ dyne/cm}^2$

10^{-16} Atm

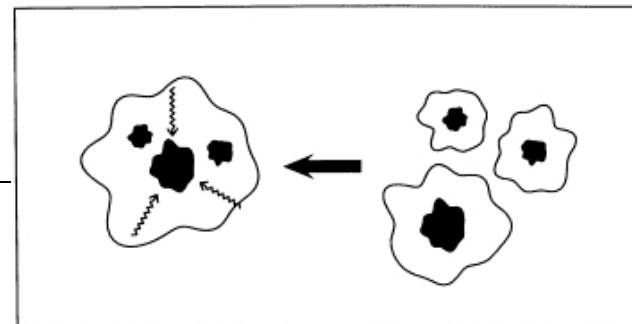
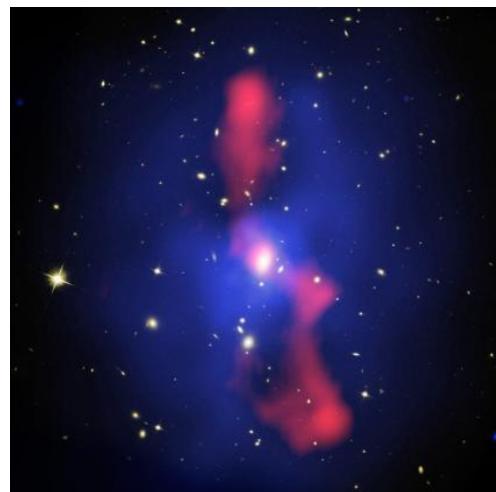
Why Study Galaxy Clusters?

- **Cosmology**
 - evolve by gravity: largest structures evolve slowly--dominated by dark matter, sensitive to matter and energy density of the Universe
 - fair sample hypothesis: baryons, dark matter, baryon fraction
 - observational goal: measuring masses; observational proxies for mass (L_x , T, etc.)
- **Galaxy formation & Evolution**
 - closed boxes: retain byproducts of stellar evolution; history of star formation
 - galaxies at same redshift: uniform sampling in space & time
 - merging & perturbation: dynamical evolution of galaxies
- **Cosmic cooling problem**
 - only few percent of baryons have condensed in clusters
 - cooling flows/feedback

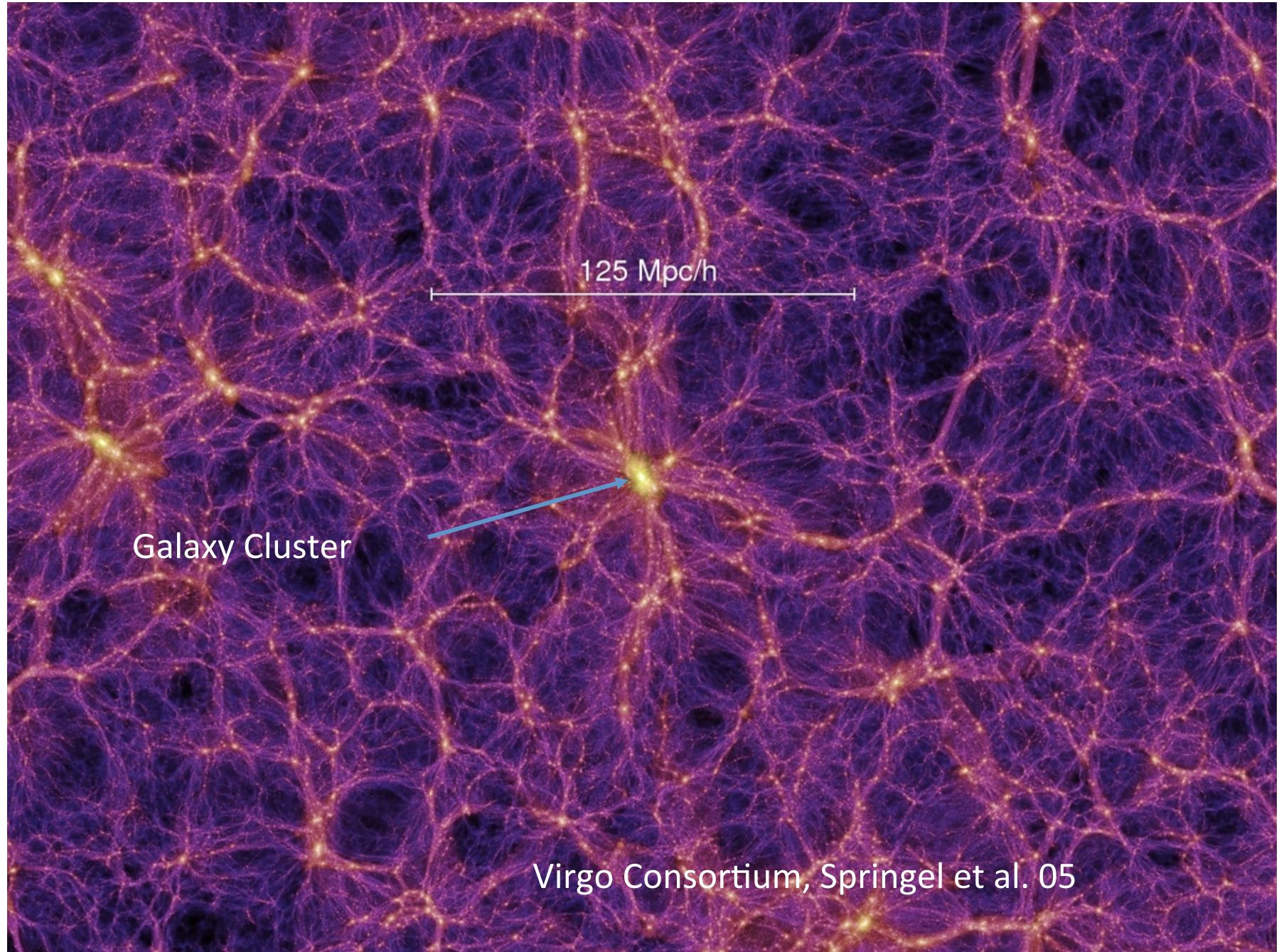
Gaussian density perturbations



Hierarchical clustering

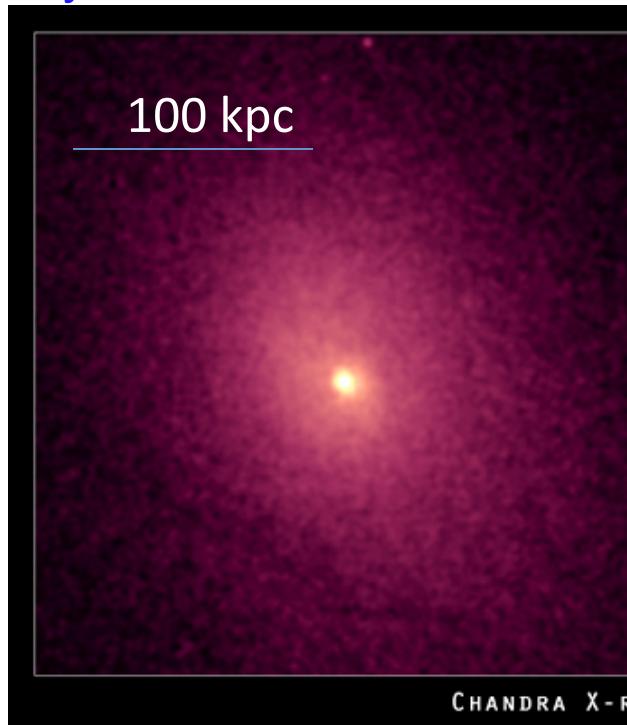


Number density of clusters: amplitude of perturbations, a
Degree of clustering: Ω_m

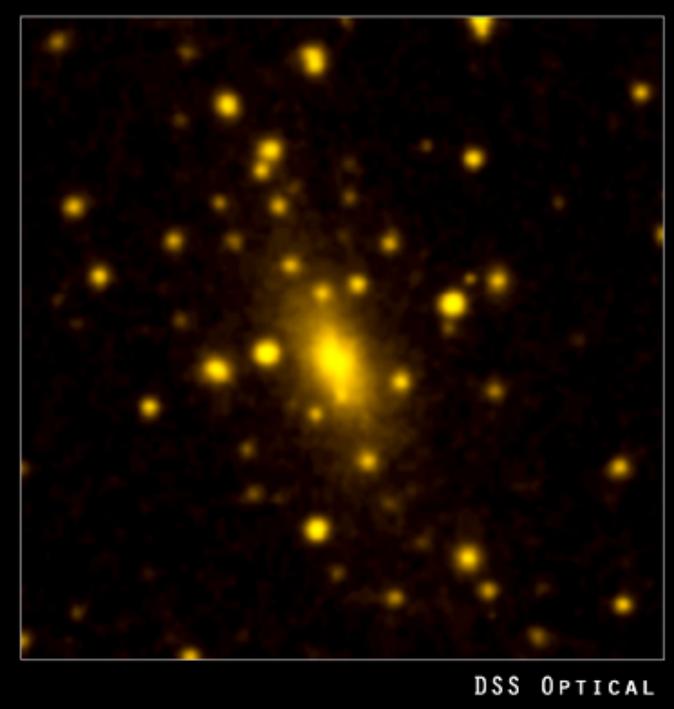


Galaxy Cluster at X-ray & Visual wavelengths

X-ray



visual



$$L_x = 10^{43} - 10^{45} \text{ erg s}^{-1}$$

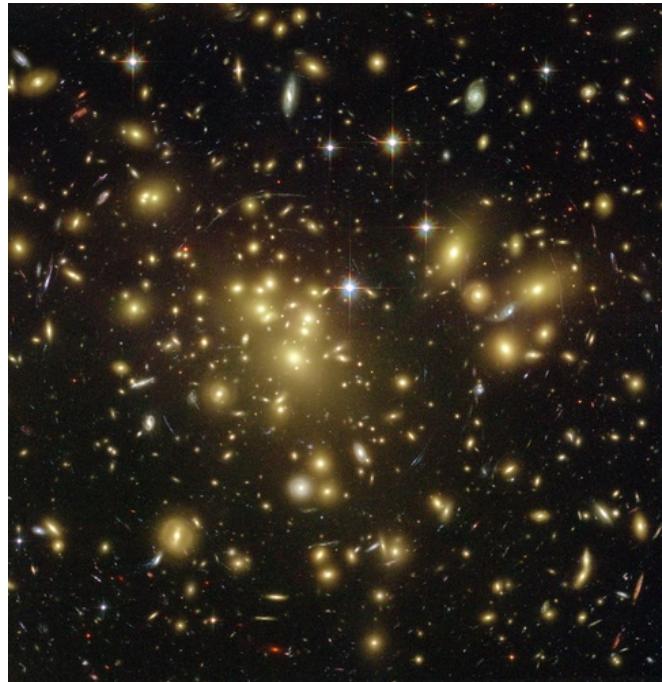
$$kT_{\text{gas}} \approx \mu m_p \sigma^2 \quad T_{\text{gas}} = 10^7 - 10^8 \text{ K}$$

$$\rho_{\text{gas}} \approx 10^{-1} - 10^{-4} \text{ cm}^{-3} \quad \text{Bremsstrahlung emissivity: } \quad \epsilon \propto n_e n_i \sqrt{T} e^{-h\nu/kT}$$

X-ray emission $\sim n_e^2$ unambiguous tracer of deep potential well ie., clusters

What is a galaxy Cluster?

Abell 1689 RC =4



Richness Classes

$M_3 + 2$

0	30-49
1	50-79
2	80-129
3	130-199
4	200-299
5	300>

Sizes 1-3 Mpc

$$\sigma \approx 300 - 1000 \text{ km s}^{-1}$$

$$M_{cl} \approx 10^{13} - 10^{15} M_\odot$$

$$M_{vir} \approx 3\sigma^2 r/G \approx 10^{15} h^{-1} M_\odot$$

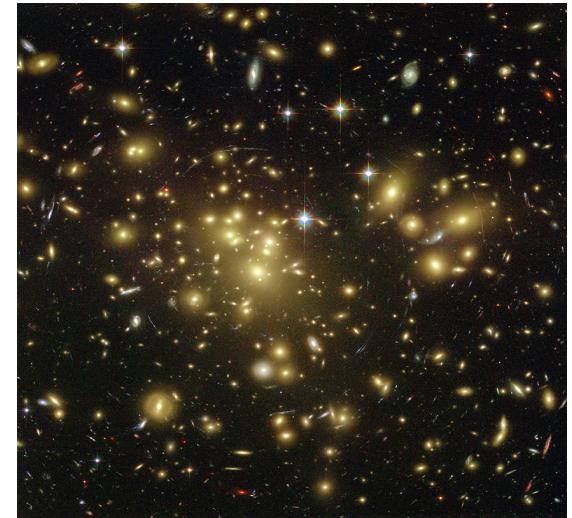
$$\text{for } \sigma \approx 1000 \text{ km s}^{-1}$$

Gravitationally bound systems

- Crossing time for the galaxies:

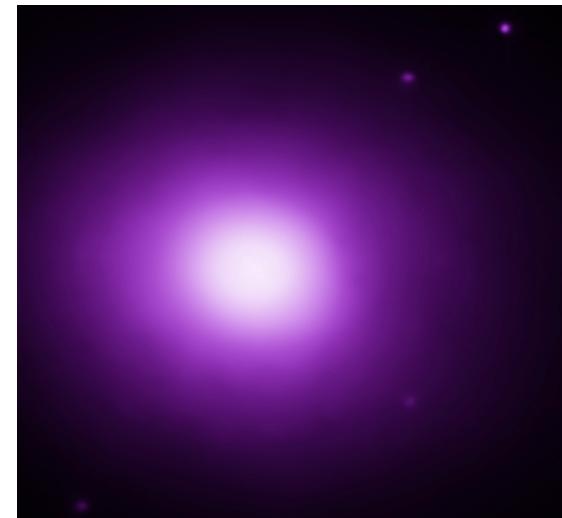
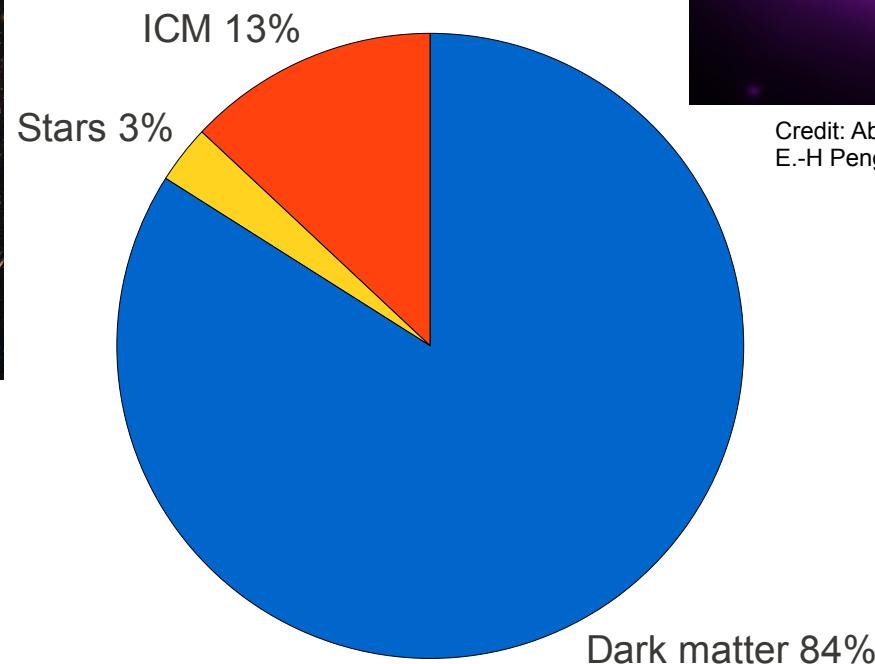
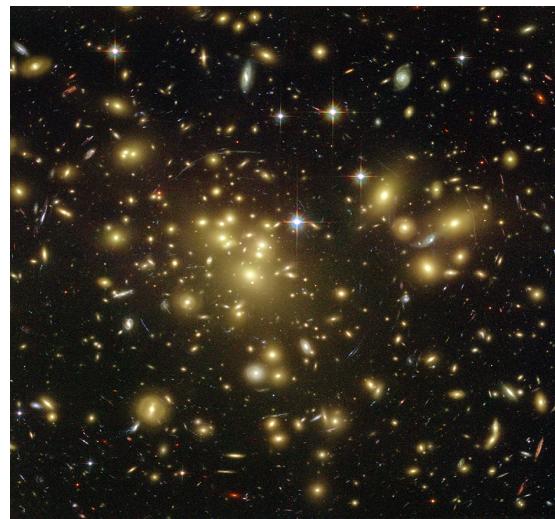
$$t_{cr, gal} = \frac{R}{\langle v^2 \rangle^{1/2}}$$

$$t_{cr, gal} \approx 2 \times 10^9 \text{ yr}$$



→ clusters are **gravitationally bound**

Galaxy cluster mass partitioning



Credit: Abell 1689, X-ray: NASA/CXC/MIT/
E.-H Peng et al.; Optical: NASA/STScI

?

Virial Theorem

- For gravitationally bound systems:

$$2E_K + E_P = 0$$

$$M_{tot} = \frac{\langle v^2 \rangle R_{tot}}{G}$$

- Assuming random uncorrelated velocity vectors:

$$M_{tot} = \frac{3\sigma_r^2 R_{tot}}{G}$$

- For velocities 1000km/s and $R = 1\text{Mpc}$ $\Rightarrow M_{tot} = 10^{15}\text{Msun}$

Sound crossing time

- Sound crossing time for X-ray gas:

$$t_{cr, \text{gas}} = \frac{D}{c_s}$$

- For an ideal gas the adiabatic sound speed is:

$$c_s^2 = \frac{dP}{d\rho} = \gamma \frac{P}{\rho} = \gamma \frac{k_B T}{\mu m_H}$$

- For a monatomic gas the adiabatic index is 5/3.

$$t_{cr, \text{gas}} \approx \text{few} \times 10^8 \text{ yr}$$

- Much shorter than age of cluster $\sim 10^{10}$ yrs

Hydrostatic Equilibrium

- Hydrostatic equilibrium: compression due to gravity is balanced by a pressure gradient force.

$$\nabla P = -\rho \nabla \Phi , \quad \frac{1}{\rho} \frac{dP}{dr} = -\frac{GM(r)}{r^2} ,$$

- Assumptions:
 - Sound crossing time much less than the age of the cluster
 - Timescale for elastic collisions much shorter than timescale for heating or cooling

$$t_s \equiv \frac{D}{c_s} \approx 6.6 \times 10^8 \text{ yr} \left(\frac{T}{10^8 \text{ K}} \right)^{-1/2} \left(\frac{D}{1 \text{ Mpc}} \right)$$

Kinetic Equilibrium

Timescales for electrons and ions to form a Maxwellian distribution electrons via Coulomb collisions:

$$t_{\text{eq}}(e, e) \approx 3 \times 10^5 \text{ yr} \left(\frac{T}{10^8 \text{ K}} \right)^{3/2} \left(\frac{n_e}{10^{-3} \text{ cm}^{-3}} \right)^{-1} \text{ yr}$$

$$t_{\text{eq}}(p, p) \approx (m_p / m_e)^{1/2} t_{\text{eq}}(e, e) \sim 43 t_{\text{eq}}(e, e)$$

$$t_{\text{eq}}(p, e) \approx (m_p / m_e) t_{\text{eq}}(e, e) \sim 1870 t_{\text{eq}}(e, e)$$

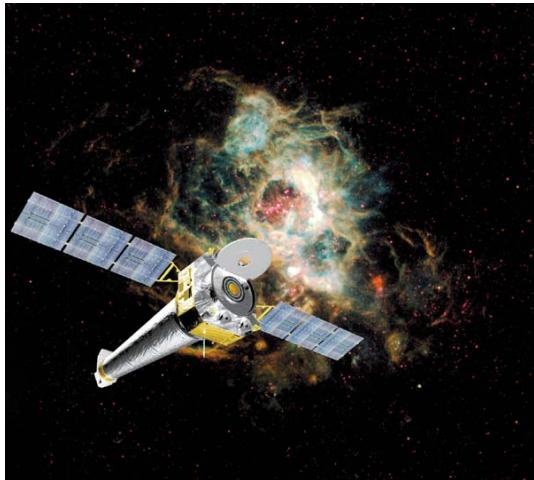
For typical conditions

$$t_{\text{eq}}(e, e) \sim 10^5 \text{ yr}, t_{\text{eq}}(p, p) \sim 4 \times 10^6 \text{ yr}, \text{ and } t_{\text{eq}}(p, e) \sim 2 \times 10^8 \text{ yr}$$

Electrons & ions in equipartition, such that $T = T_e = T_p$

Detecting Clusters with X-rays

Chandra



E = 0.5-10 keV
EA = 600 cm² @ 1.5 keV
FOV = 16.9x16.9 arcmin
PSF = 0.5 arcsec FWHM

XMM-Newton



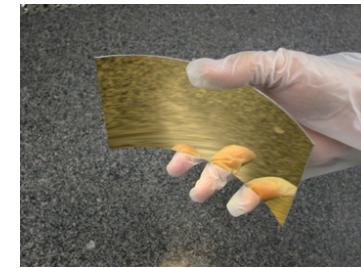
E = 0.1-15 keV
EA = 922 cm² @ 1 keV
FOV = 33 x 33 arcmin
PSF ~ 4 arcsec FWHM

Suzaku



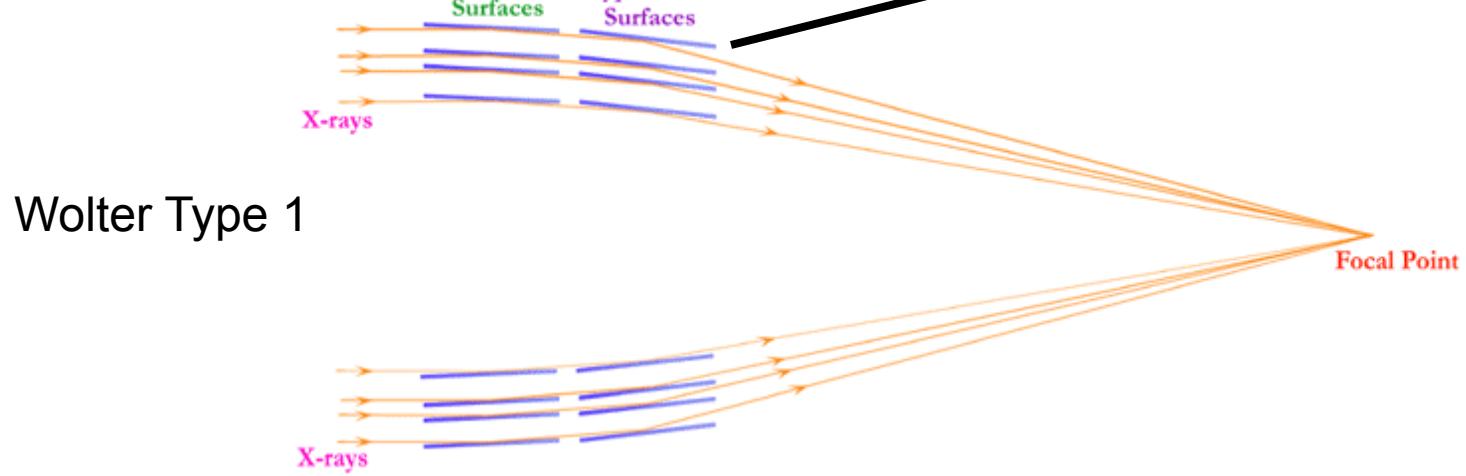
E = 0.4-600 keV
EA = 900 cm² @ 1.5 keV
FOV = 20 arcmin @ 1keV
PSF ~ 2 arcmin (HPD)

Chandra's mirrors



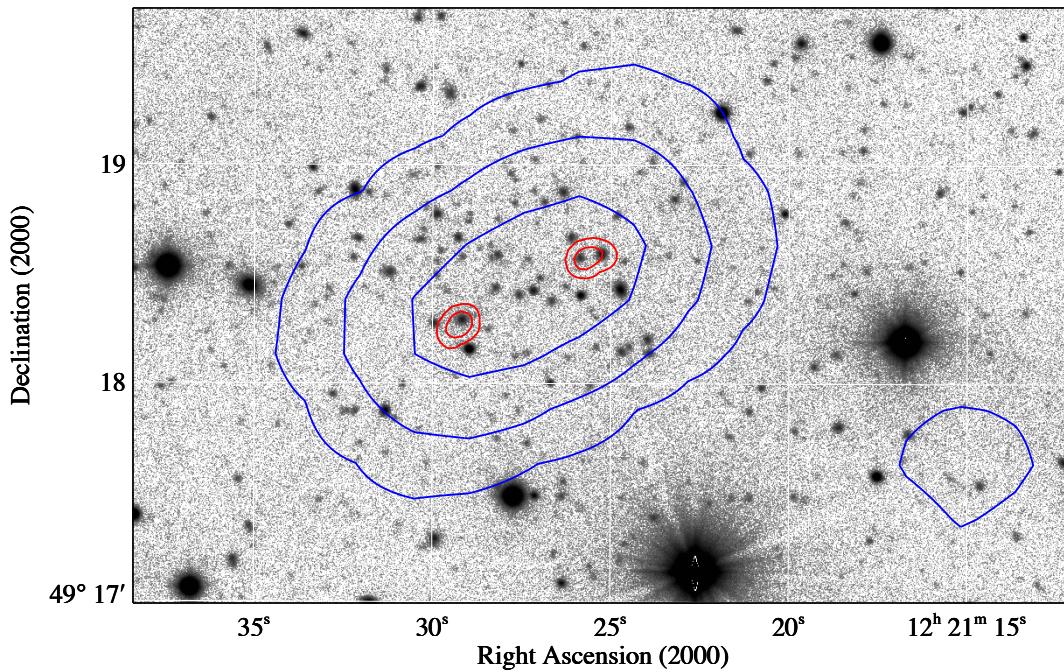
Grazing incidence optics

critical angle
few degrees



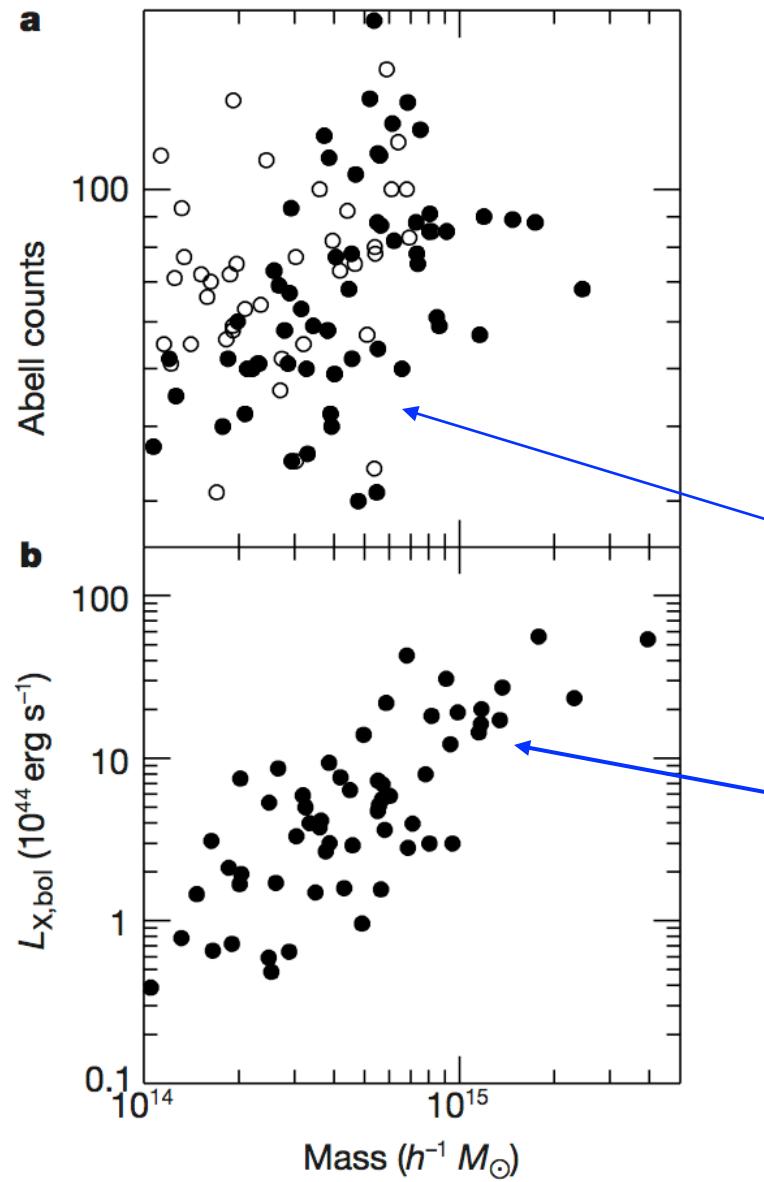
Hundreds of Clusters Detected with *ROSAT*

NRAO-VLA Sky Survey (NVSS)
ROSAT X-ray Imaging

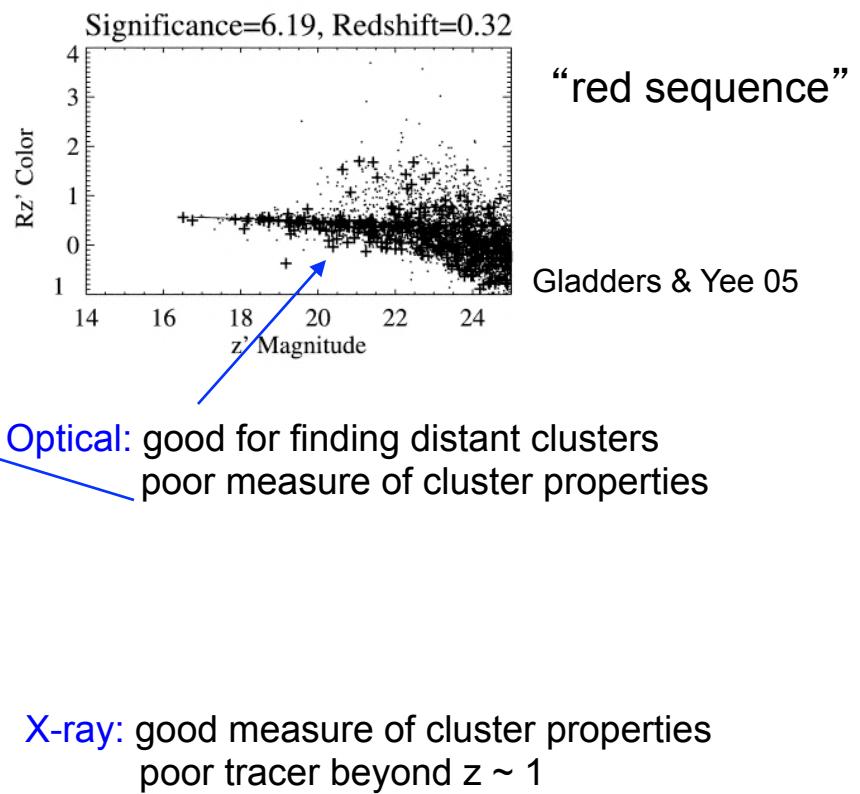


J1221+4918
 $z = 0.7$
 $L_x = 1.2 \times 10^{45} \text{ erg s}^{-1}$
 $kT = 6.5 \text{ keV}$

Ma + 11



Borgani & Guzzo 01



Thermal Bremsstrahlung

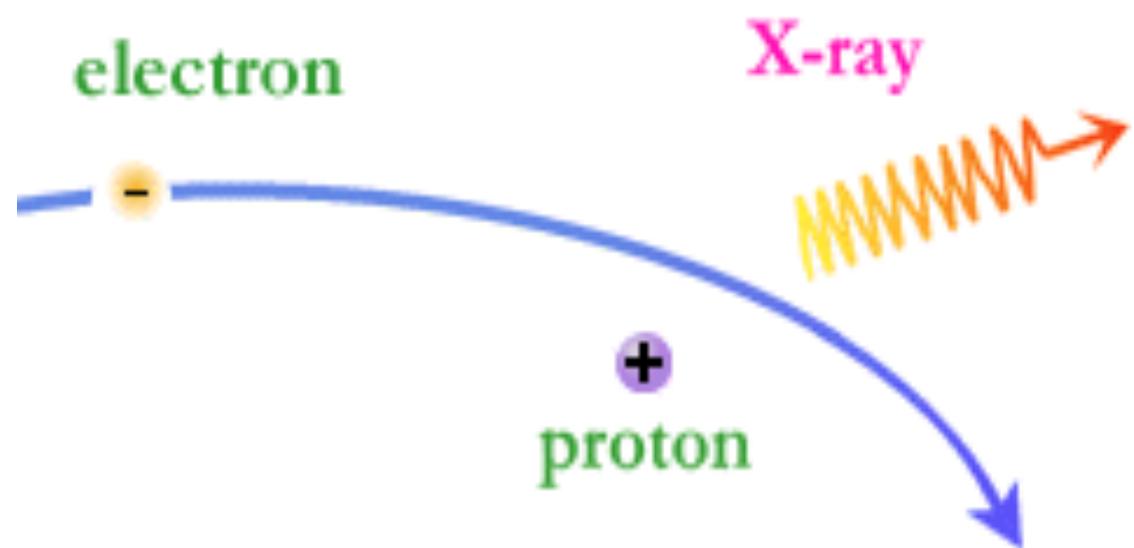
Primary source of X-ray emission from clusters due to diffuse intracluster plasma at 10^8K

Power radiated by accelerating charge:

Larmor's equation

$$P = \frac{2e^2 a^2}{3c^3}$$

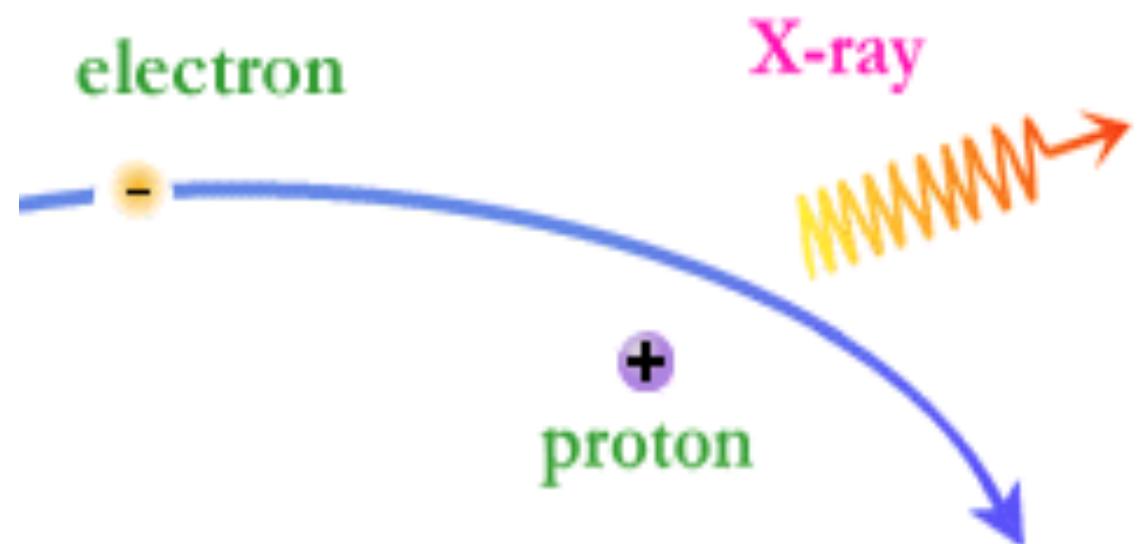
a = acceleration



Thermal Bremsstrahlung

- Power per unit frequency and per unit volume:

$$\varepsilon^{ff}(\nu, i) = \frac{2^5 \pi e^6}{3m_e c^3} \left(\frac{2\pi}{3k_B m_e} \right)^{1/2} T^{-1/2} Z_i^2 g_{ff}(Z_i, T, \nu) n_e n_i e^{-h\nu/kT}$$



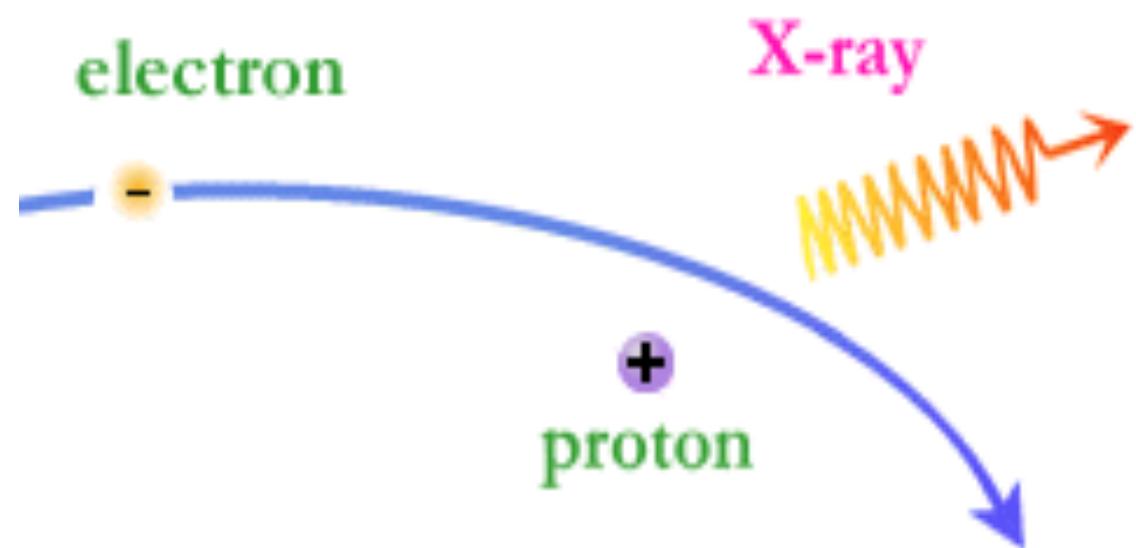
See Rybicki + Lightman or Longair

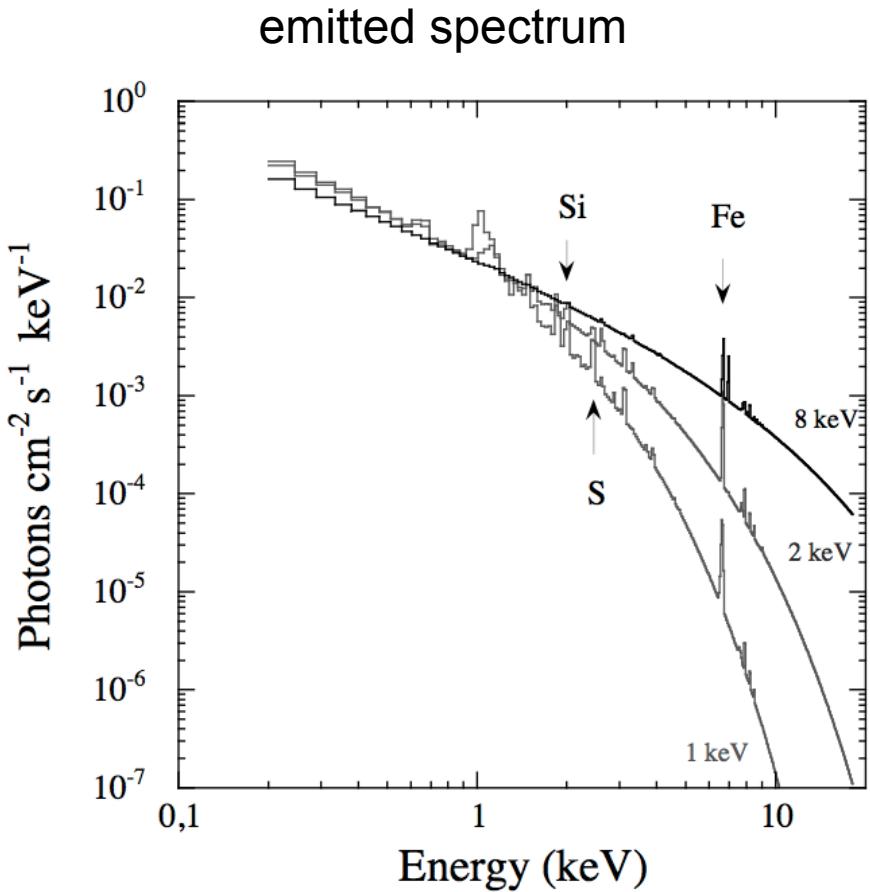
Thermal Bremsstrahlung Total Power

- Radiated power per unit volume:

$$L \propto n_e n_H T^{1/2}$$

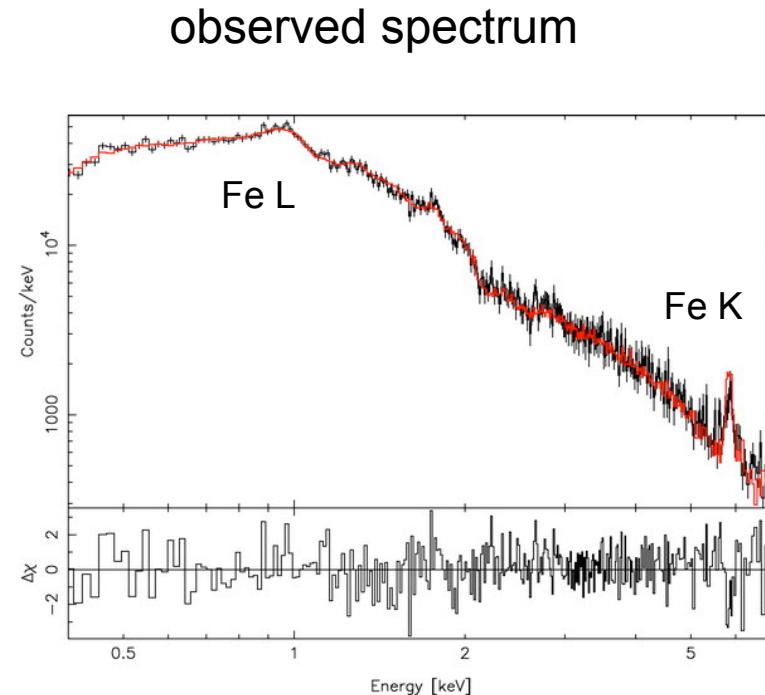
- Two body collision \rightarrow power strongly dependent on density²





$$\varepsilon \propto g(E,T,Z) T^{-1/2} \exp(-E/kT)$$

- ionization equilibrium
- temperature sensitivity
- metallicity sensitivity
- $Z \sim 1/3$ solar



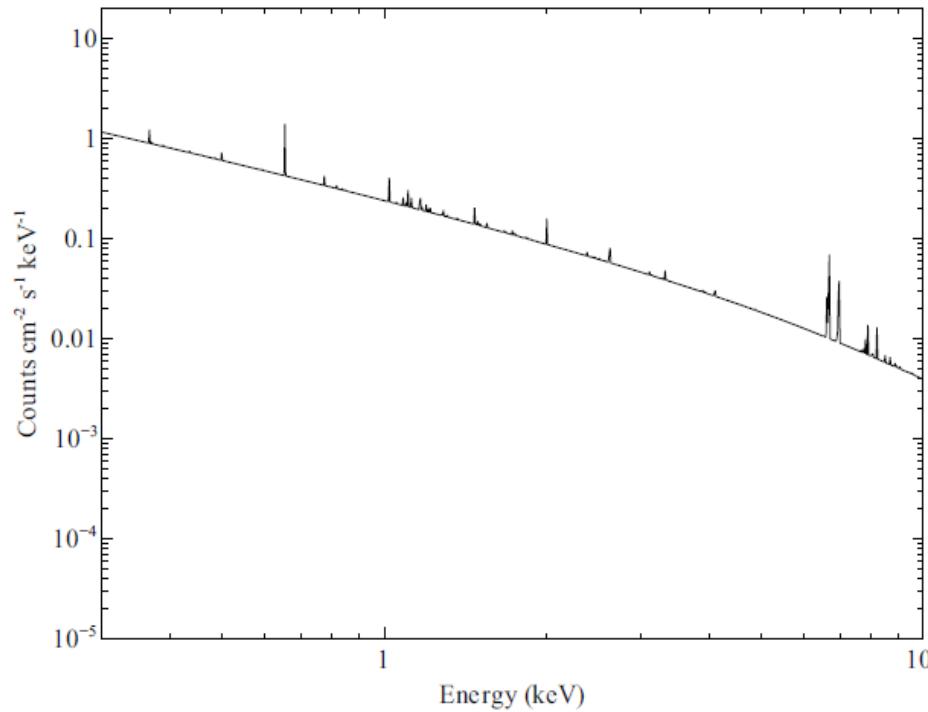
$$s \propto \varepsilon(E,T,Z) \otimes R(E,\phi)$$

Normalization: $EM \equiv \int n_p n_e dV$

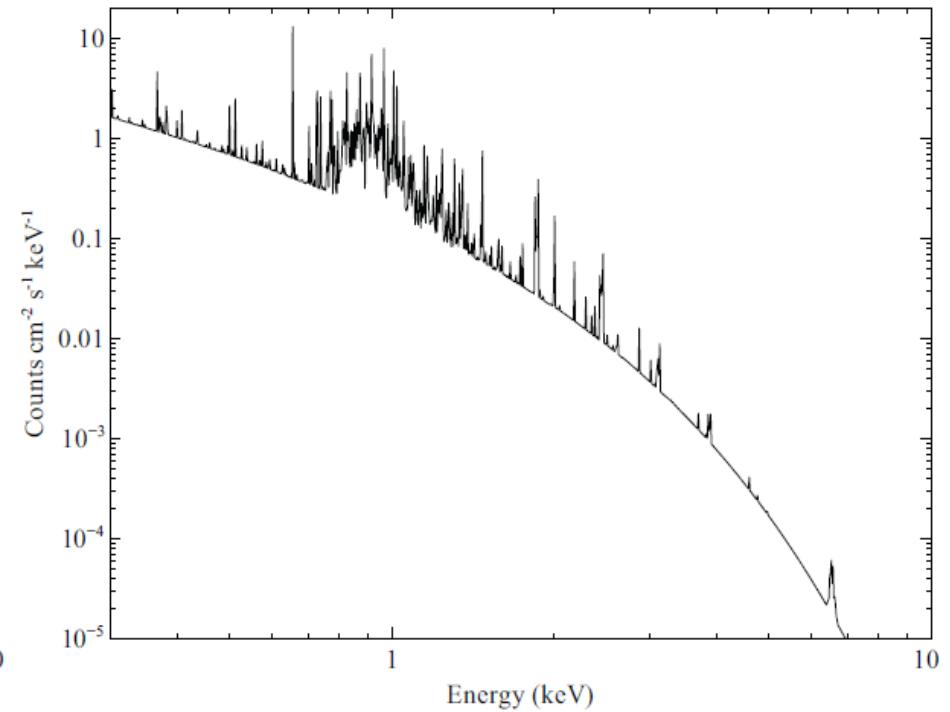
Arnaud 05

Thermal Bremsstrahlung Spectrum

- Model cluster spectrum, 1/3 solar abundance



8 keV cluster



1 keV group

$$\varepsilon^{ff}(\nu, i) = \frac{2^5 \pi e^6}{3m_e c^3} \left(\frac{2\pi}{3k_B m_e} \right)^{1/2} T^{-1/2} Z_i^2 g_{ff}(Z_i, T, \nu) n_e n_i e^{-h\nu/kT}$$

Line Emission

Emissivity from collisionally excited line (electrons)

$$\int \epsilon_{\nu}^{\text{line}} d\nu = n(X^i) n_e \frac{h^3 \nu \Omega(T) B}{4 \omega_{gs}(X^i)} \left[\frac{2}{\pi^3 m_e^3 k T} \right]^{1/2} e^{-\Delta E/kT}$$

dE - excitation energy of line above ground state

B – branching ratio (probability that excited state decays to given state)

Omega – Collision strength

omega_gs - statistical weight of ground state

Luminosity per frequency

$$L_{\nu} = \Lambda_{\nu}(T, \text{Abundances}) \int n_e n_p dV$$

Emissivity:

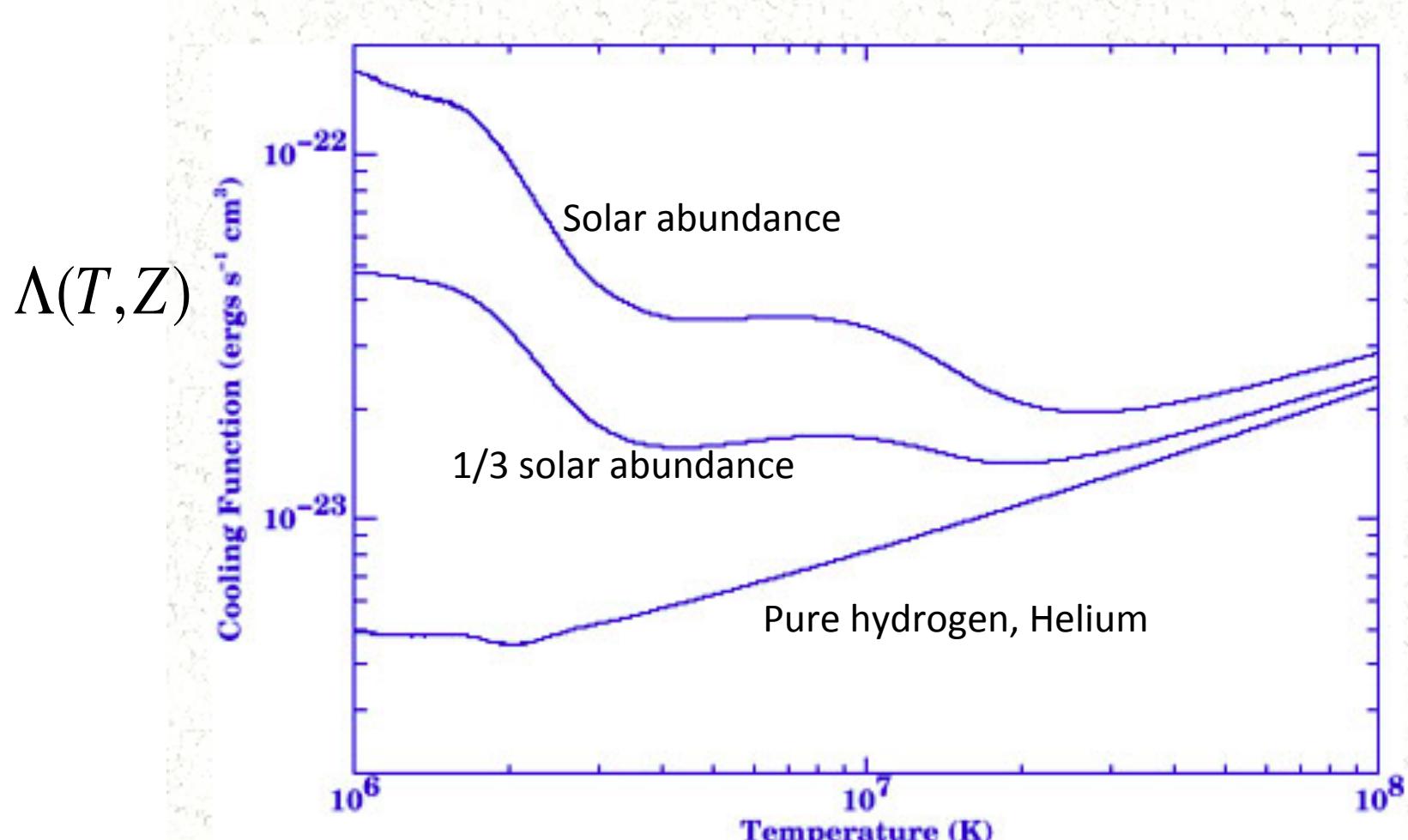
$$\epsilon_{\nu} = \Lambda_{\nu}(z, T) n_e n_p \quad (\text{erg/cm}^3)$$

Significant Soft X-ray Lines

IMPORTANT LINE BLENDS IN CLUSTER SOFT X-RAY SPECTRA

Ion	Wavelengths	Temperature	Ion	Wavelengths	Temperature
Å			keV		
Fe XXIV	10.6, 11.2	0.9 → 4.0	O VIII	19.0, 16.0	>0.2
Fe XXIII	11.0, 11.4	0.8 → 2.0	O VII	21.6, 22.0	0.1 → 0.2
	12.2		Si XIV	6.2	>1.0
Fe XXII	11.8, 12.2	0.6 → 1.5	Si XIII	6.6, 6.7	0.2 → 1.0
Fe XXI	12.2, 12.8	0.5 → 1.0	Al XIII	7.2	>1.2
Fe XX	12.8, 13.5	0.4 → 1.0	Al XII	7.8, 7.9	0.3 → 1.2
Fe XIX	13.5, 12.8	0.3 → 0.9	Mg XII	8.4	>0.7
Fe XVIII	14.2, 16.0	0.3 → 0.8	Mg XI	9.2, 9.3	0.1 → 0.6
Fe XVII	15.0, 17.1	0.2 → 0.6	Ne X	12.2	>0.4
	15.3, 16.8		Ne IX	13.5, 13.7	0.1 → 0.3
N VII	24.8	>0.1			
C VI	33.7	>0.1			

X-ray cooling function

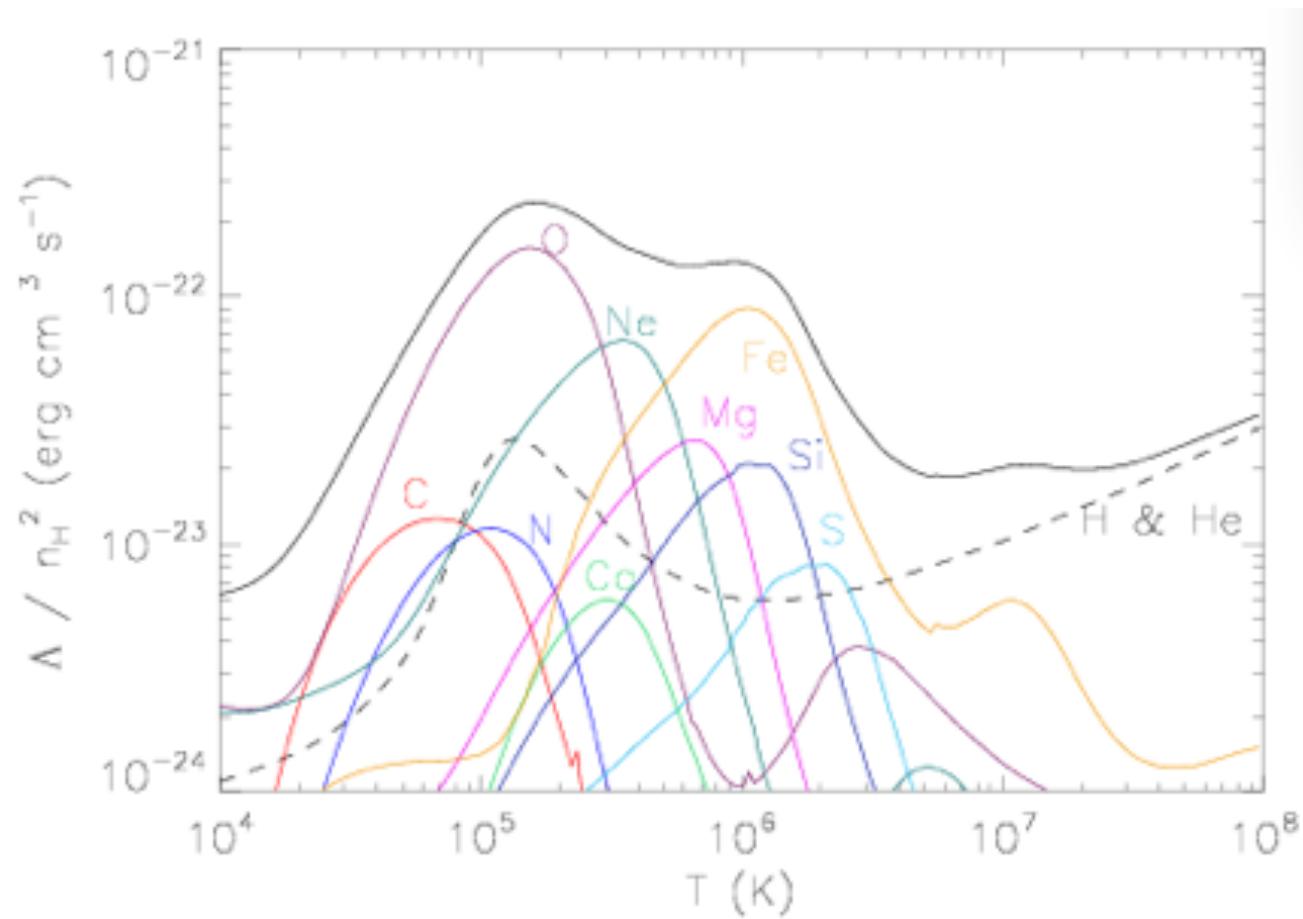


$$\epsilon_\nu = \Lambda_\nu(z, T) n_e n_p$$

Peterson & Fabian 06

X-ray cooling function

$$\Lambda(T, Z)$$



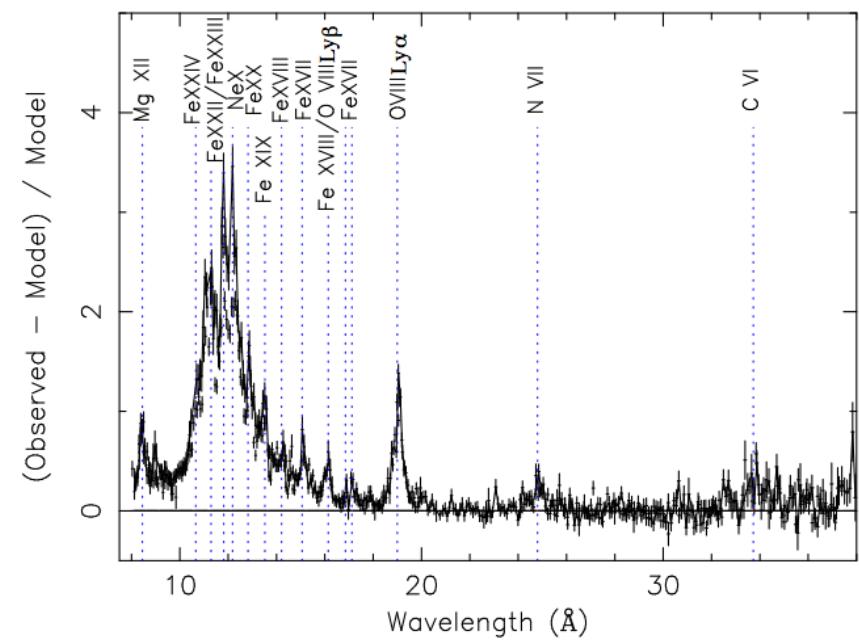
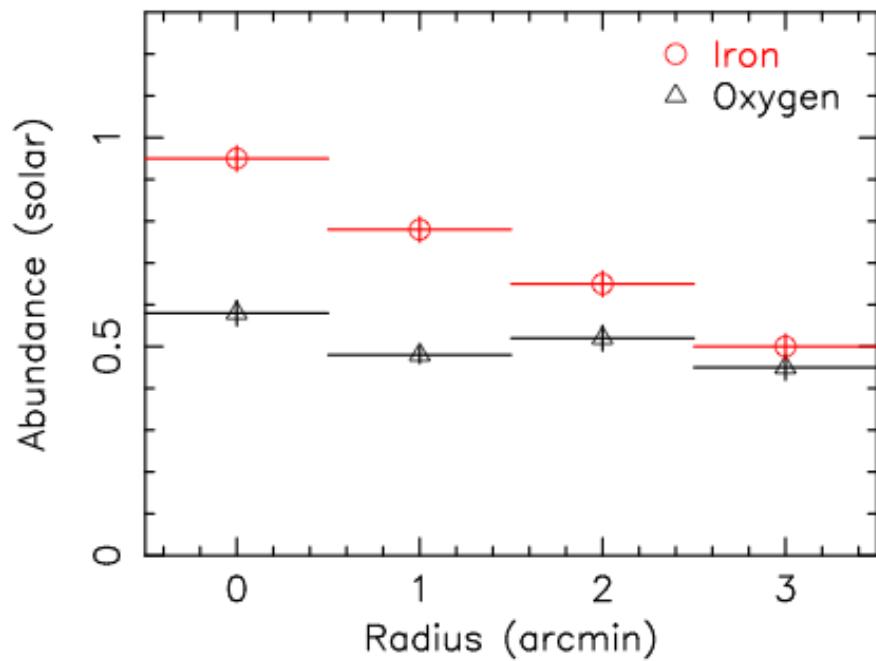
$$\varepsilon_\nu = \Lambda_\nu(z, T) n_e n_p$$

Abundances from X-ray emission

- Line emissivity proportional n_e^2 and elemental abundance
- Thermal bremsstrahlung proportional to n_e^2
- ratio of line emission to brems continuum indep of density
- Temperature sensitive to brems shape and line ratios (low T)
- Hence, line equivalent width gives elemental abundance

M87 Observed with the Reflection Grating Spectrometer

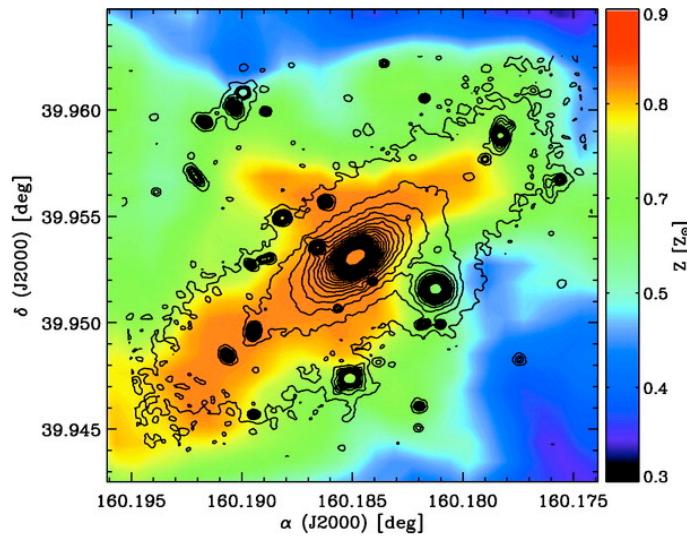
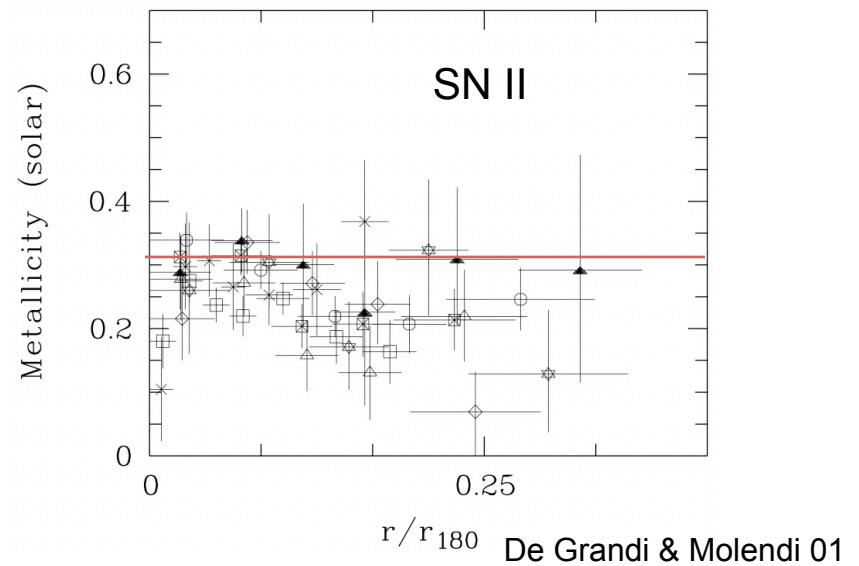
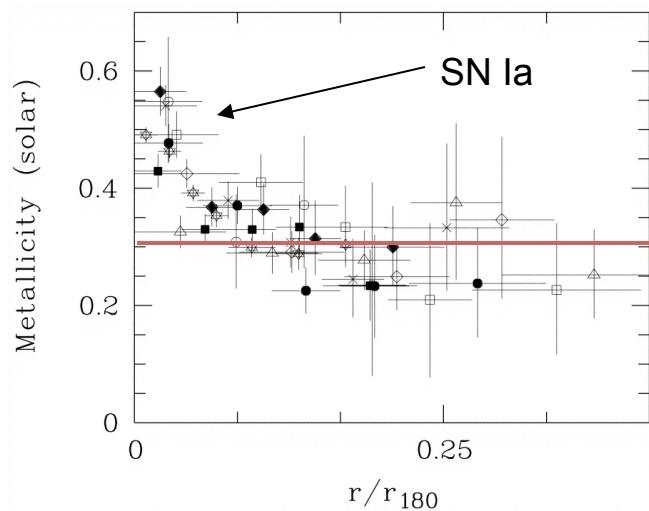
Early metal enrichment by core collapse SN, ongoing Sn Ia



Werner .. Kaastra, De Plaa...06

60% core collapse
40% Sn Ia

Metallicity Distribution



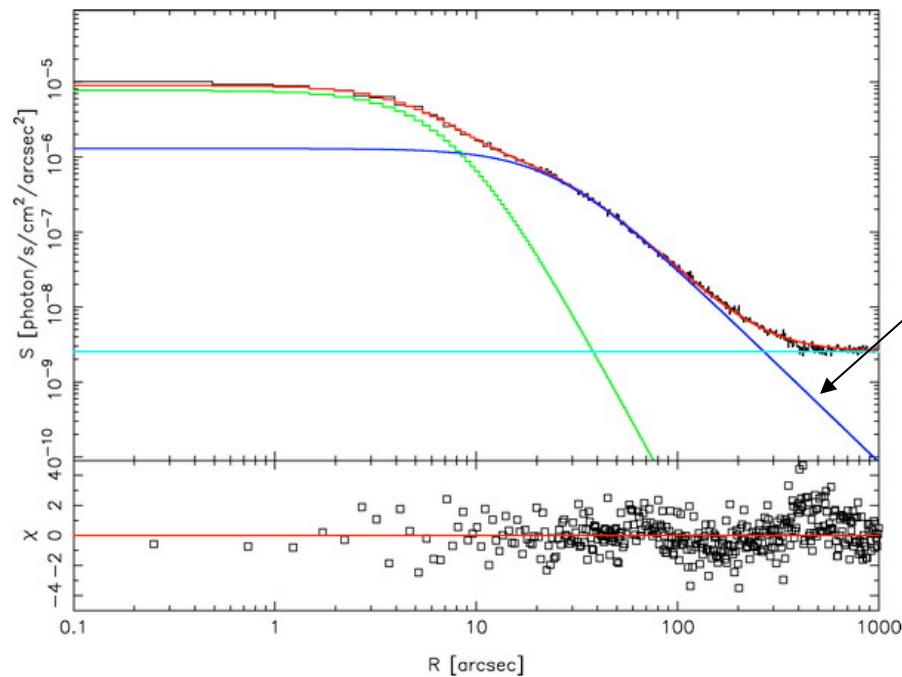
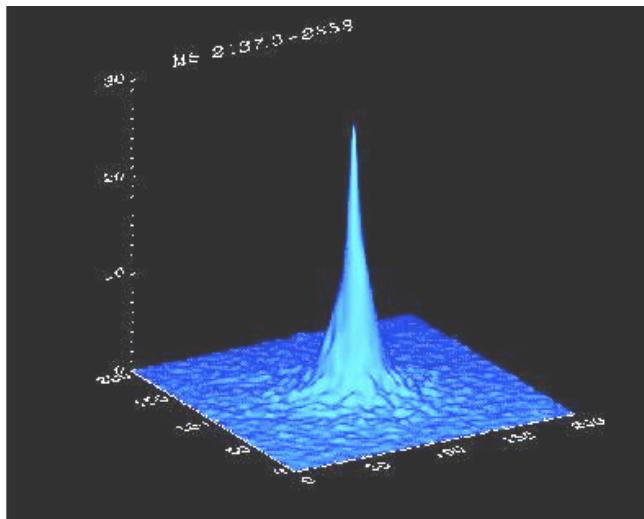
Wise et al. 04

α element enhanced (Si, S, Ne, Mg)

early enrichment $z \gg 0.5$

top-heavy IMF?

X-ray surface brightness profile



Isothermal Beta Model: Cavaliere & Fusco-Femiano (76)

surface brightness: $I_x \propto [1 + (r/r_c)]^{-3\beta}$ $\beta = \frac{\mu m_p \sigma_r^2}{kT}$ $\beta \approx 2/3$

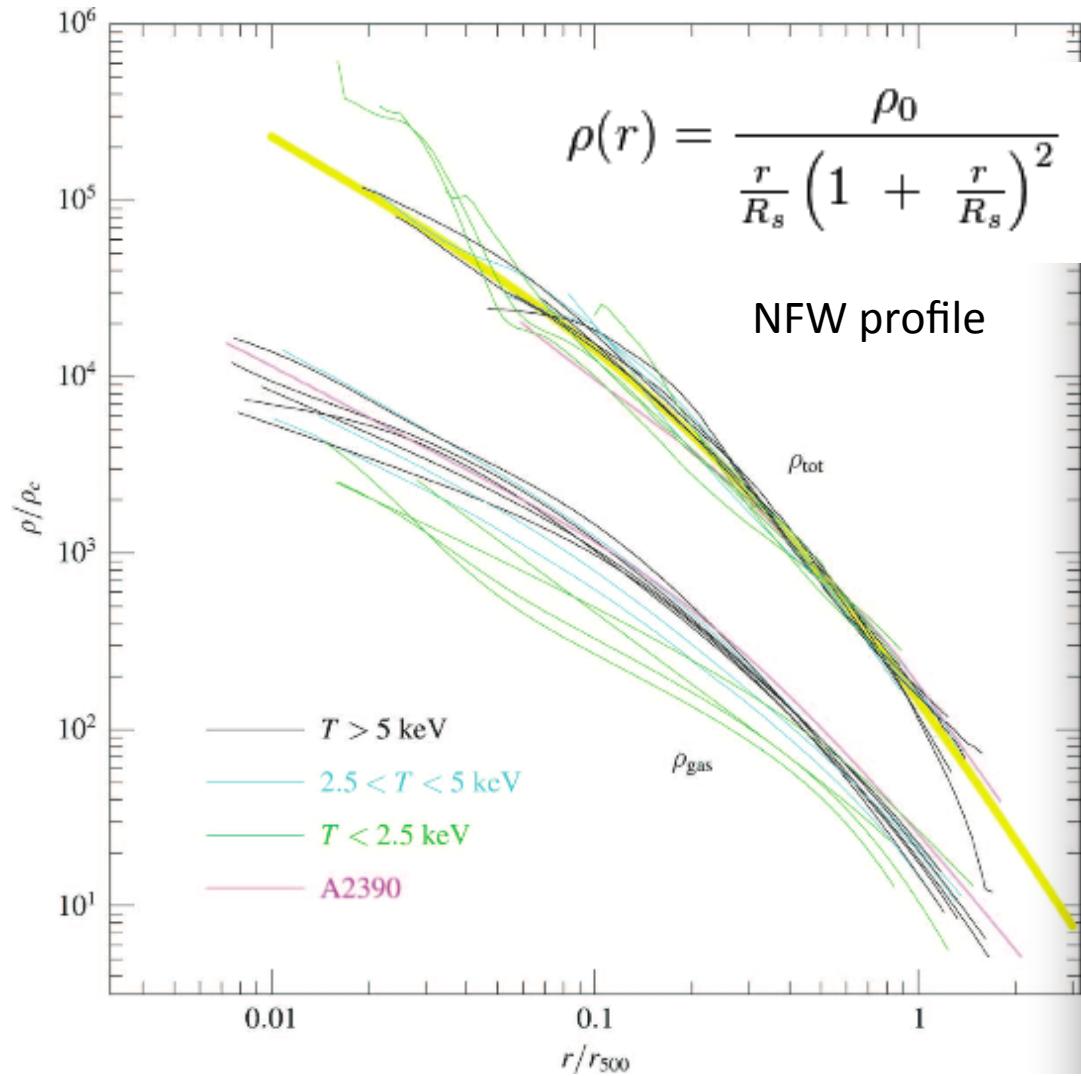
electron density: $n_e(r) = n_0 [1 + (r/r_c)^2]^{-3\beta/2}$

ratio of energy
per unit mass in
galaxies to that in
gas

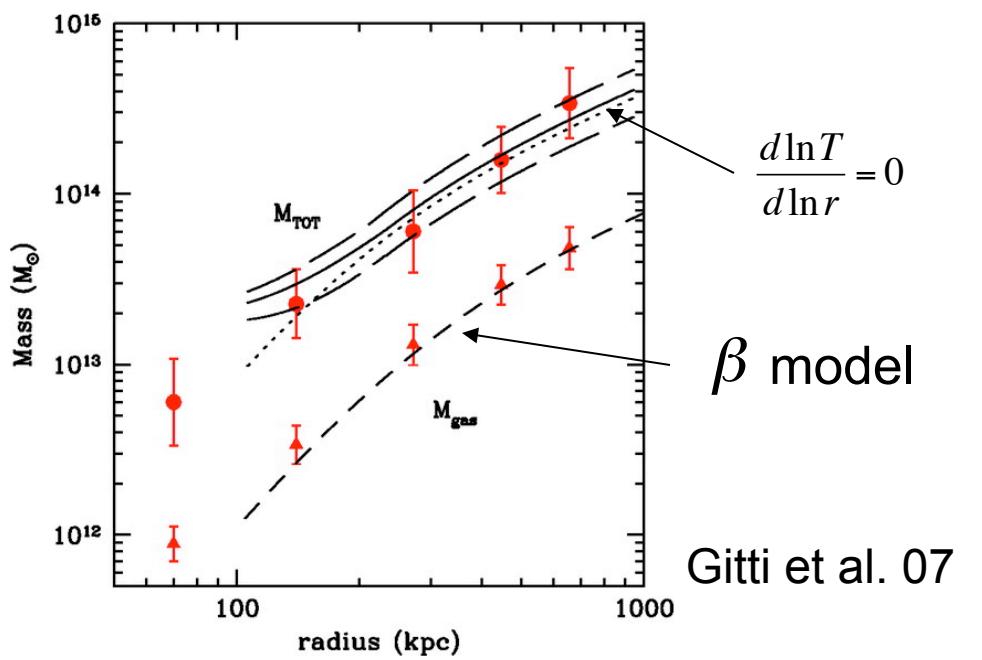
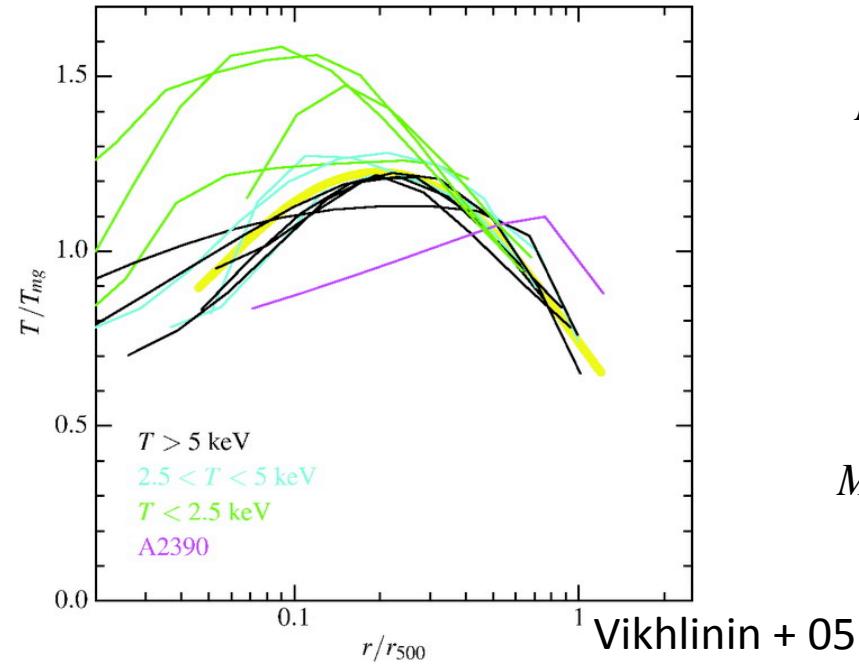
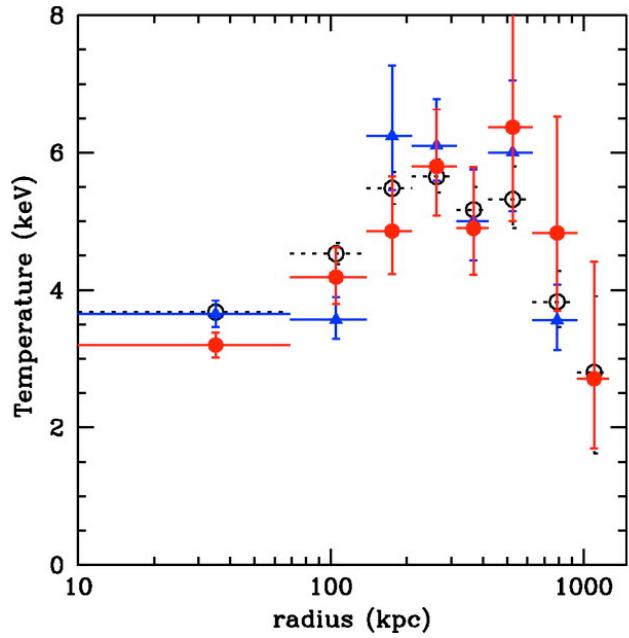
$$I_c = 0.35 I_0 \quad n_c = 0.5 n_0$$

See Vikhlinin + 06 for more sophisticated models

Cluster Density Profiles - Vikhlinin 06



Mass Profiles



Gitti et al. 07

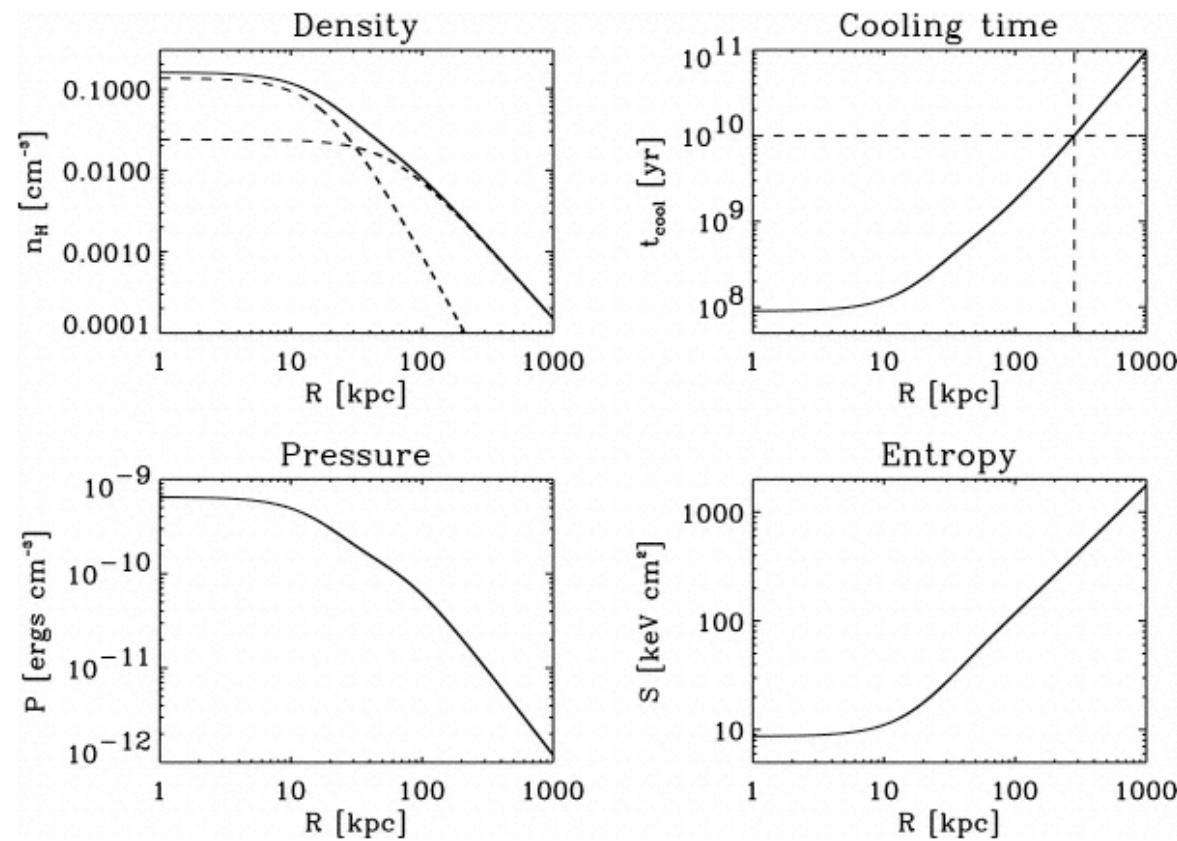
$$M_{tot}(< r) = \frac{kTr}{G\mu m_p} \left(\frac{d \ln \rho}{d \ln r} + \frac{d \ln T}{d \ln r} \right)$$

$$\rho(r) = \rho_0 \left[1 + \left(\frac{r}{r_c} \right)^2 \right]^{-3\beta/2}$$

$$M_{tot}(< r) = \frac{kr^2}{G\mu m_p} \left(\frac{3\beta rT}{r^2 + r_c^2} - \frac{dT}{dr} \right)$$

Assume
 - hydrostatic equil.
 - spherical symmetry
 - only gas pressure

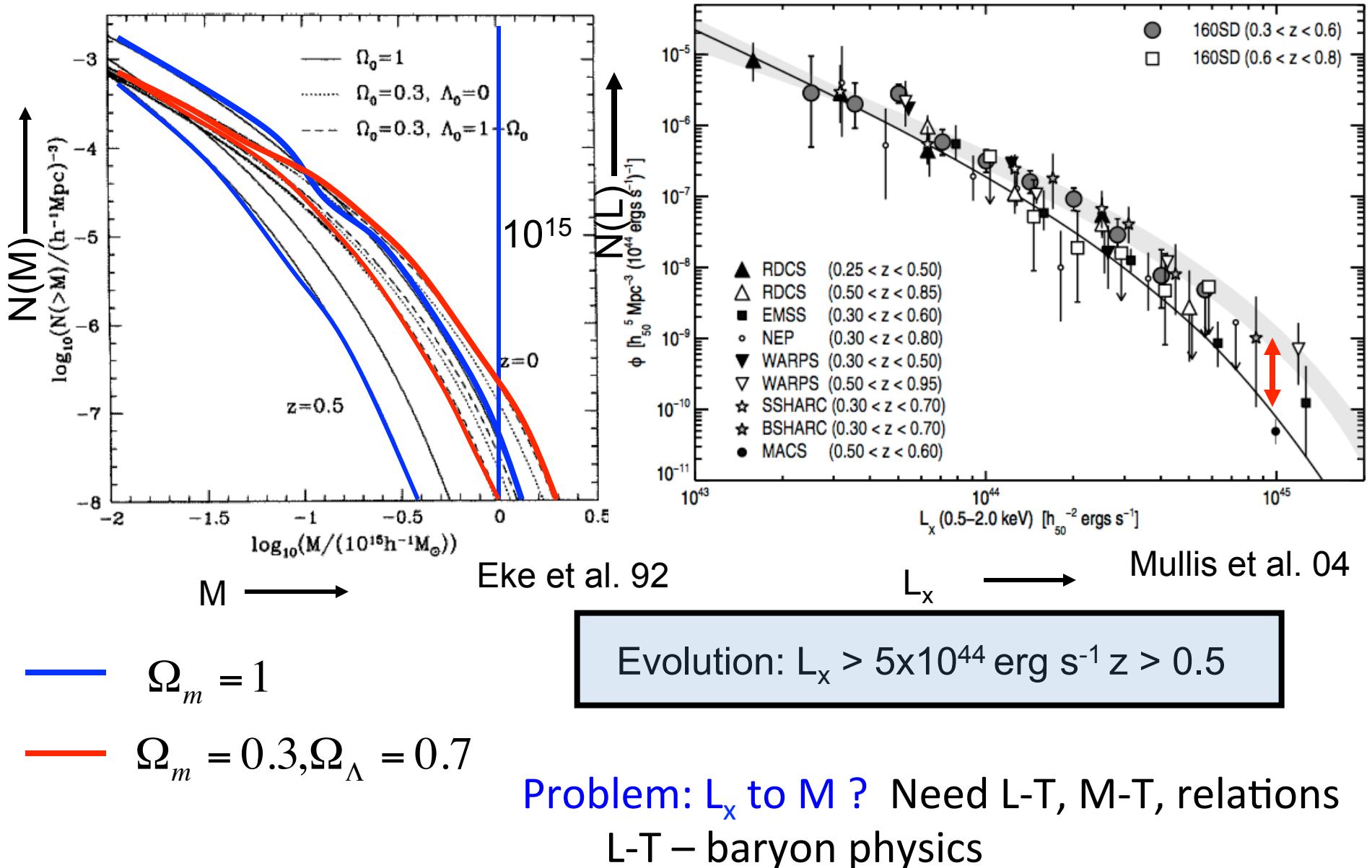
Typcal Gas Profiles



H

Clusters & Cosmology

mass & luminosity function sensitivity to $\Omega_{m,\Lambda}$



Cluster Cosmology requires: Simulations + Observations

Is cluster matter distribution only a function of **Mass and Redshift?**



Compare observations to scaling relations

Deviations: models deal primarily with DM mass
observations deal with luminous baryons
simulated & observed parameters (eg. L,T) do not
always probe same quantity

Mass Function constrains

$$\Omega_m, \sigma_8$$

mean matter density
in units of critical density,
amplitude of primordial
density fluctuations

strategy: find observational proxies for mass

L_x - easy to measure, sensitive to non-gravitational processes

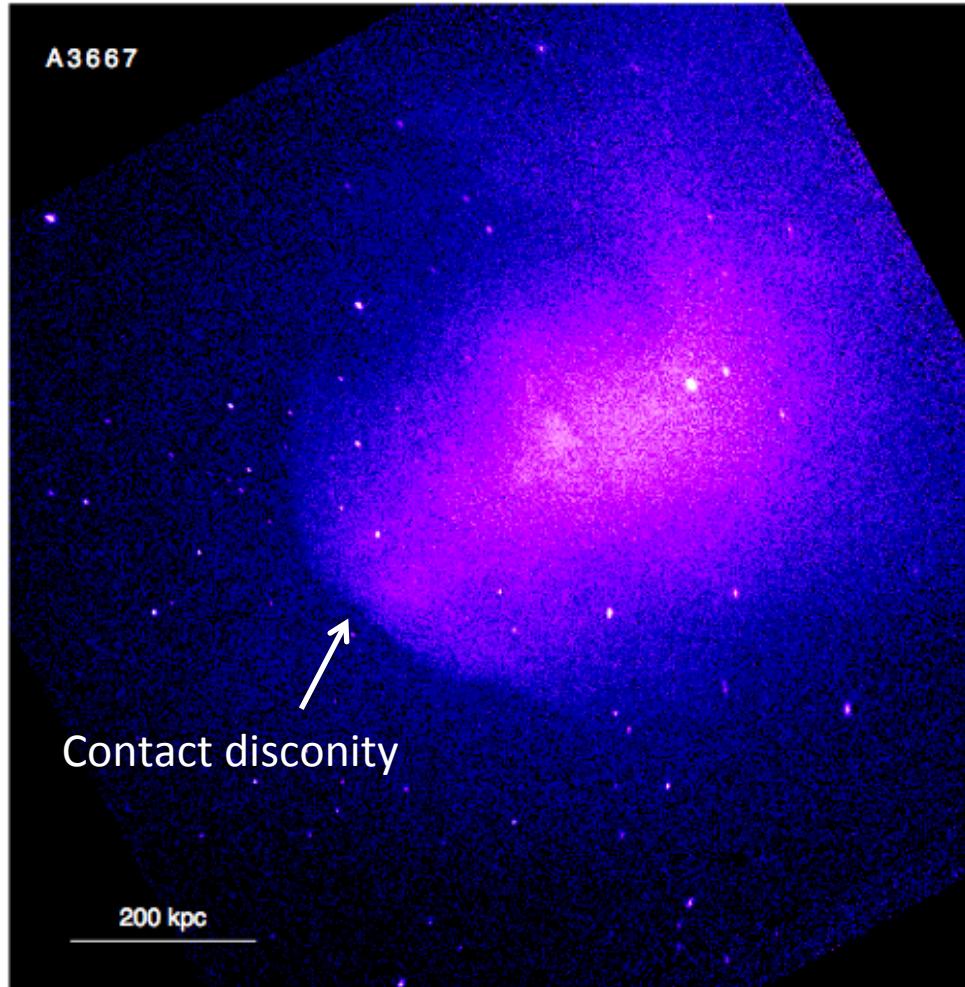
T_x - difficult to measure except for brightest objects

M - even tougher



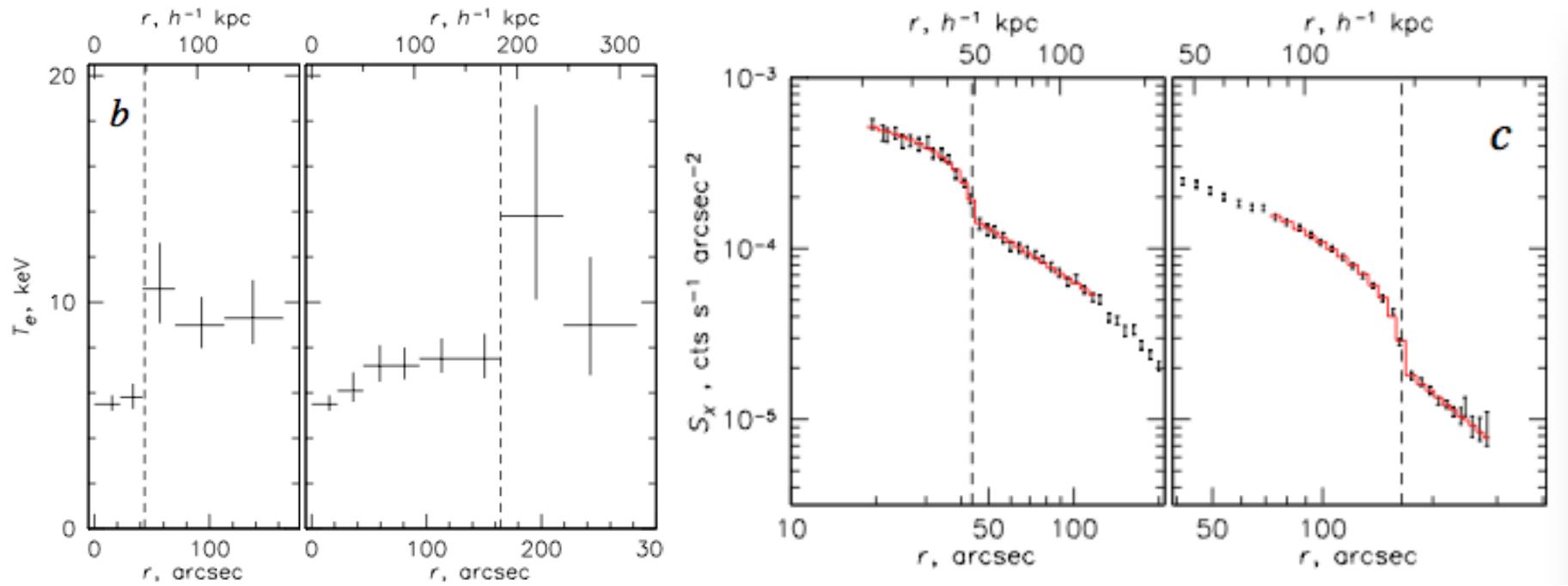
Clusters Grow by Mergers

Cold Fronts – mildly transonic mergers/sloshing



Markevich & Vikhlinin 06

Contact Discontinuity: Abell 2142



Markevich & Vikhlinin 06

Cluster Scaling Relations

Simple assumptions: gravity, adiabatic compression, shocks

self similarity: scale free, cluster defined by its **mass and redshift**

$$M \propto T_x^{3/2} (1+z)^{-1}$$

$$L_x \propto T_x^2 (1+z)^{3/2}$$

$$L_x \propto M^{4/3} (1+z)^{7/2}$$

Implications:

Same internal structure, but distant clusters denser, smaller, more luminous

Observables in reverse order of difficulty: L_x, T_x, M

Departures from scaling relations: additional physics (cooling, feedback, etc.)

Evolution of scaling laws: Ettori et al. 04

Are clusters self-similar ? Do they scale as expected? Do they evolve as expected?

Relationship between Mass, Temperature, Luminosity

examples:

$$M \propto TR$$

$$L_x \propto \rho^2 V \Lambda$$

$$R \propto T^{1/2}$$

$$M \propto \rho V \propto M_{gas}$$

$$M \propto T^{3/2}$$

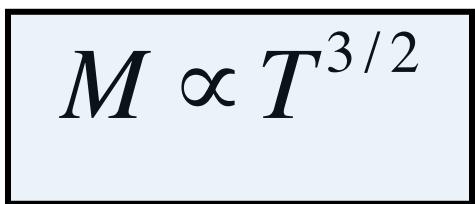
$$\Lambda \propto T^{1/2} \quad (T > 2 \text{ keV})$$

$$L_x \propto M \rho T^{1/2}$$

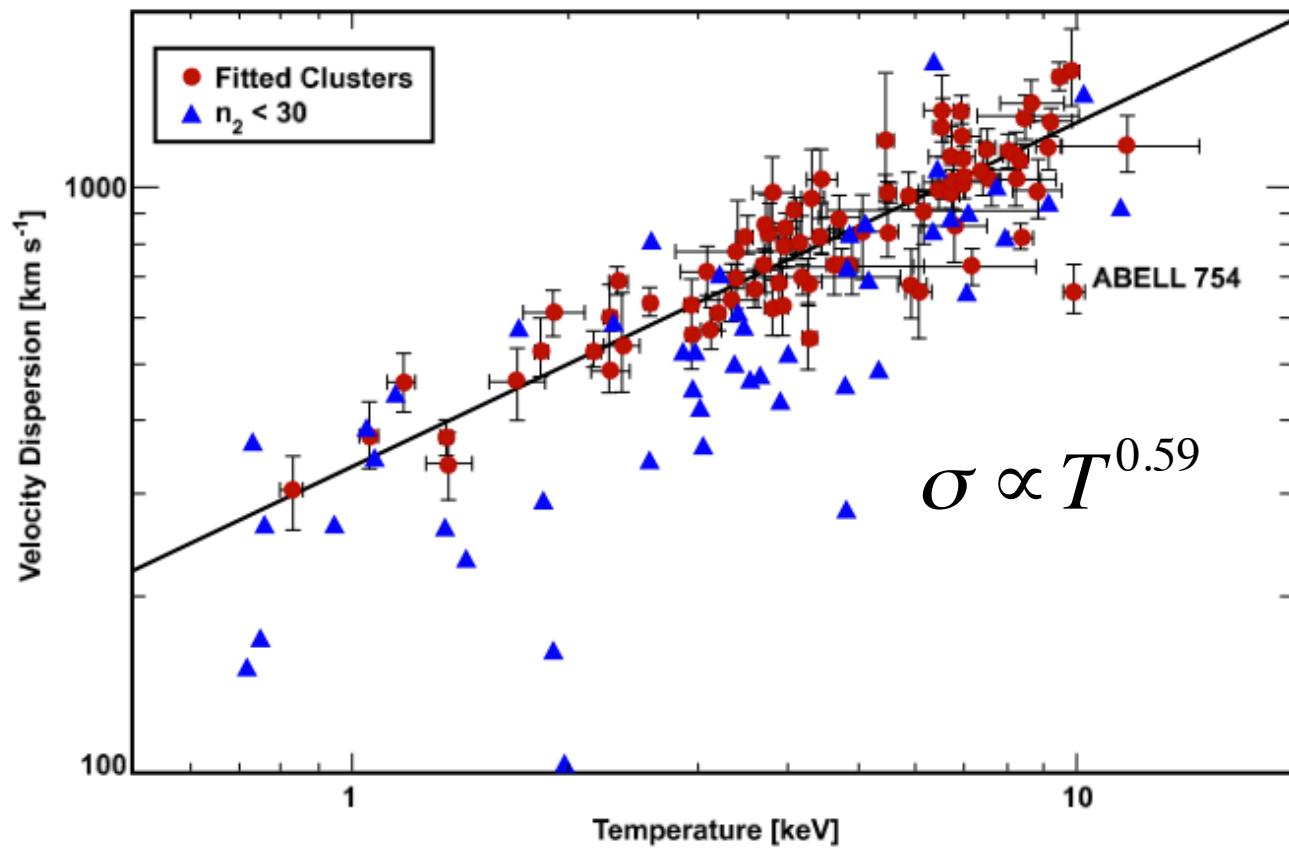
$$L \sim M^{4/3}$$

$$M \propto T^{1.5}$$

$$L_x \propto T^2$$



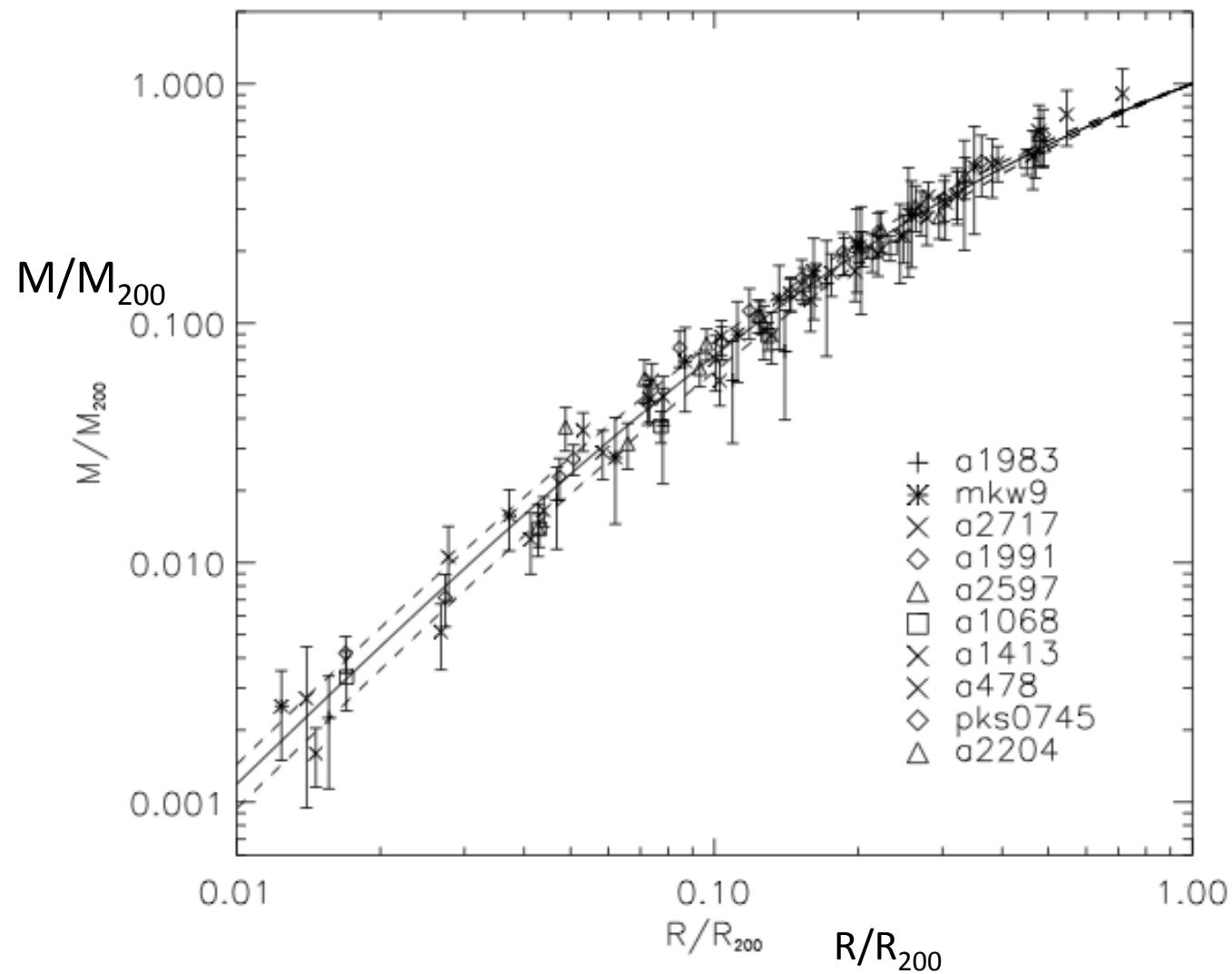
Gas Temperature is related to velocity dispersion, Mass



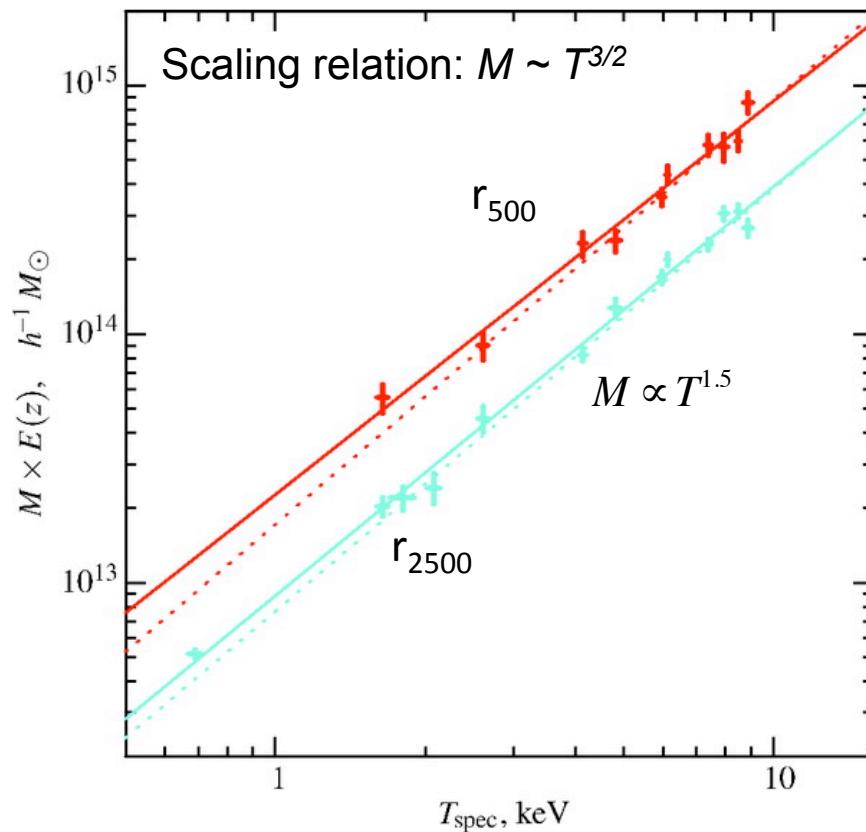
Scaling relations $\sigma \propto T^{1/2}$
Gas temperature $\Rightarrow GM/R$

Mushotzky 2004

Cluster Mass Profiles closely follow Universal NFW Profile Good News for Cluster Cosmology



Mass-Temperature Relation: Temperature good Mass proxy
 Problem: Hard to measure in abundance



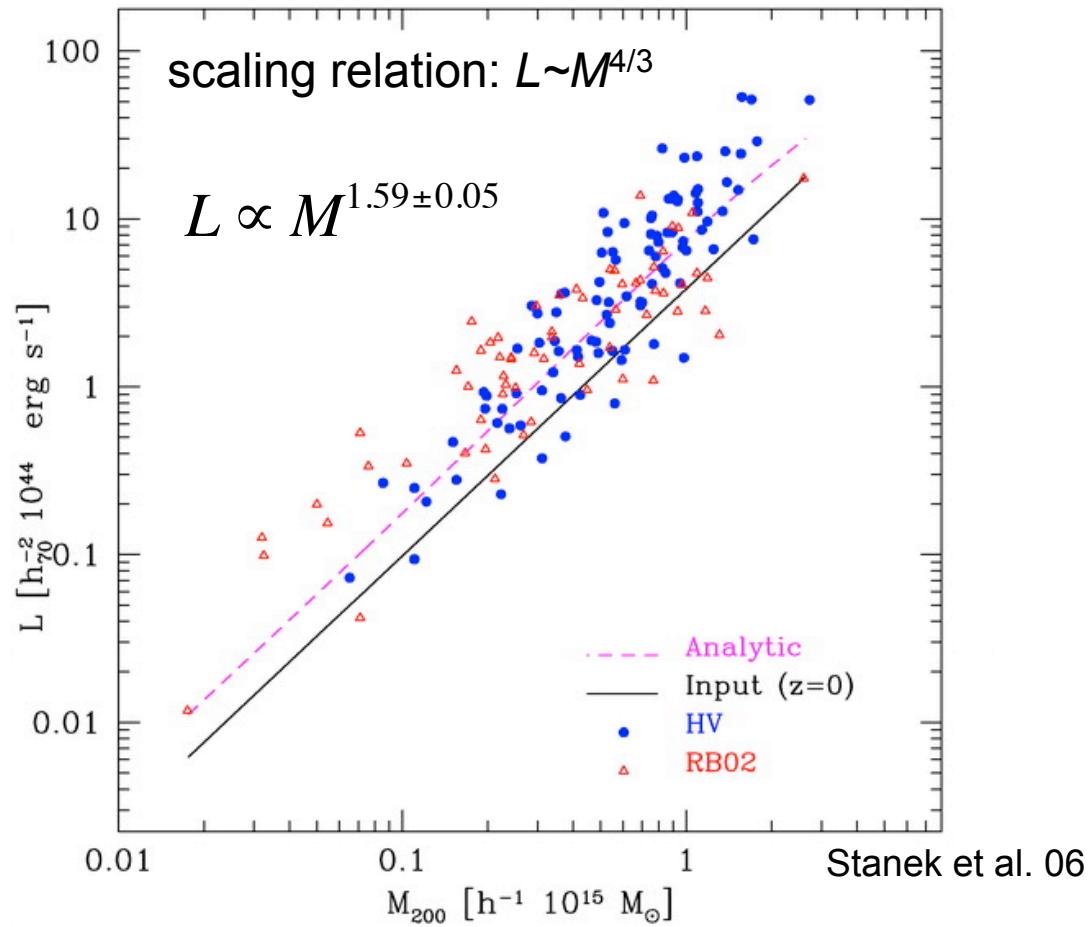
Vikhlinin et al. 06
 Ettori et al. 04
 Arnaud 05

Issues:

- small deviations from self similarity for cool clusters
- normalization varies by ~30% between studies
- least affected by non-gravitational physics
- consistent with self-similar evolution -Ettori et al. 04

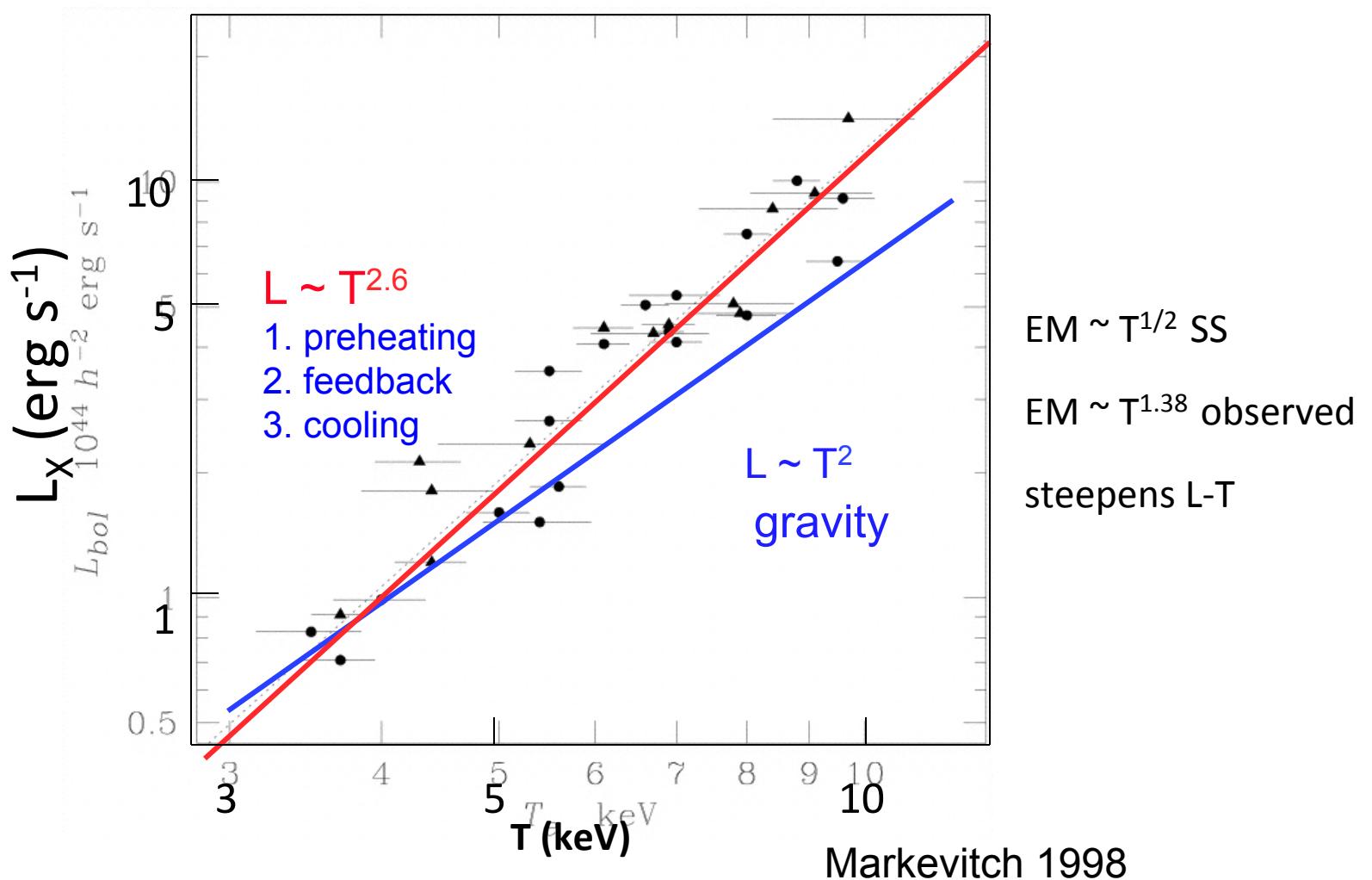
- Beta model fits
- inner core
- reference radius
- redshift range

Mass-luminosity Relation: clusters too luminous for their masses



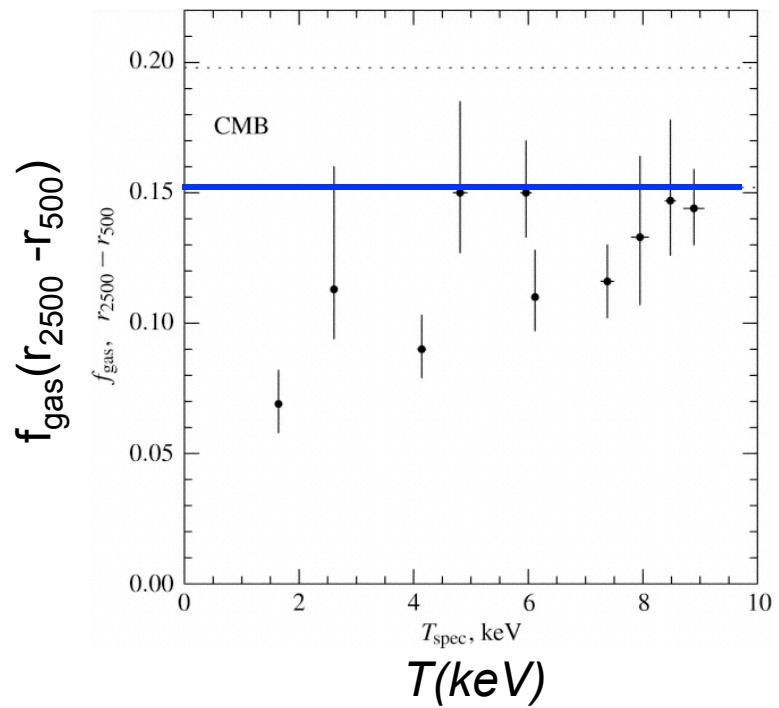
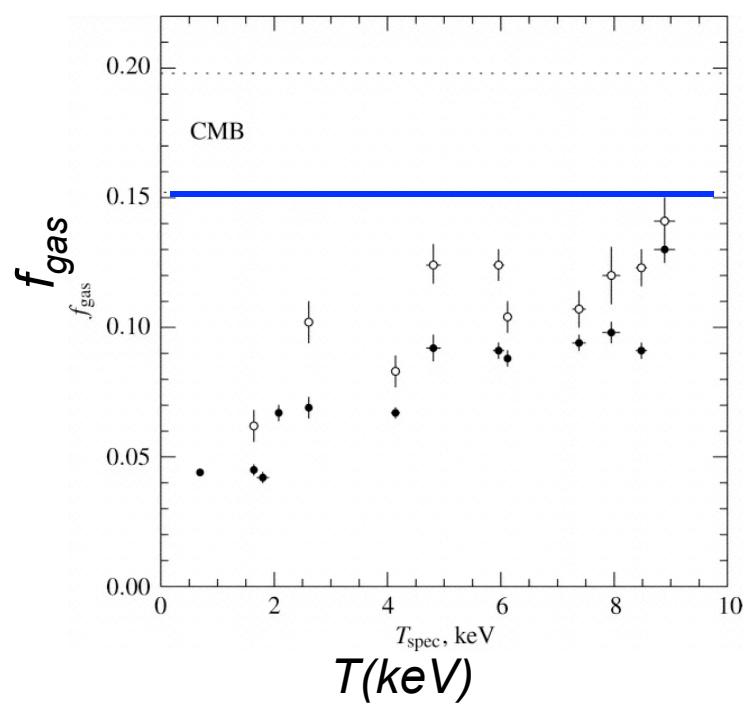
$$L \propto M^p \quad P = 1.59 \pm 0.05 \text{ Stanek et al. 06} \quad \text{too steep!}$$
$$= 1.88 \pm 0.43 \text{ Ettori et al. 04}$$

Luminosity-Temperature Relation: too steep



Evolution see Branchesi, Gioia et al. 07

Baryon fraction as cosmological probe



assumption: universal f_b

$$f_b = \Omega_b / \Omega_m$$

$$f_{\text{gal}} \sim 0.016 \quad \therefore f_{\text{gas}} / f_{\text{gal}} \approx 10$$

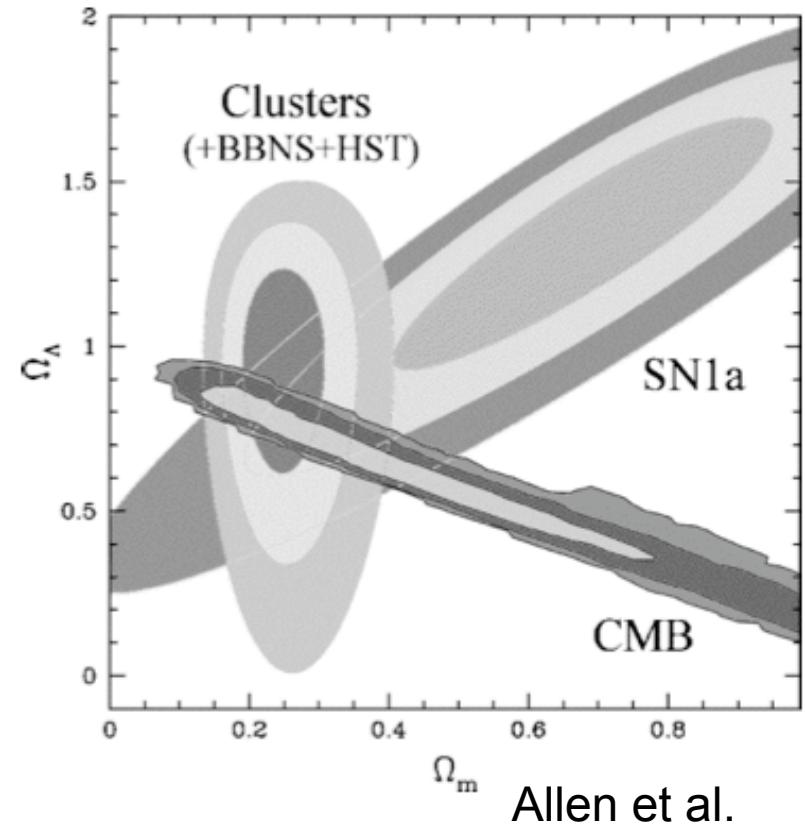
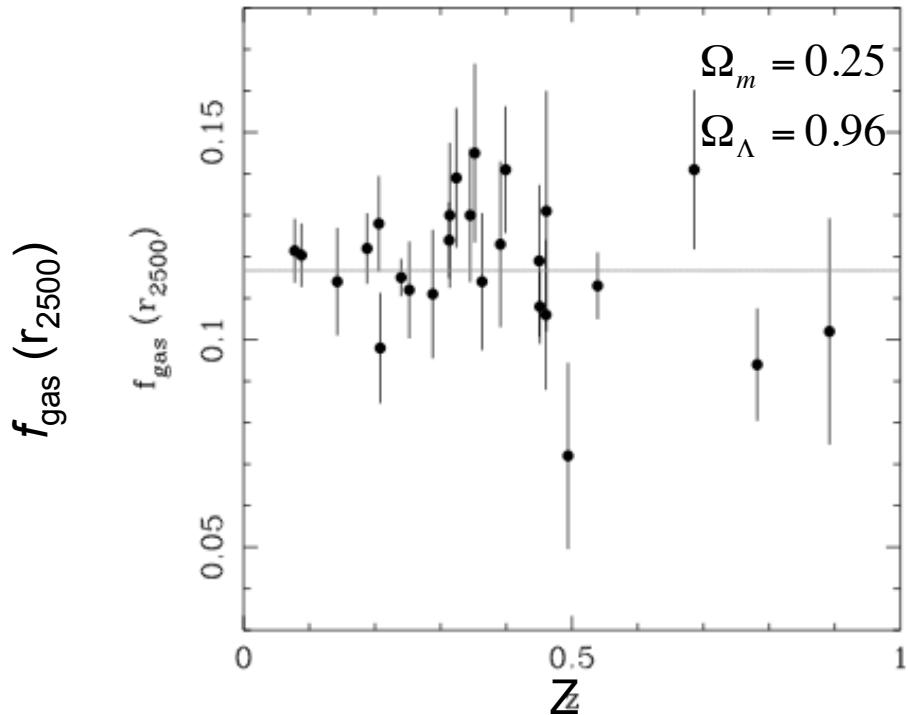
$$f_b = f_{\text{gas}} + f_{\text{gal}}$$

$$\Omega_b = 0.045 h_{70}^{-2} \quad (\text{from BBN \& WMAP})$$

Vikhlinin et al. 06
Ettori & Fabian 99
Ettori et al. 03

$$\Omega_m = \Omega_b / f_b = 0.30^{+0.04}_{-0.03}$$

Constraints on Ω_m, Ω_Λ from f_{gas}



assume $f_{\text{gas}} = \text{constant}$

$$f_{\text{gas}} \propto (1+z)^2 d(h(z))_A^{3/2}$$

$$h(z) = \sqrt{\Omega_m (1+z)^3 + \Omega_\Lambda}$$

- geometrical test
- precise mass estimates over z
- biases in aperture size, f constancy
- priors: $H_0, \Omega_b h^2$

Summary

Clusters are a rich area of studies at all wavelengths

- Clusters close to self-similar population (M-T)
- Dark matter profile universal: theory solid
- Gas scaling steeper than predicted:
problem: baryon physics: cooling, heating